

# Unloading the dice

Minimising biases in Full Configuration Interaction Quantum Monte Carlo

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# Part I

## Introduction: FCIQMC

# Full Configuration Interaction

- Want to solve the eigenvalue problem  $H|\Psi\rangle = E|\Psi\rangle$  where  $|\Psi\rangle = \sum_I C_I |D_I\rangle$
- Traditionally Diagonalise matrix of  $\langle D_I | \hat{H} | D_J \rangle$
- Unfortunately this turns out to be impossibly large  $\binom{M}{N_e}$  by  $\binom{M}{N_e}$ .
- **Instead use Monte Carlo to sample this large space. FCIQMC:**

G. H. Booth, A. J. W. Thom, A. Alavi, J. Chem. Phys. **131**, 054106 (2009)

# FCIQMC: Huge Systems to FCI accuracy

With an additional approximation people have studied:

- Molecular Systems with  $10^{29}$  configurations

C. Daday *et. al.*, J. Chem. Theory Comput. **8**, 4441 (2012)

- A Uniform Electron Gas with a Hilbert Space of  $10^{108}$  determinants.

J. J. Shepherd *et. al.*, Phys. Rev. B **85**, 081103 (2012)

- This would take about  $10^{109}$  bytes of storage space.
- The world wide web is about  $10^{21}$  bytes.

# Full Configuration Interaction Quantum Monte Carlo: Projector

- Apply  $e^{-(\hat{H}-S)\delta\tau}$  (stochastically) repeatedly to some initial vector  $|I\rangle$ .
- $\delta\tau$  is the time step,  $S$  is the shift (controls normalisation).
- As long as  $\langle\Psi|I\rangle \neq 0$ : the ground state emerges.
- Can use the approximation:  $e^{-(\hat{H}-S)\delta\tau} \approx 1 - (\hat{H} - S)\delta\tau$
- Exact ground state still emerges if  $\delta\tau < 2/(E_{max} - E_{FCI})$ .

J. S. Spencer *et. al.*, J. Chem. Phys. **136**, 054110 (2012)  
(This is just a stochastic powers method)

# Full Configuration Interaction Quantum Monte Carlo: Sampling

- We use integer weights of determinants in the vector, although this is not the only choice.

F. R. Petruzielo *et. al.*, Phys. Rev. Lett. **109**, 230201 (2012)

- Positive (negative) coefficients are represented by a number of positive (negative) psips ( $\Psi$  particles).
- For this configuration:

$$|\Psi(\tau)\rangle = +0.25 |D_0\rangle + 0.5 |D_1\rangle - 0.25 |D_2\rangle$$

J . B. Anderson J. Chem. Phys. **63**, 1499 (1975)

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- Positive (negative) coefficients are represented by a number of positive (negative) psips ( $\Psi$  particles).
- For this configuration:

$$|\Psi(\tau)\rangle = \text{green circle} |D_0\rangle + \text{green circle} |D_1\rangle + \text{red circle} |D_2\rangle$$

J . B. Anderson J. Chem. Phys. **63**, 1499 (1975)

- Move  $\tau$  forwards by  $\delta\tau$ :

$$|\psi(\tau + \delta\tau)\rangle = \left(1 - (\hat{H} - S)\delta\tau\right) |\psi(\tau)\rangle$$

$$|\psi(\tau + \delta\tau)\rangle = \left(1 - \sum_{i \neq j} (\langle D_i | H | D_j \rangle) \delta\tau - \sum_i (\langle D_i | H | D_i \rangle - S) \delta\tau\right) |\psi(\tau)\rangle$$



# Full Configuration Interaction Quantum Monte Carlo

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## Spawning

$$|\Psi(\tau)\rangle = \begin{array}{c} \circ \\ \circ \end{array} |D_0\rangle + \cdots \begin{array}{c} \circ \\ \circ \\ \circ \end{array} |D_i\rangle + \cdots \circ |D_j\rangle + \cdots$$

