

# An overview of vibrations in solids

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**TCM**

QMC in the Apuan Alps IX  
27 July 2014

## Why vibrations in a QMC conference?

- ▶ If the required accuracy makes QMC necessary, then vibrations could also be important.
- ▶ QMC band gaps are static, the effects of electron-phonon coupling may be important.

# Outline

The vibrational energy in solids

- Theoretical background

- Applications

Vibrational coupling in solids

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- Applications

Conclusions

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## Why is the vibrational problem difficult?

$$\hat{H}_{\text{vib}} = -\frac{1}{2} \sum_{\mathbf{R}_p, \alpha} \frac{1}{m_\alpha} \nabla_{p\alpha}^2 + V(\mathbf{r}_\alpha)$$

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- ▶ 3*N*-dimensional function

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- ▶  $3N$ -dimensional function
- ▶ Each data point requires an electronic energy calculation

# Harmonic approximation

- ▶ Vibrational Hamiltonian in  $\{\mathbf{r}_\alpha\}$  (or  $\{\mathbf{u}_\alpha\}$ ):

$$\hat{H}_{\text{vib}}^{\text{har}} = -\frac{1}{2} \sum_{\mathbf{R}_p, \alpha} \frac{1}{m_\alpha} \nabla_{p\alpha}^2 + \frac{1}{2} \sum_{\mathbf{R}_p, \alpha; \mathbf{R}_{p'}, \beta} \mathbf{u}_{p\alpha} \Phi_{p\alpha; p'\beta} \mathbf{u}_{p'\beta}$$

- ▶ Normal mode analysis:  $\{\mathbf{u}_{p\alpha}\} \longrightarrow \{q_{\mathbf{k}s}\}$
- ▶ Vibrational Hamiltonian in  $\{q_{\mathbf{k}s}\}$ :

$$\hat{H}_{\text{vib}}^{\text{har}} = \sum_{\mathbf{k}, s} \left( -\frac{1}{2} \frac{\partial^2}{\partial q_{\mathbf{k}s}^2} + \frac{1}{2} \omega_{\mathbf{k}s}^2 q_{\mathbf{k}s}^2 \right)$$

## Principal axes approximation to the BO energy surface

$$V(\mathbf{q}) = V(0) + \sum_{\mathbf{k}, s} V_{\mathbf{k}s}(q_{\mathbf{k}s}) + \frac{1}{2} \sum_{\mathbf{k}, s} \sum'_{\mathbf{k}', s'} V_{\mathbf{k}s; \mathbf{k}'s'}(q_{\mathbf{k}s}, q_{\mathbf{k}'s'}) + \dots$$

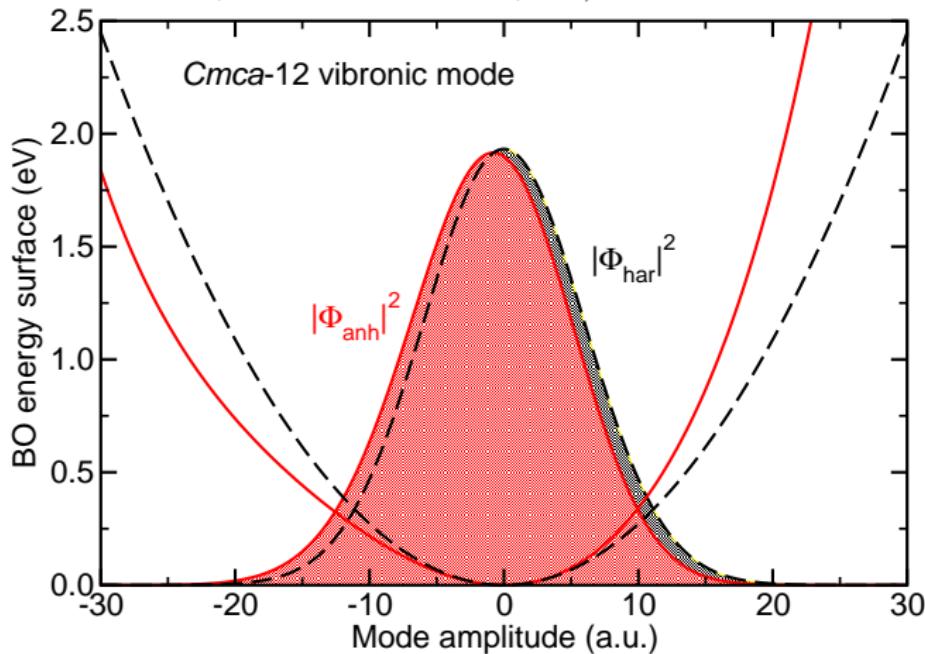
- ▶ **Static lattice** DFT total energy.
- ▶ DFT total energy along frozen **independent mode**.
- ▶ DFT total energy along frozen **coupled modes**.

Features:

- ▶ Can be improved systematically.
- ▶ Subspace with higher  $N$ -body terms (e.g. perovskites).
- ▶ Estimate of error in anharmonic energy.

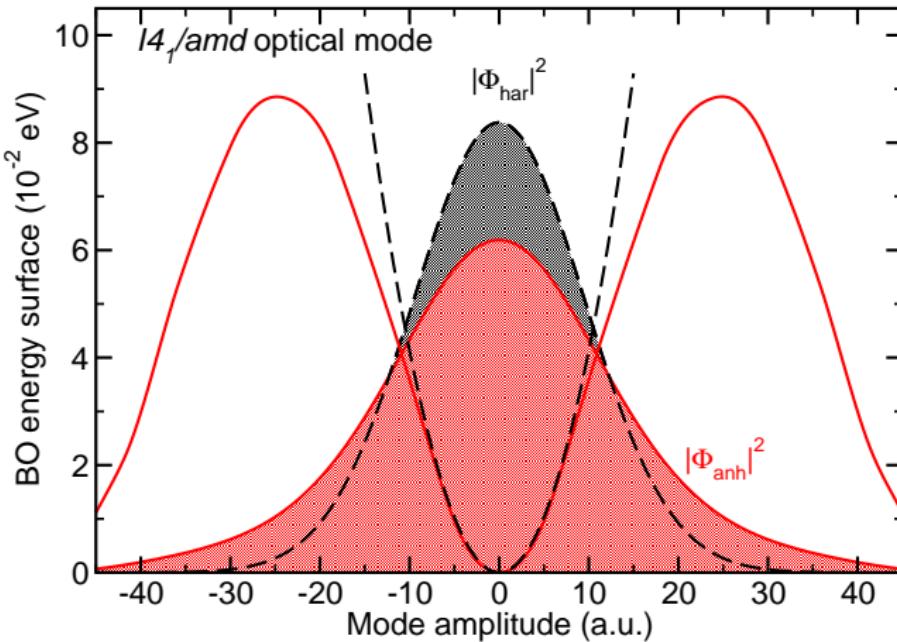
## Independent term (I)

$$V(\mathbf{q}) = V(0) + \sum_{\mathbf{k}, s} V_{\mathbf{k}s}(q_{\mathbf{k}s}) + \frac{1}{2} \sum_{\mathbf{k}, s} \sum'_{\mathbf{k}', s'} V_{\mathbf{k}s; \mathbf{k}'s'}(q_{\mathbf{k}s}, q_{\mathbf{k}'s'}) + \dots$$



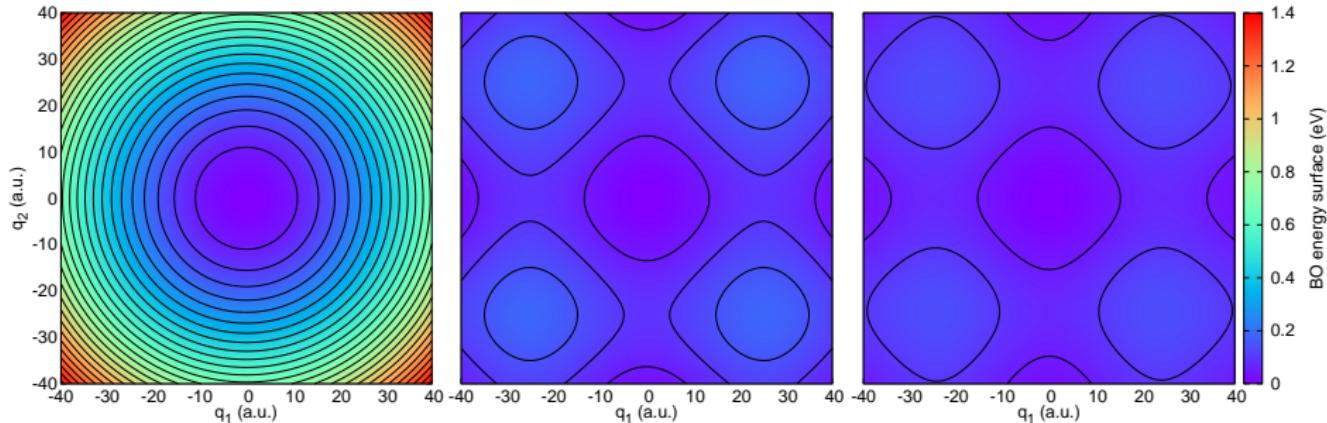
## Independent term (II)

$$V(\mathbf{q}) = V(0) + \sum_{\mathbf{k}, s} V_{\mathbf{k}s}(q_{\mathbf{k}s}) + \frac{1}{2} \sum_{\mathbf{k}, s} \sum'_{\mathbf{k}', s'} V_{\mathbf{k}s; \mathbf{k}'s'}(q_{\mathbf{k}s}, q_{\mathbf{k}'s'}) + \dots$$