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# Full Configuration Interaction Quantum Monte Carlo

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## Spawning

$$|\Psi(\tau)\rangle = \begin{array}{c} \text{green} \\ \text{green} \end{array} |D_0\rangle + \dots + \begin{array}{c} \text{green} \\ \text{green} \\ \text{green} \end{array} |D_i\rangle + \dots + \begin{array}{c} \text{green} \\ \text{green} \\ \text{red} \end{array} |D_j\rangle + \dots$$

- with probability  $\langle D_i | H | D_j \rangle \delta\tau$
- if  $\langle D_i | H | D_j \rangle \delta\tau < 0$  psp has same sign as parent and vice-versa.

# Full Configuration Interaction Quantum Monte Carlo

- Move  $\tau$  forwards by  $\delta\tau$ :

$$|\psi(\tau+\delta\tau)\rangle = \left( 1 - \sum_{i \neq j} (\langle D_i | H | D_j \rangle) \delta\tau - \sum_i (\langle D_i | H | D_i \rangle - S) \delta\tau \right) |\psi(\tau)\rangle$$

## Diagonal Death

$$|\Psi(\tau)\rangle = \begin{array}{c} \circ \\ \circ \end{array} |D_0\rangle + \cdots + \begin{array}{c} \circ \\ \circ \\ \circ \end{array} |D_i\rangle + \cdots + \begin{array}{c} \circ \end{array} |D_j\rangle + \cdots$$

- Move  $\tau$  forwards by  $\delta\tau$ :

$$|\psi(\tau+\delta\tau)\rangle = \left( \mathbf{1} - \sum_{i \neq j} (\langle D_i | H | D_j \rangle) \delta\tau - \sum_i (\langle D_i | H | D_i \rangle - S) \delta\tau \right) |\psi(\tau)\rangle$$

## Diagonal Death

$$|\Psi(\tau)\rangle = \begin{array}{c} \circ \\ \circ \end{array} |D_0\rangle + \dots \begin{array}{c} \circ \\ \circ \end{array} |D_i\rangle + \dots \begin{array}{c} \circ \\ \circ \end{array} |D_j\rangle + \dots$$

- Death occurs with probability  $(\langle D_i | H | D_i \rangle - S) \delta\tau$
- Cloning occurs (population becomes more negative or positive) if  $(\langle D_i | H | D_i \rangle - S) > 0$

# Full Configuration Interaction Quantum Monte Carlo

- Move  $\tau$  forwards by  $\delta\tau$ :

$$|\psi(\tau+\delta\tau)\rangle = \left( 1 - \sum_{i \neq j} (\langle D_i | H | D_j \rangle) \delta\tau - \sum_i (\langle D_i | H | D_i \rangle - S) \delta\tau \right) |\psi(\tau)\rangle$$

## Annihilation

$$|\Psi(\tau)\rangle = \begin{array}{c} \circ \\ \circ \end{array} |D_0\rangle + \dots + \begin{array}{c} \circ \\ \circ \\ \circ \end{array} |D_i\rangle + \dots + \begin{array}{c} \circ \\ \circ \\ \circ \end{array} |D_j\rangle + \dots$$

- Move  $\tau$  forwards by  $\delta\tau$ :

$$|\psi(\tau+\delta\tau)\rangle = \left( 1 - \sum_{i \neq j} (\langle D_i | H | D_j \rangle) \delta\tau - \sum_i (\langle D_i | H | D_i \rangle - S) \delta\tau \right) |\psi(\tau)\rangle$$

## Annihilation

- psips with opposite signs on the same determinant annihilate.

$$|\Psi(\tau)\rangle = \begin{array}{c} \circ \\ \circ \end{array} |D_0\rangle + \dots \begin{array}{c} \circ \\ \circ \\ \circ \end{array} |D_i\rangle + \dots |D_j\rangle + \dots$$

## Projected Energy

- Project onto a reference state:

$$E = \frac{\langle D_0 | \hat{H} e^{-(\hat{H}-S)\tau} | D_0 \rangle}{\langle D_0 | e^{-(\hat{H}-S)\tau} | D_0 \rangle} = \frac{\langle D_0 | \hat{H} | \Psi_0 \rangle}{\langle D_0 | \Psi_0 \rangle}$$

- Average over every step after the simulation has equilibrated.

## Shift

- $S$  is initially set to a constant until the total number of psips  $N$  reaches the desired level.
- After which  $S$  is updated every  $A$  steps according to:

$$S(\tau + A\delta\tau) = S(\tau) - \frac{\gamma}{A\delta\tau} \log \frac{N(\tau + A\delta\tau)}{N(\tau)}$$



## Part II

# FCIQMC and Markov Chain Monte Carlo

# Markov Chain Monte Carlo

- A **Markov** chain is a stochastic system in which each step **only** depends on the previous step.
- This means the system can be described by a matrix  $\Gamma_{\alpha\beta}$  stochastic matrix: The probability that the system transitions from state  $\alpha$  to state  $\beta$  in one '**Markov**' step.
- Under some conditions  $\Gamma$  has an eigenvector  $\gamma$  with eigenvalue one.
- $\gamma_\alpha$  probability the system is in a state  $\alpha$  after equilibration.
- FCIQMC seems to be described by an infinite number of states as the shift can take on any value.

## An Example:

$$\alpha = (S, \text{green circle} |D_0\rangle + \text{green circle} |D_1\rangle + \text{red circle} |D_2\rangle)$$

# The Shift Again

- The shift is only function of the number of psips  $N$  on the last shift update step and the number on the step before the shift turned on  $N_s$  :

$$S(\tau + A\delta\tau) = S(\tau) - \frac{\gamma}{A\delta\tau} \log \frac{N(\tau + A\delta\tau)}{N(\tau)}$$

$$S(\tau + A\delta\tau) = S(\tau - A\delta\tau) - \frac{\gamma}{A\delta\tau} \log \frac{N(\tau)}{N(\tau - A\delta\tau)} - \frac{\gamma}{A\delta\tau} \log \frac{N(\tau + A\delta\tau)}{N(\tau)}$$

$$S(\tau) = S_0 - \frac{\gamma}{A\delta\tau} \log \frac{N(\tau)}{N_s}$$

- Thus we make one 'Markov' step every  $A$  steps of  $\delta\tau$ .
- $A = 1$  for mathematical simplicity

# The FCIQMC Markov Chain: for two Determinants

We want  $\Gamma_{\alpha,\beta}$ . The population of psips on each determinant is changed by:

- 1 Death of psips on the determinant.
- 2 Spawning from the other determinant onto this determinant.

These events are Binomially distributed probability:

$$B(n, N, p) = \binom{N}{n} p^n (1-p)^{N-n} \quad (1)$$

Sum over these binomial coefficients such that the change could have happened.

# The FCIQMC Markov Chain: for two Determinants

## An Example

$$\alpha = \begin{array}{c} \circ \\ \circ \\ \circ \\ \circ \end{array} |D_a\rangle + \begin{array}{c} \circ \\ \circ \end{array} |D_b\rangle$$

$$\beta = \begin{array}{c} \circ \\ \circ \\ \circ \end{array} |D_a\rangle + \begin{array}{c} \circ \\ \circ \end{array} |D_b\rangle$$

- If  $\langle D_b|H|D_a\rangle < 0$  and  $\langle D_a|H|D_a\rangle - S > 0$  then the probability that the population changes on **a**:

$$B(1, 4, P_{da})B(0, 2, P_{sa}) + B(2, 4, P_{da})B(1, 2, P_{sa}) \\ + B(3, 4, P_{da})B(2, 2, P_{sa})$$

- After which we just do the same for **b** and product the two probabilities to compute  $\Gamma_{\alpha\beta}$ .