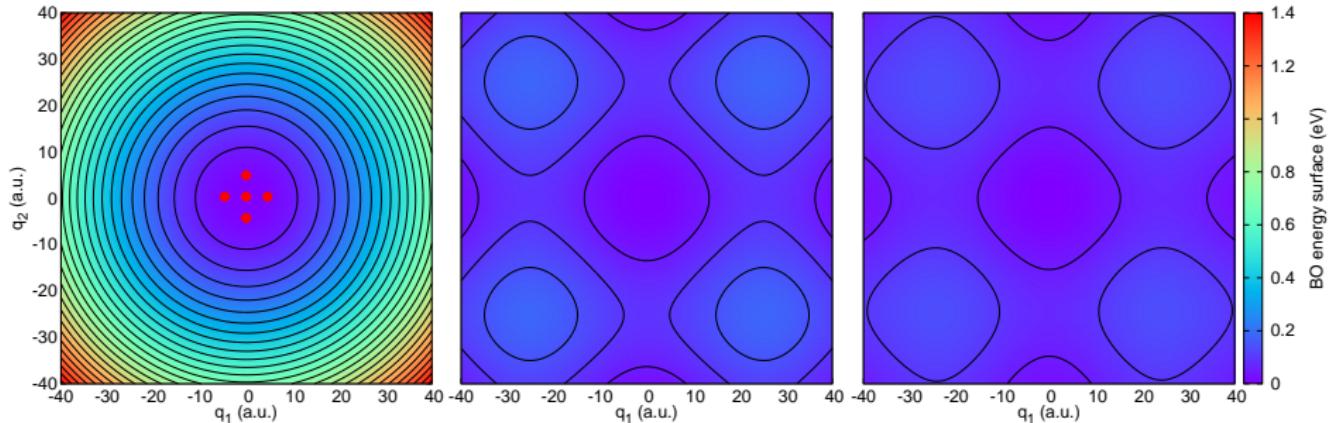
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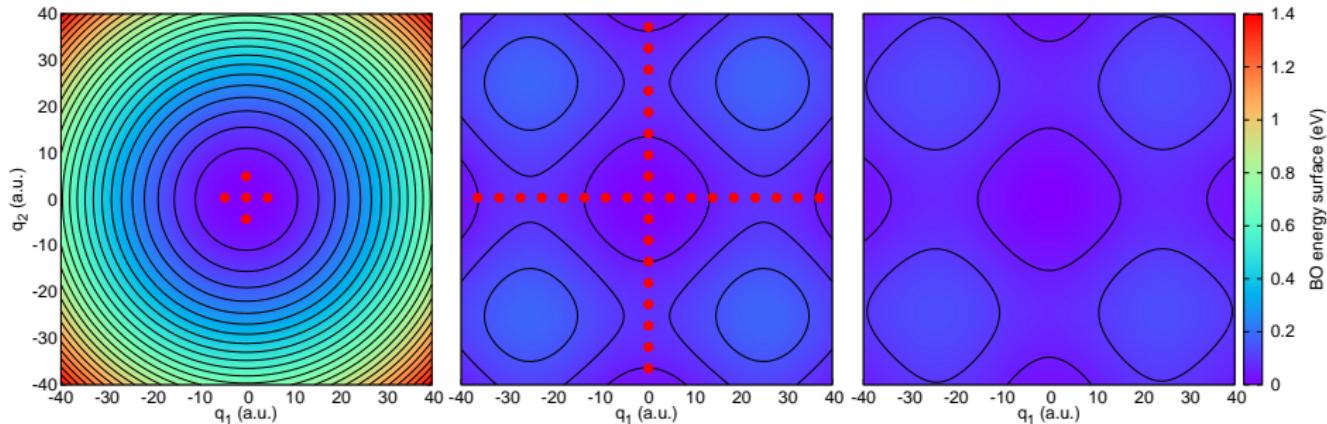
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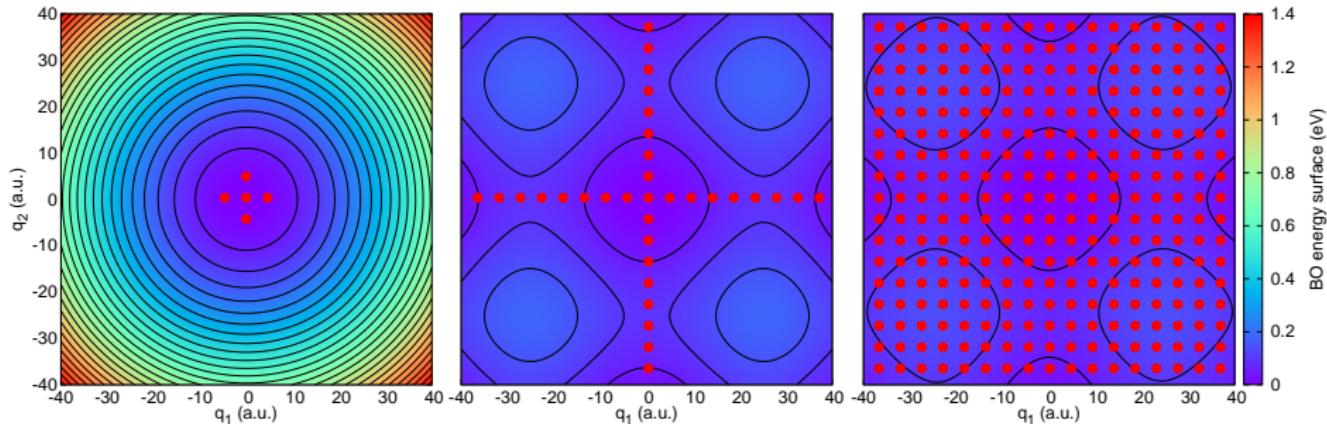
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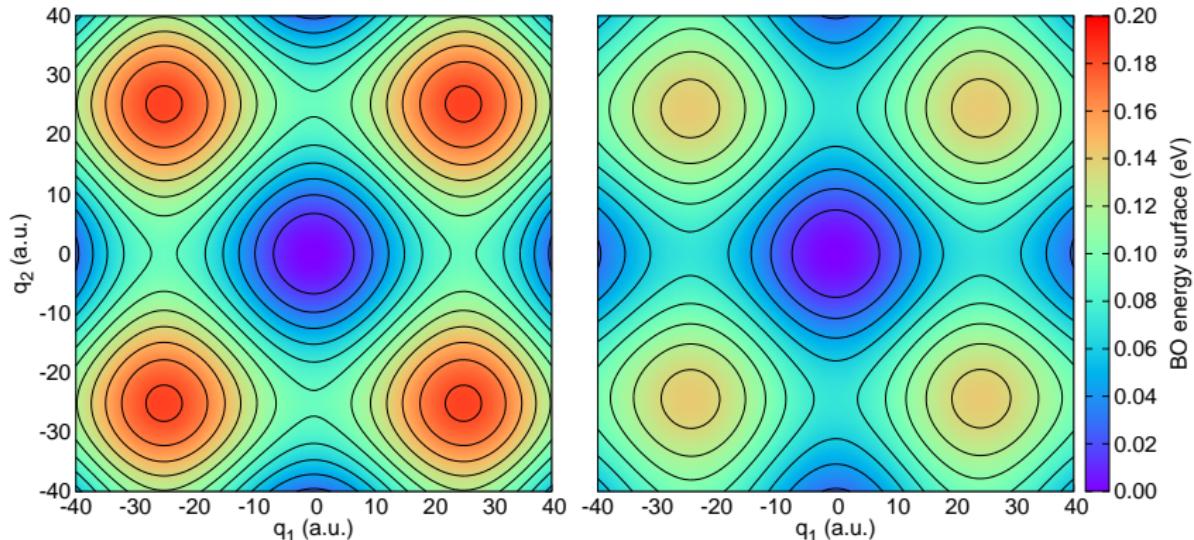
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# Vibrational self-consistent field equations

- ▶ Vibrational Schrödinger equation:

$$\left( \sum_{\mathbf{k},s} -\frac{1}{2} \frac{\partial^2}{\partial q_{\mathbf{ks}}^2} + V(\mathbf{q}) \right) \Phi(\mathbf{q}) = E \Phi(\mathbf{q})$$

- ▶ **Ansatz:**  $\Phi(\mathbf{q}) = \prod_{\mathbf{k},s} \phi_{\mathbf{ks}}(q_{\mathbf{ks}})$
- ▶ Self-consistent equations:

$$\left( -\frac{1}{2} \frac{\partial^2}{\partial q_{\mathbf{ks}}^2} + \bar{V}_{\mathbf{ks}}(q_{\mathbf{ks}}) \right) \phi_{\mathbf{ks}}(q_{\mathbf{ks}}) = \lambda_{\mathbf{ks}} \phi_{\mathbf{ks}}(q_{\mathbf{ks}})$$

$$\bar{V}_{\mathbf{ks}}(q_{\mathbf{ks}}) = \left\langle \prod'_{\mathbf{k}',s'} \phi_{\mathbf{k}'s'}(q_{\mathbf{k}'s'}) \middle| V(\{q_{\mathbf{k}''s''}\}) \middle| \prod'_{\mathbf{k}',s'} \phi_{\mathbf{k}'s'}(q_{\mathbf{k}'s'}) \right\rangle$$

## Second order perturbation theory

- ▶ Second order perturbation theory (similar to MP2).
- ▶ Measures the accuracy of the mean-field approach.
- ▶ So far small MP2 corrections.
- ▶ Can use other methods: whole electronic structure hierarchy.

# Anharmonic free energy

- ▶ Anharmonic vibrational excited states:

$$|\Phi^S(\mathbf{q})\rangle = \prod_{\mathbf{k},s} |\phi_{\mathbf{k}s}^{S_{\mathbf{k}s}}(q_{\mathbf{k}s})\rangle$$

where  $\mathbf{S}$  is a vector with elements  $S_{\mathbf{k}s}$ .

- ▶ Anharmonic free energy:

$$\mathcal{F}_{\text{anh}} = -\frac{1}{\beta} \ln \sum_{\mathbf{S}} e^{-\beta E_{\mathbf{S}}}$$

Single calculation for insulators.

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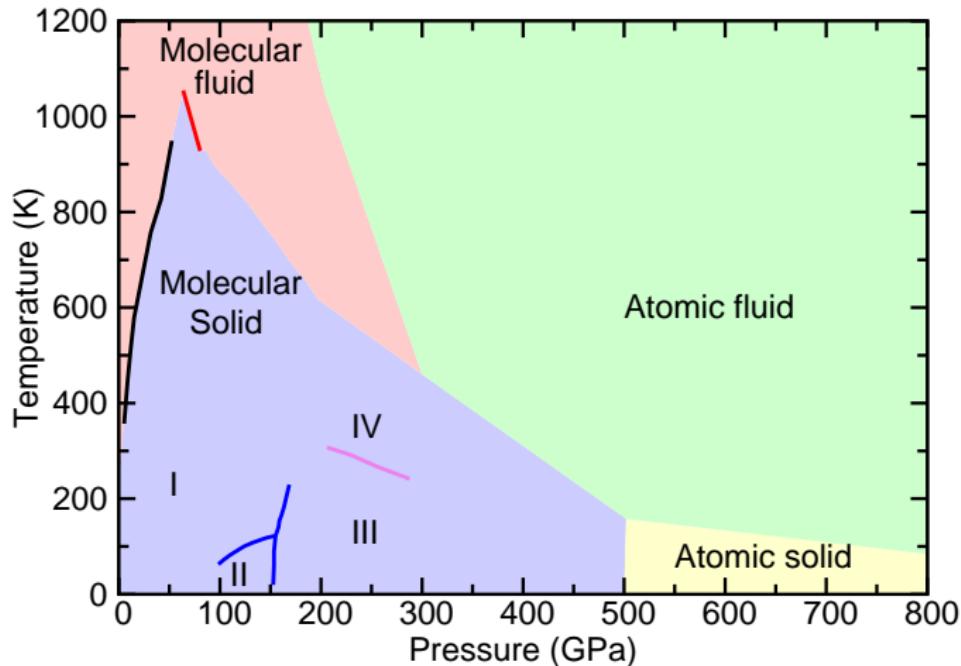
## Conclusions

# Solid hydrogen

- ▶ Most abundant element in the Universe.
- ▶ Hydrogen at high pressure in planetary interiors and stars.
- ▶ Possibility of exotic phases:  
high-temperature superconductivity,  
zero-temperature quantum fluid, ...

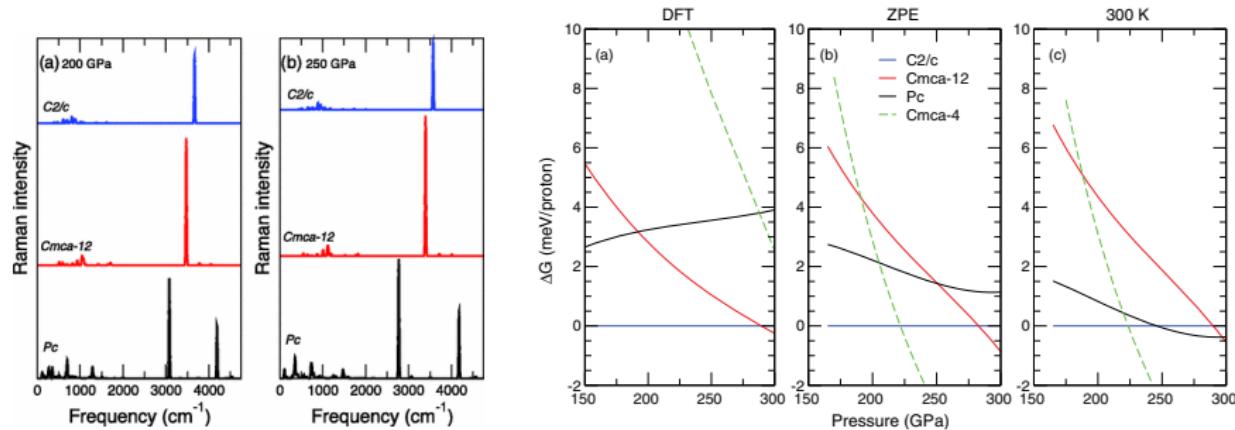


# The phase diagram of high pressure hydrogen



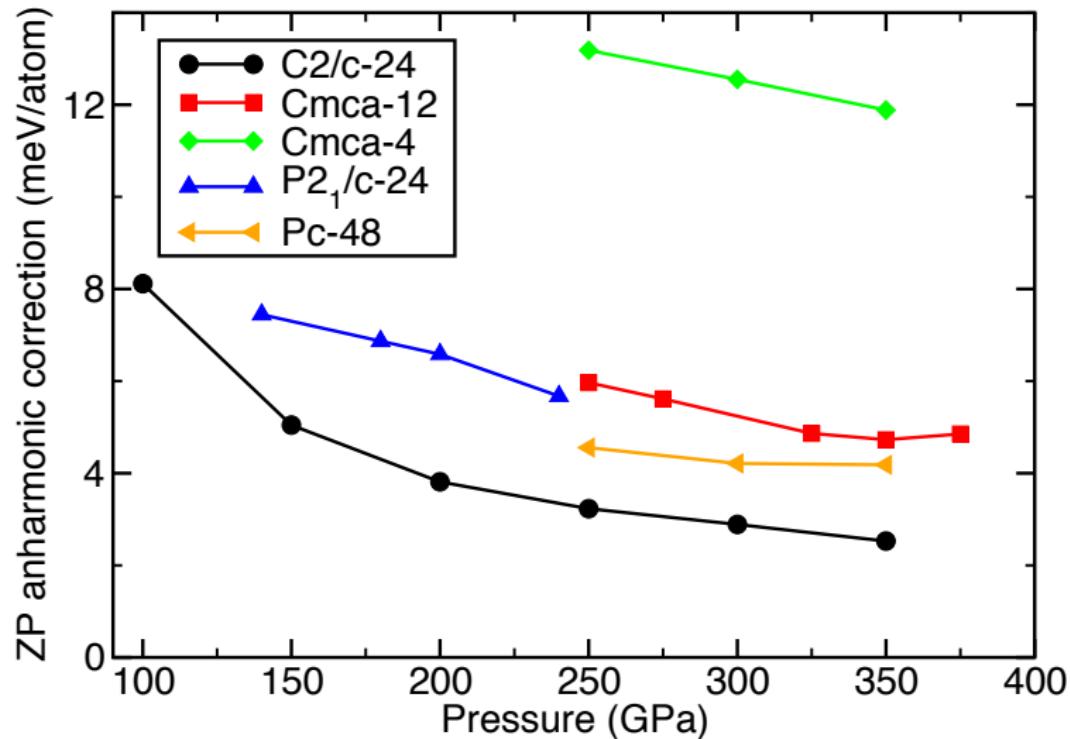
- Goncharov et al., Phys. Rev. Lett. **80**, 101 (1998)
- Datchi et al., Phys. Rev. B **61**, 6535 (2000)
- Gregoryanz et al., Phys. Rev. Lett. **90**, 175701 (2003)
- Deemyad and Silvera, Phys. Rev. Lett. **100**, 155701 (2008)
- Howie et al., Phys. Rev. Lett. **108**, 125501 (2012)