# Conclusions: Markov Chain Monte Carlo and FCIQMC

- FCIQMC is a Markov chain.
- $\xi$ ,  $\delta\tau$ ,  $N_s$ ,  $A$  and the system specify the chain.
- The size of the stochastic matrix scales terribly with system size.
- We can diagonalise the Stochastic Matrix to compute the stationary distribution.

## Part III

# Population Control Bias in FCIQMC

# Population control

- Let's scale  $\gamma$  when we change  $\delta\tau$  and  $A$ , to fix  $\xi$ :

$$\xi = \frac{\gamma}{A\delta\tau} \quad (2)$$

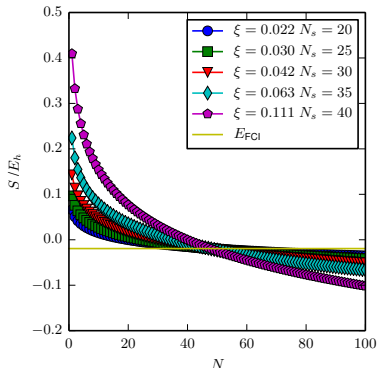
- If we assume that  $S = E_{FCI}$  at  $\langle N \rangle$

$$E_{FCI} = -\xi \log \frac{\langle N \rangle}{N_0}$$

$$\langle N \rangle = N_0 e^{-\frac{\xi}{E_{FCI}}}$$

- Thus we can fix  $\langle N \rangle$  if we know the correlation energy.

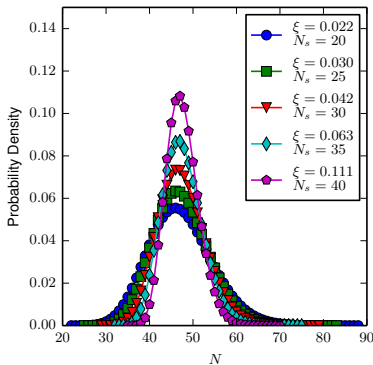
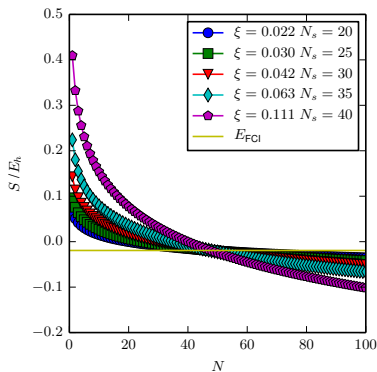
## H2 in a STO-3g Basis set





# Distribution of the Number of psips

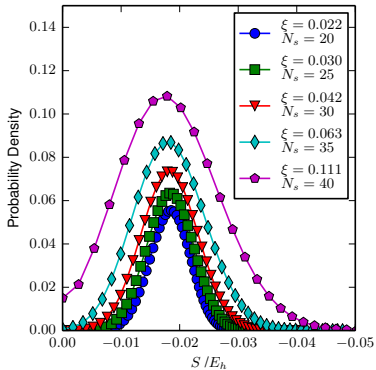
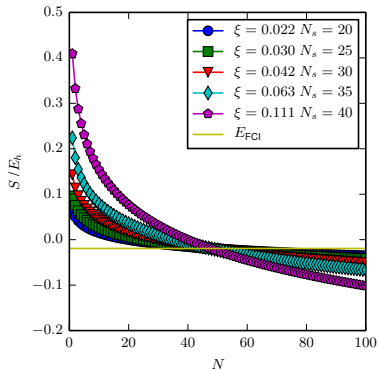
H2 in a STO-3g Basis set  $r_0 = 0.7122 \text{ \AA}$



- If  $\xi$  is big then the population is well controlled.