# The determination of phase IV

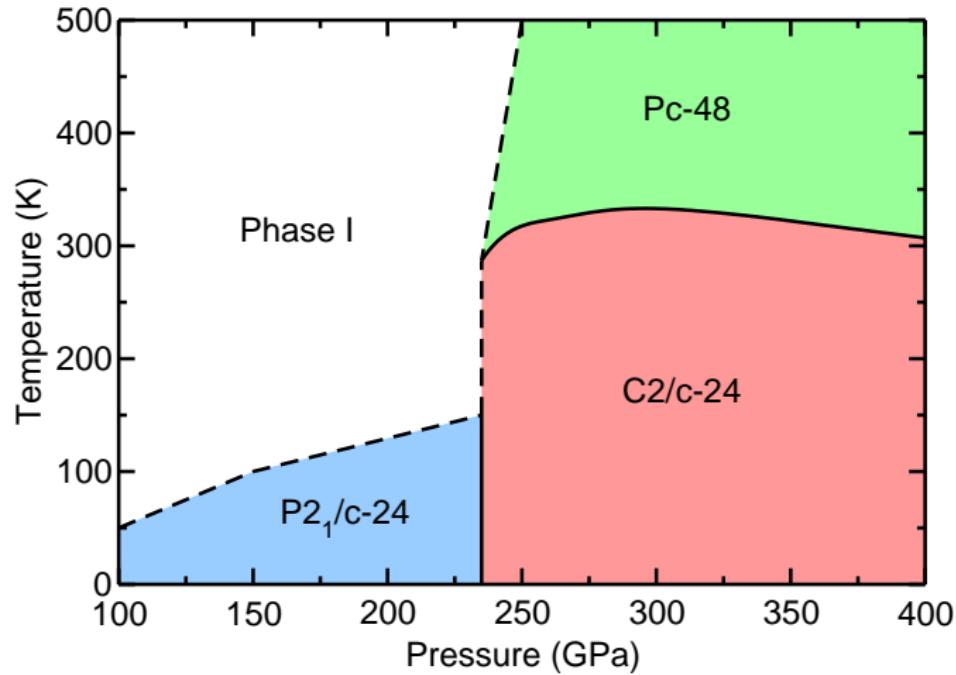


- Pickard, Martinez-Canales, and Needs, Phys. Rev. B **85**, 214114 (2012)
- Pickard, Martinez-Canales, and Needs, Phys. Rev. B **86**, 059902 (2012)

# Anharmonic vibrations in solid molecular hydrogen



# The phase diagram of solid molecular hydrogen

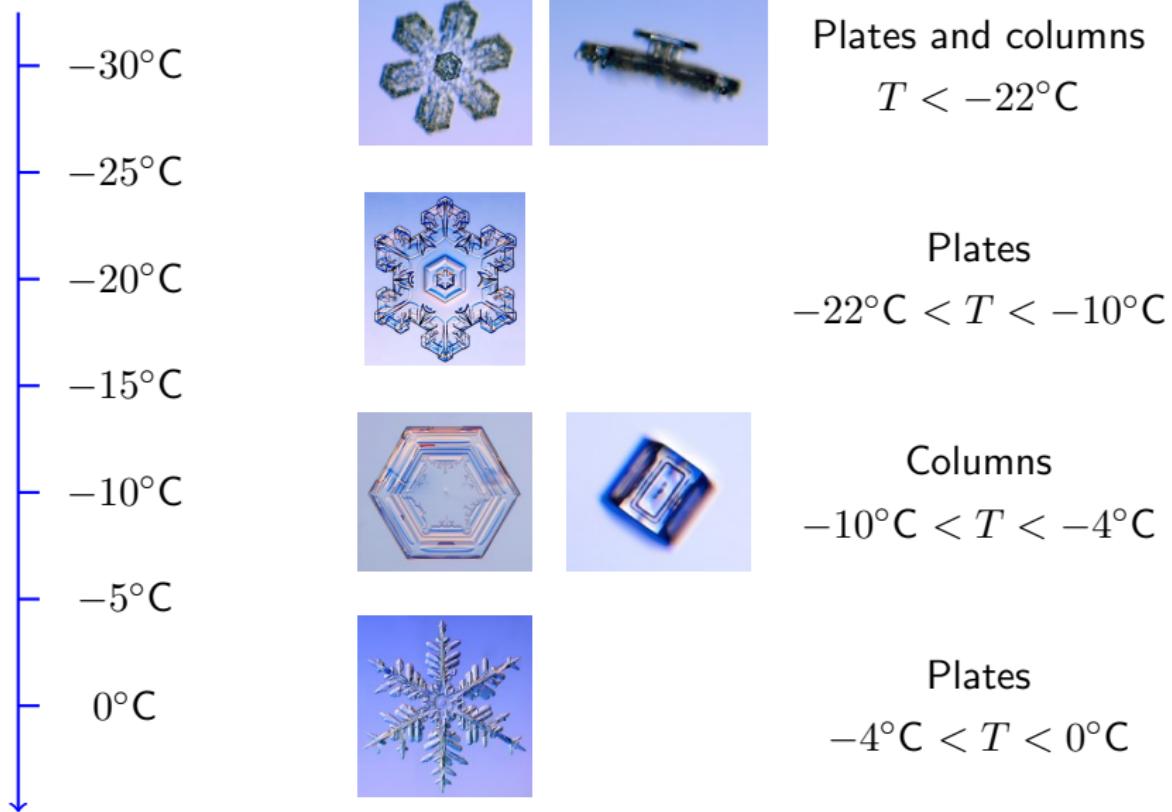


Lloyd-Williams, Monserrat, López Ríos, Drummond, and Needs

See talk on Tuesday at 12:00am:

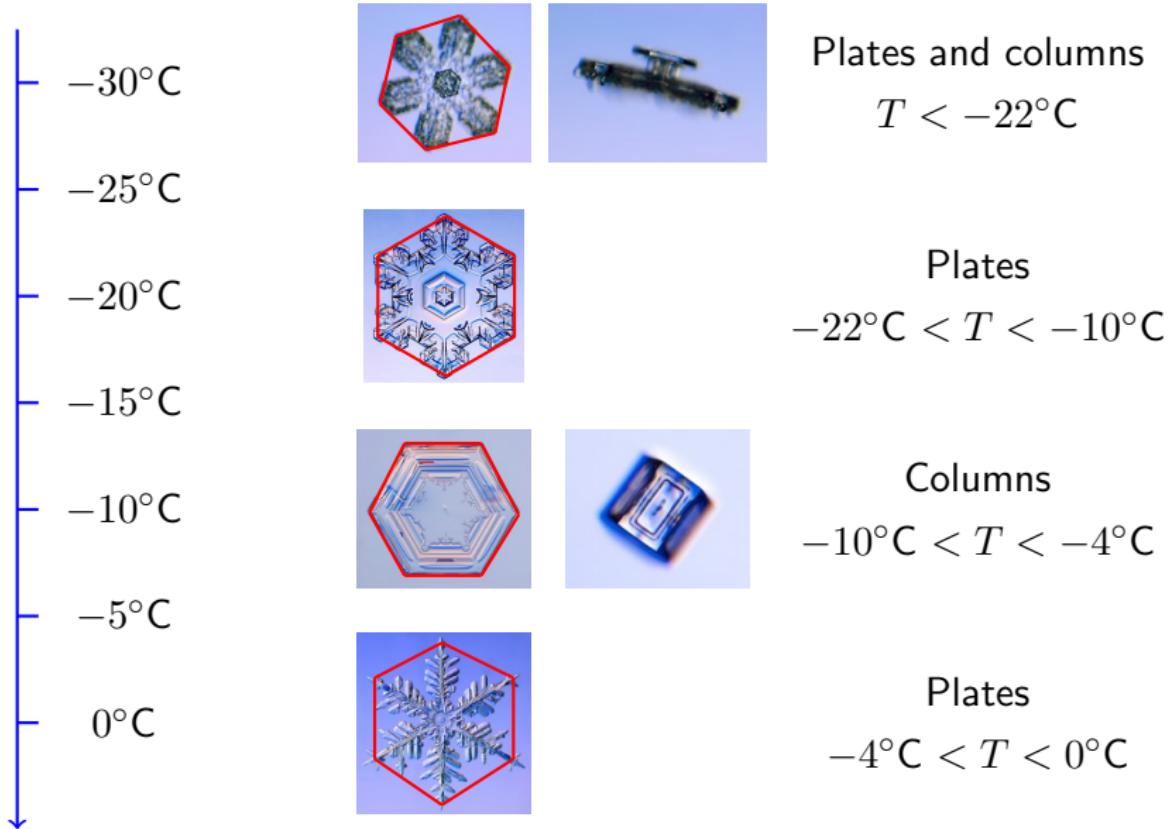
*DFT and QMC calculations of solid molecular hydrogen*

# Why are snowflakes hexagonal?



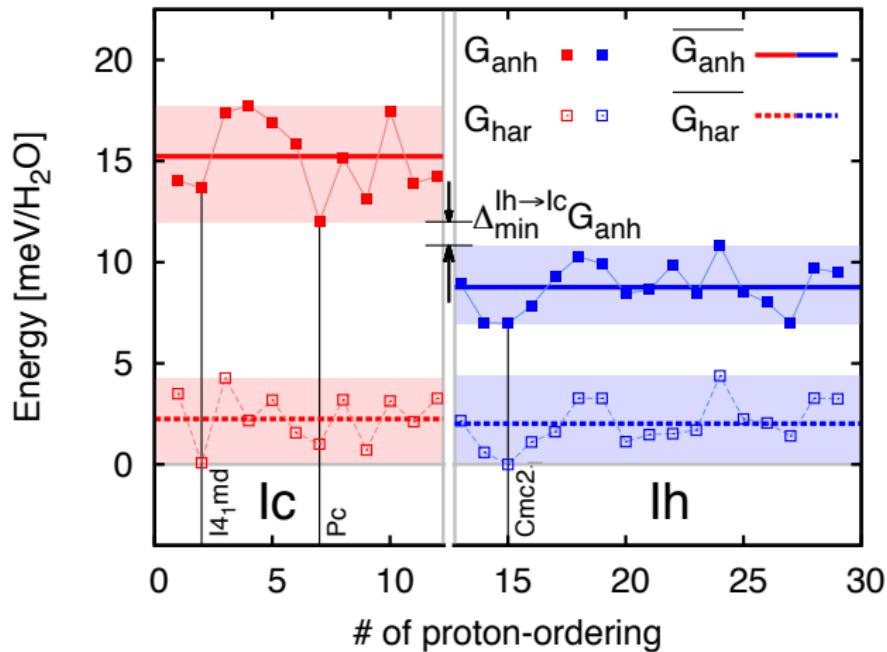
Pictures by Kenneth G. Libbrecht ([SnowCrystals.com](http://SnowCrystals.com))

# Why are snowflakes hexagonal?



Pictures by Kenneth G. Libbrecht (SnowCrystals.com)

# The relative stability of hexagonal and cubic ice



Engel, Monserrat, and Needs

See talk on Wednesday at 12:00am:

*Anharmonic nuclear motion and the relative stability of hexagonal and cubic ice*

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## Vibrational coupling to generic physical observables

$$\langle \mathcal{A} \rangle = \langle \Phi(\mathbf{q}) | \mathcal{A}(\mathbf{q}) | \Phi(\mathbf{q}) \rangle$$

# Vibrational coupling to generic physical observables

$$\langle \mathcal{A} \rangle = \langle \Phi(\mathbf{q}) | \mathcal{A}(\mathbf{q}) | \Phi(\mathbf{q}) \rangle$$

- Vibrational wave function: harmonic or anharmonic

# Vibrational coupling to generic physical observables

$$\langle \mathcal{A} \rangle = \langle \Phi(\mathbf{q}) | \mathcal{A}(\mathbf{q}) | \Phi(\mathbf{q}) \rangle$$

- ▶ Vibrational wave function: harmonic or anharmonic
- ▶ Observable coupling:
  - ▶ Expansion:

$$\mathcal{A}(\mathbf{q}) = \sum_{n,\mathbf{k}} a_{n\mathbf{k}}^{(1)} q_{n\mathbf{k}} + \sum_{n,\mathbf{k}} \sum_{n',\mathbf{k}'} a_{n\mathbf{k};n'\mathbf{k}'}^{(2)} q_{n\mathbf{k}} q_{n'\mathbf{k}'} + \dots$$

$$\mathcal{A}(\mathbf{q}) = \sum_{n,\mathbf{k}} a_{n\mathbf{k}} q_{n\mathbf{k}}^2$$

- ▶ Monte Carlo sampling

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$$\mathcal{A}(\mathbf{q}) = \sum_{n,\mathbf{k}} a_{n\mathbf{k}} q_{n\mathbf{k}}^2$$

- ▶ Monte Carlo sampling
- ▶ Examples: electronic band gaps, chemical shielding tensor

# Vibrational phase space sampling

$$\langle \mathcal{A} \rangle = \langle \Phi(\mathbf{q}) | \mathcal{A}(\mathbf{q}) | \Phi(\mathbf{q}) \rangle$$

MD/PIMD

Random

Quadratic

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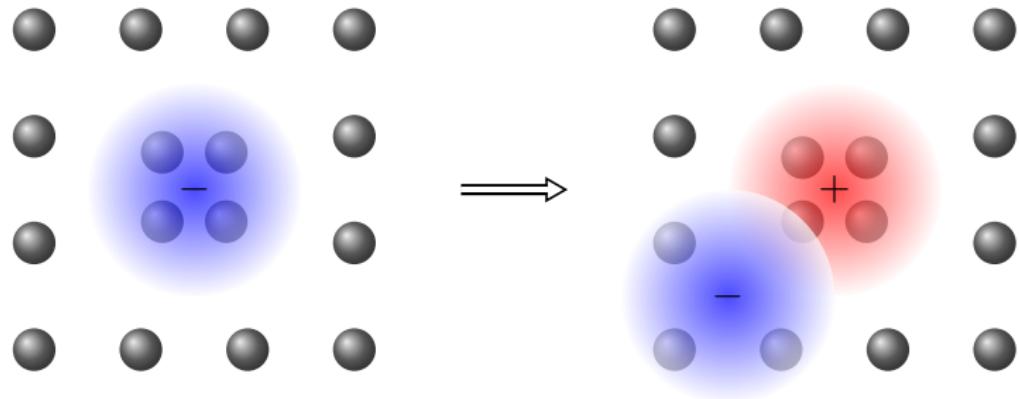
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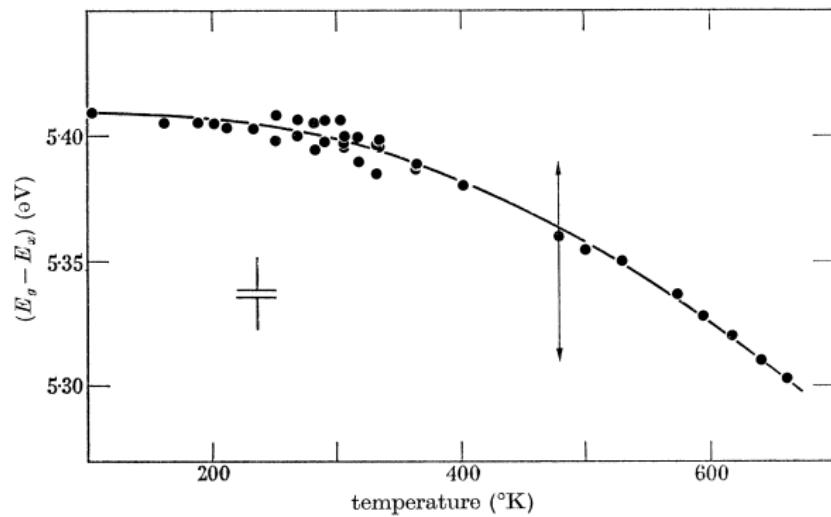
# Electron-phonon coupling in condensed matter

- ▶ Coupling to form Cooper pairs in standard superconductivity



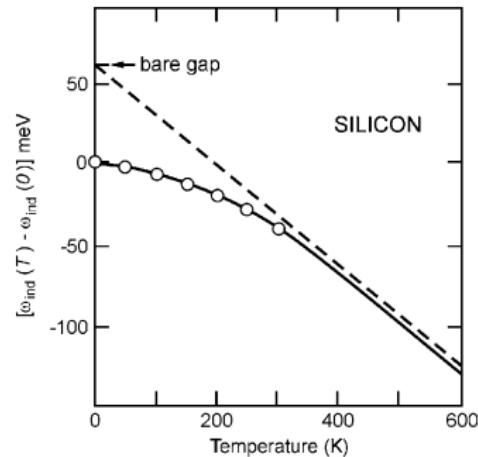
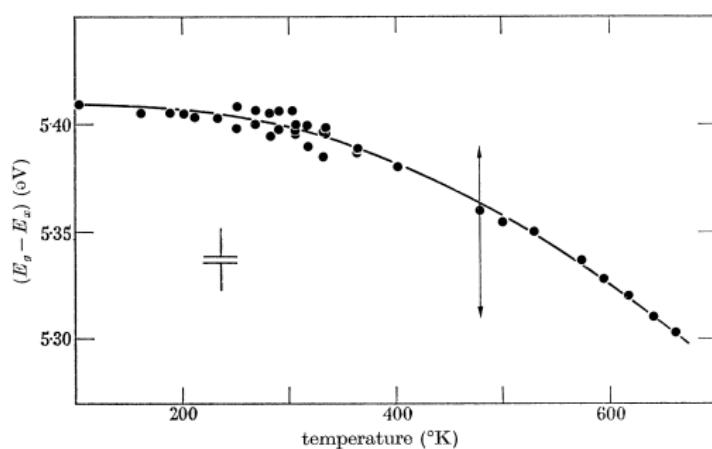
# Electron-phonon coupling in condensed matter