# Distribution of the Shift

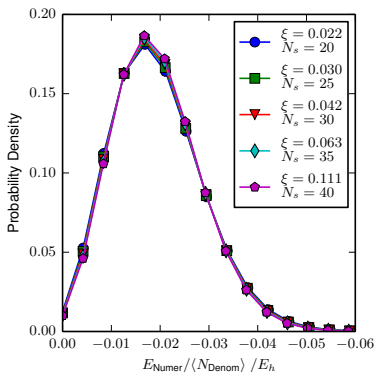
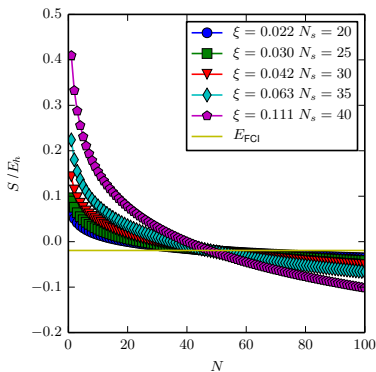
H2 in a STO-3g Basis set  $r_0 = 0.7122 \text{ \AA}$



- If  $\xi$  is big then shift has a large variance.

# Distribution of the Projected Energy

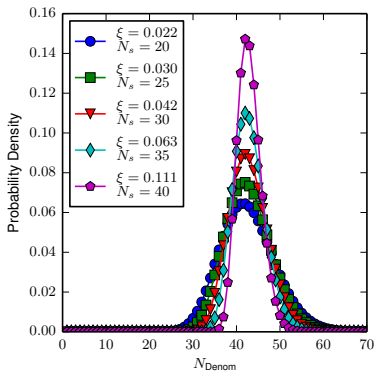
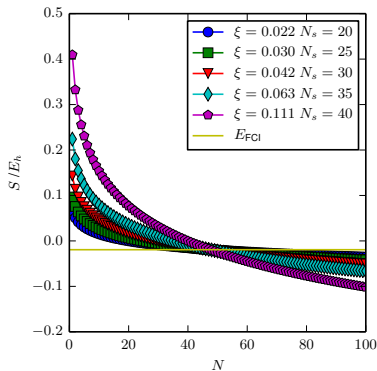
H2 in a STO-3g Basis set  $r_0 = 0.7122 \text{ \AA}$



- $\xi$  has little effect on the numerator of the projected energy.

# Distribution of the Projected Energy

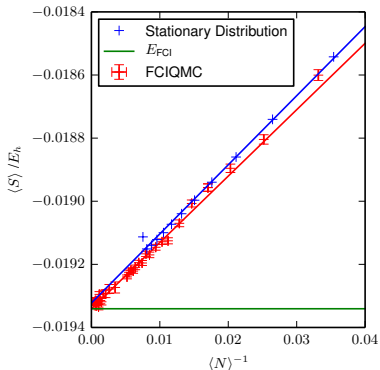
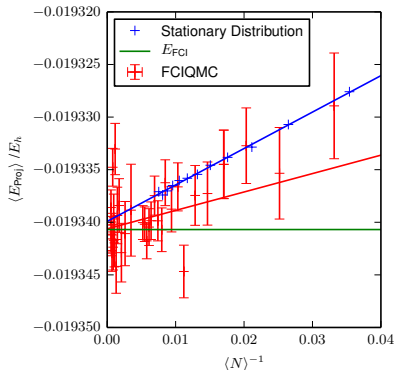
H2 in a STO-3g Basis set  $r_0 = 0.7122 \text{ \AA}$



- If  $\xi$  effects the denominator of the projected energy in the same way as the population.

# Population Control Bias in $H_2$

$H_2$  in a STO-3g Basis set  $r_0 = 0.7122 \text{ \AA}$



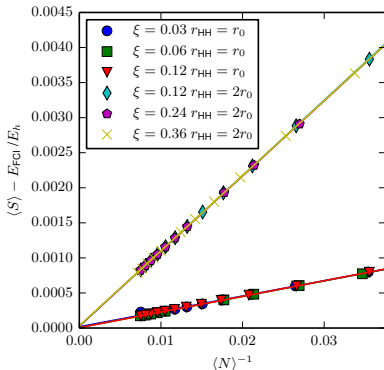
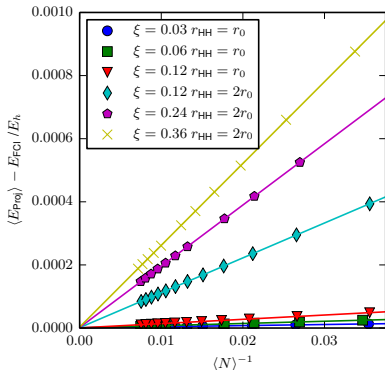
- The mean of both estimators of the energy is incorrect.