- ▶ Coupling to form Cooper pairs in standard superconductivity
- ▶ Temperature dependence of band gaps in semiconductors



→ Clark, Dean, and Harris, Proc. R. Soc. Lond. A **277**, 312 (1964)

# Electron-phonon coupling in diamond and silicon

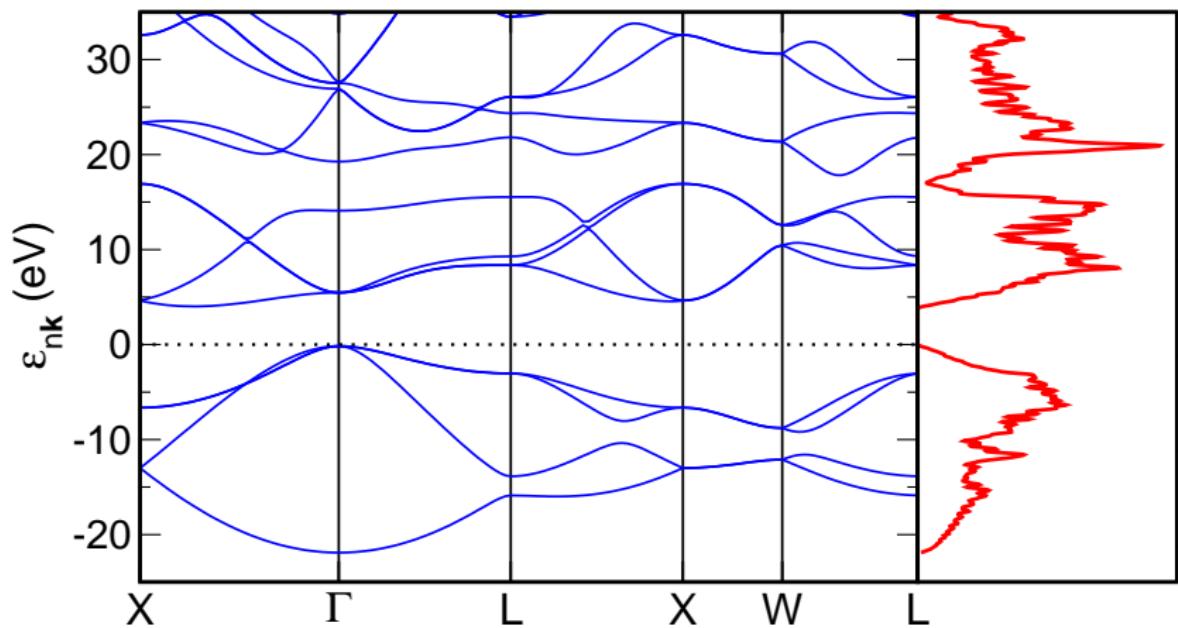


ZP band gap corrections:

- ▶ Silicon:  $-53$  meV
- ▶ Diamond: about  $-370$  meV

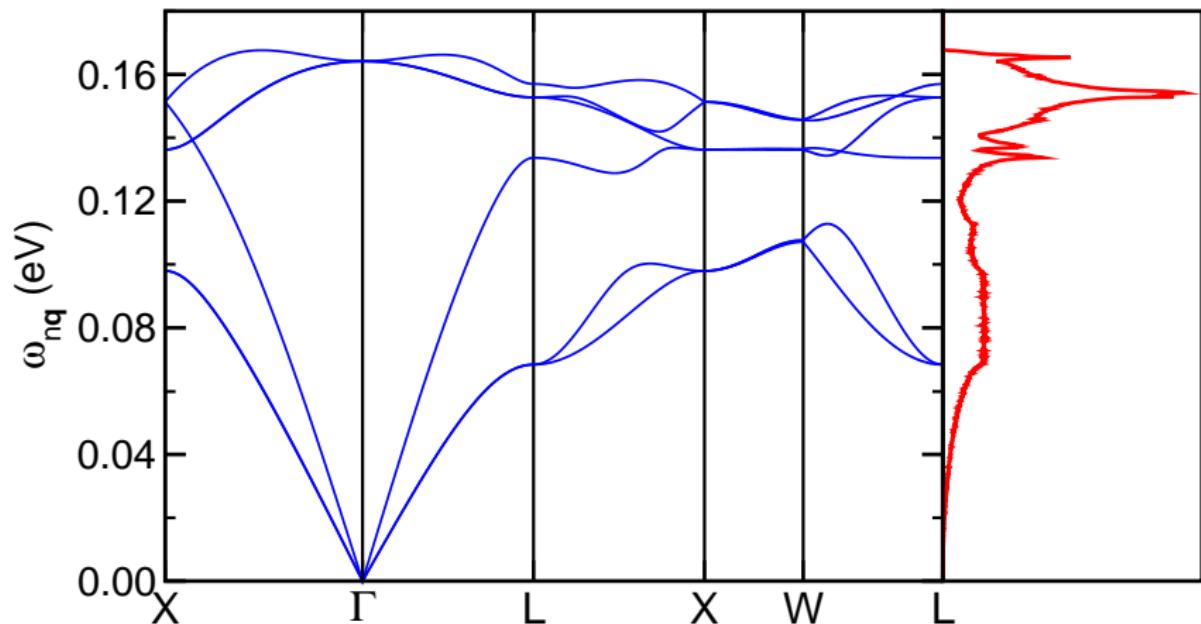
→ Clark, Dean, and Harris, Proc. R. Soc. Lond. A **277**, 312 (1964)  
→ Cardona, Solid State Comm. **133**, 3 (2005)

# Diamond band structure



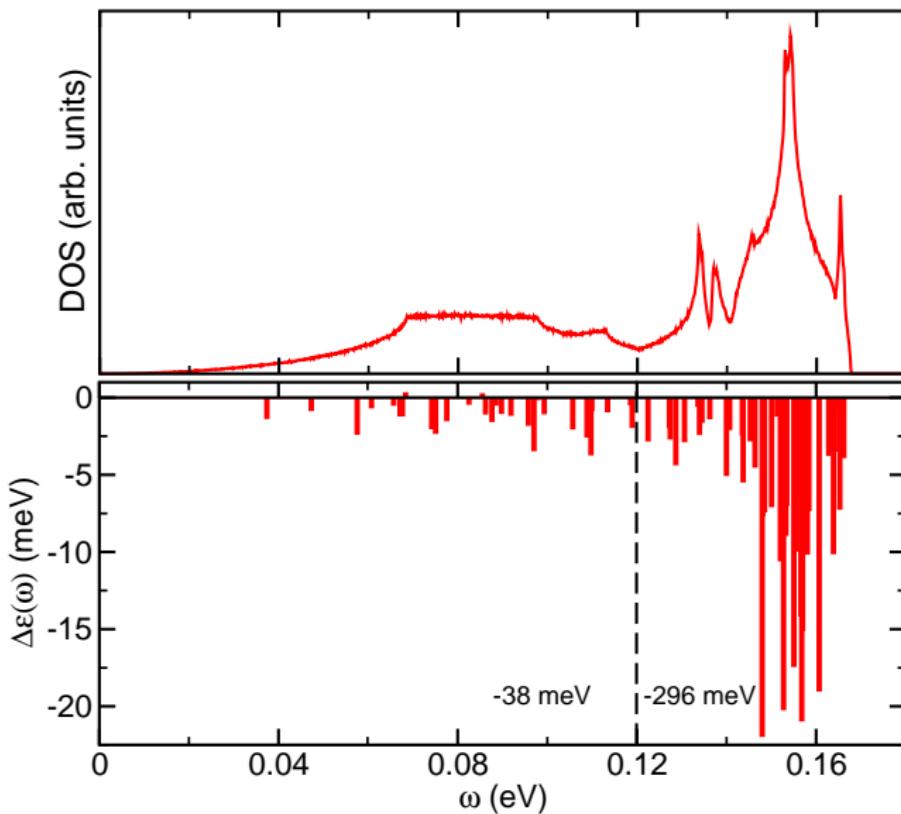
→ Monserrat and Needs, Phys. Rev. B **89**, 214304 (2014)

# Diamond phonon dispersion



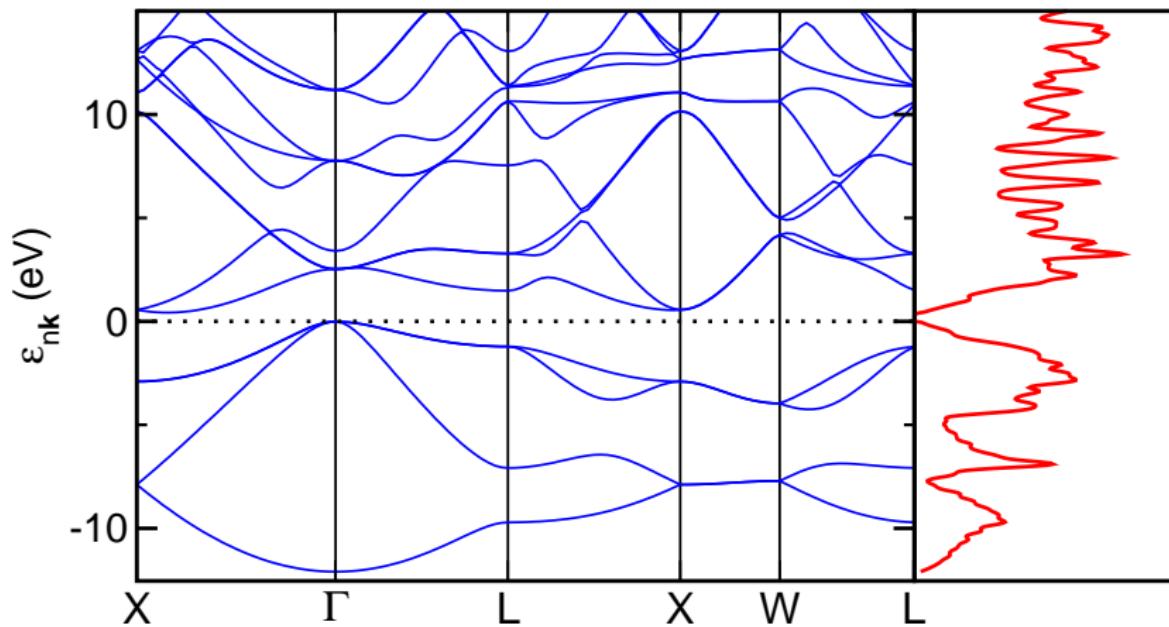
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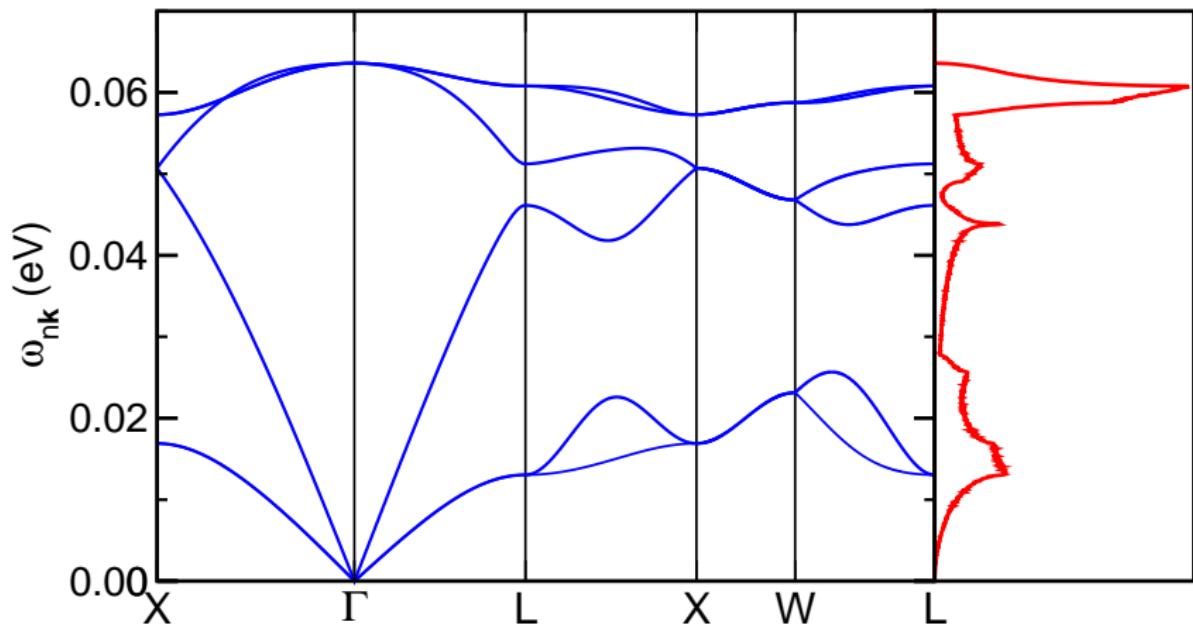
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# Silicon band structure



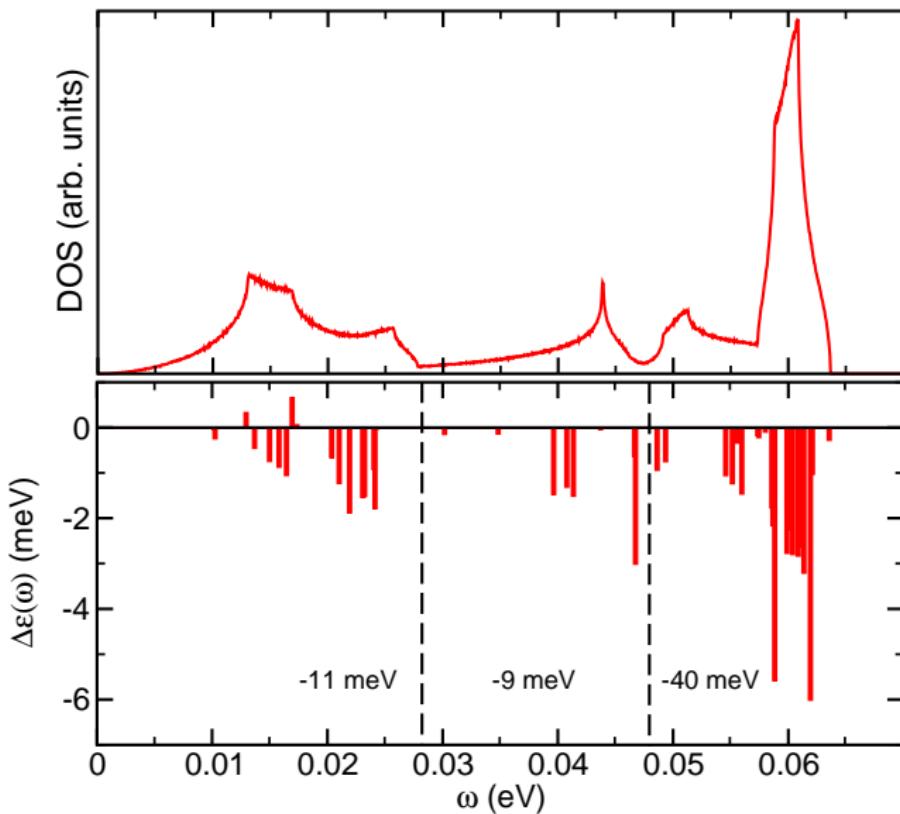
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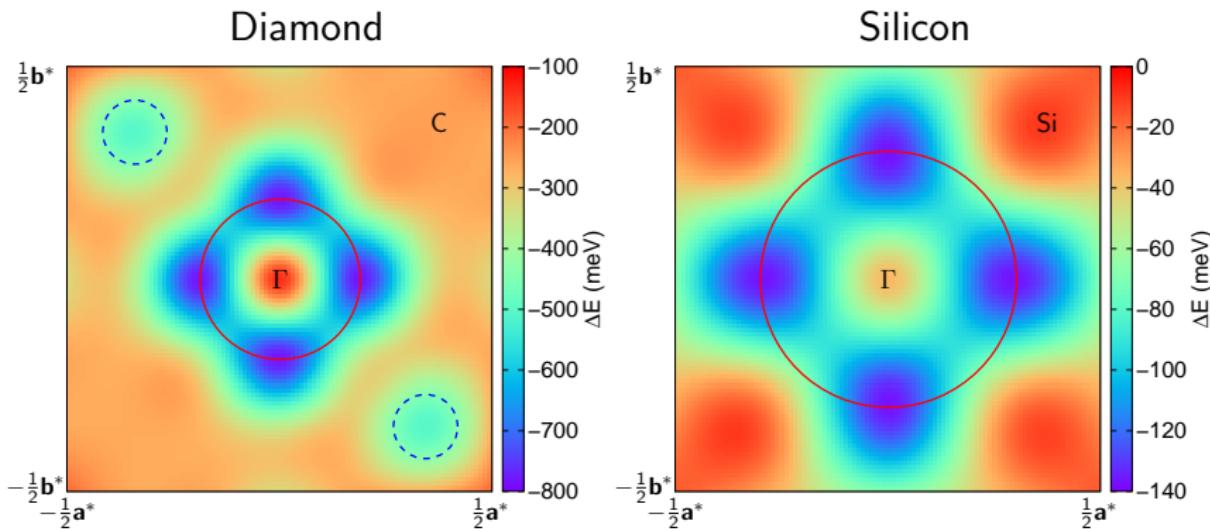
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# Silicon electron-phonon coupling