$$\mathcal{O}(1/\langle N \rangle)$$

N. Cerf *et. al.*, Phys. Rev. E **51**, 3679 (1995)

# Population Control Bias in $H_2$

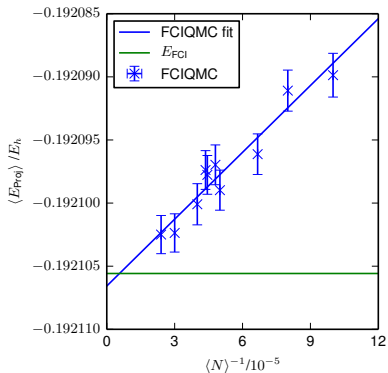
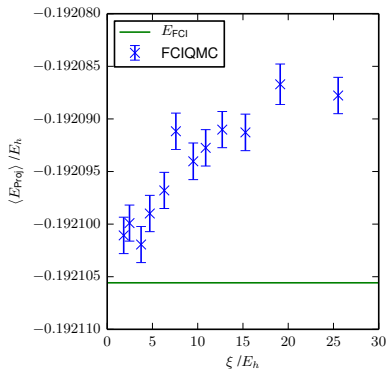
$H_2$  in a STO-3g Basis set  $r_0 = 0.7122 \text{ \AA}$



- If  $\xi$  is small then population control bias decreases.
- Strongly correlated systems seem to be more affected.

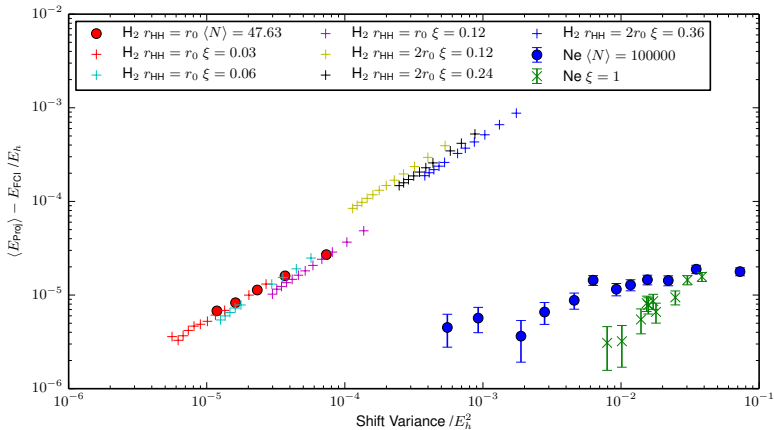
# Population Control Bias in Ne cc-pVDZ

## Ne cc-pVDZ



- We see similar thing for Ne cc-pVDZ.

## Shift Variance





# Minimise Population Control Bias

- Use as large as population as possible.
- Reduce  $\xi$  (don't let the population fall below the plateau).
- If you really care about accuracy extrapolate  $1/\langle N \rangle \rightarrow \infty$  with a fixed  $\xi$ .

# Acknowledgements

- **Thank you for listening**
- Alex Thom, Michael Bearpark, James Spencer
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- James Spencer, Alex Thom, Nick Blunt, Fionn Malone, James Shepherd, Matthew Foulkes, Thomas Rogers, Will Handley, Joseph Weston
- Imperial College High Performance Computing Service

See [arXiv:1407.1753](https://arxiv.org/abs/1407.1753) for more details.