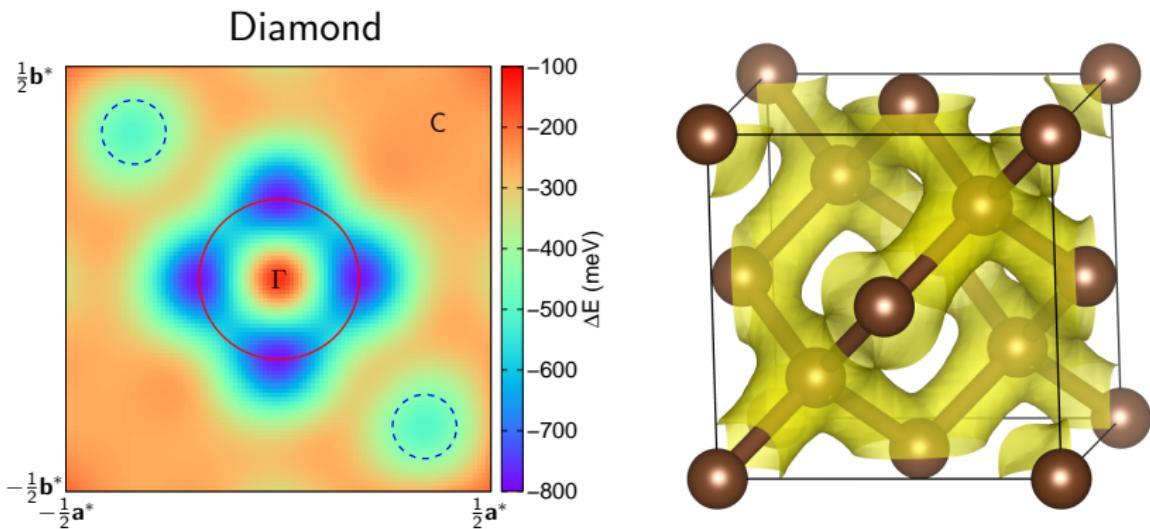
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# Electron-phonon coupling in diamond and silicon



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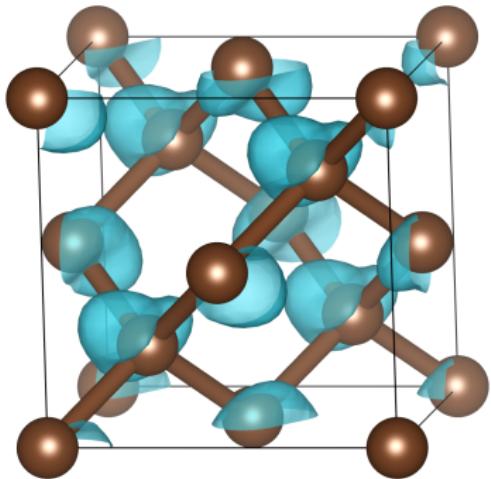
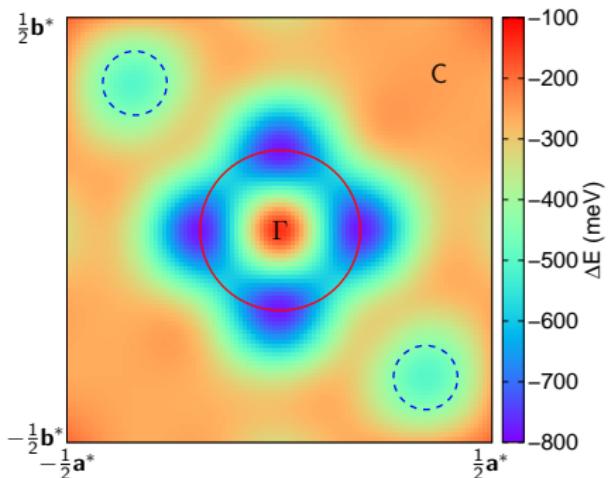
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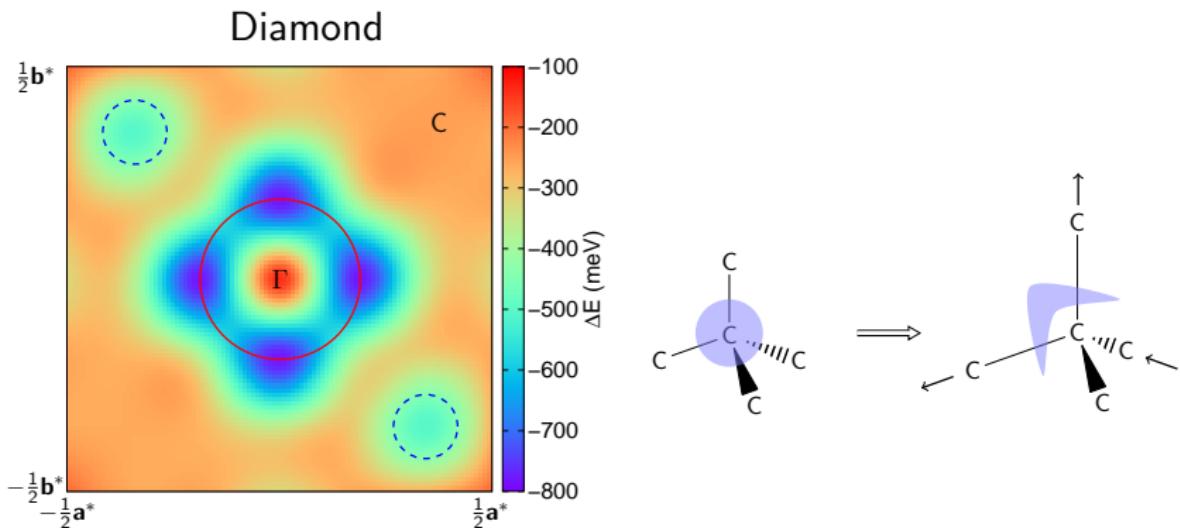
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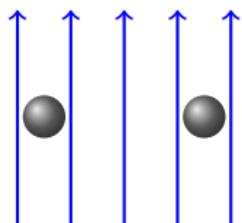
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# Nuclear magnetic resonance (NMR)



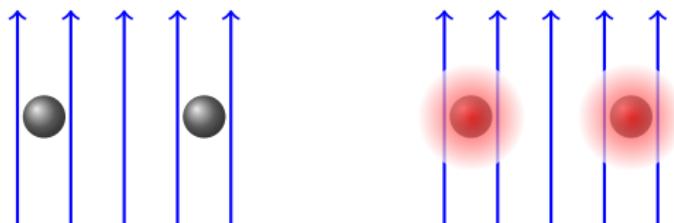
# Nuclear magnetic resonance (NMR)

Magnetic field



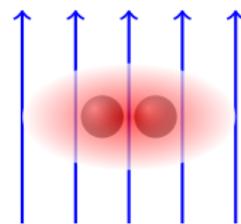
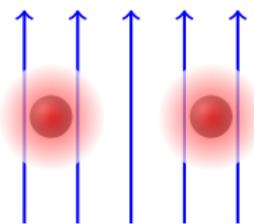
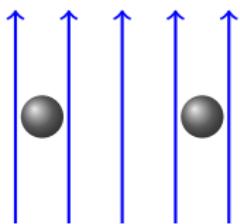
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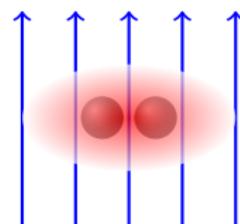
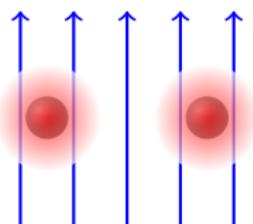
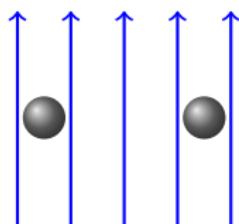
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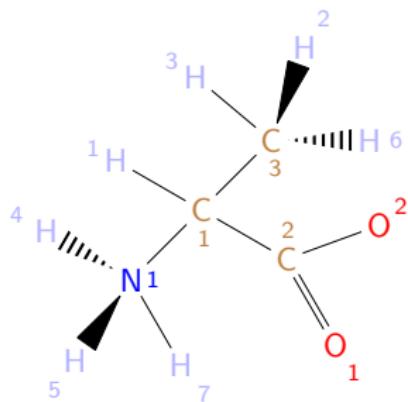
$$\mathbf{B} = \mathbf{B}_{\text{ext}} - \mathbf{B}_{\text{ind}} = (1 - \sigma) \mathbf{B}_{\text{ext}}$$

$$\mathbf{B}_{\text{ind}} = \sigma \mathbf{B}_{\text{ext}}$$

$$\mathbf{B}_{\text{ind}}(\mathbf{r}) = \frac{1}{c} \int d^3 \mathbf{r}' \mathbf{j}(\mathbf{r}') \times \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3}$$

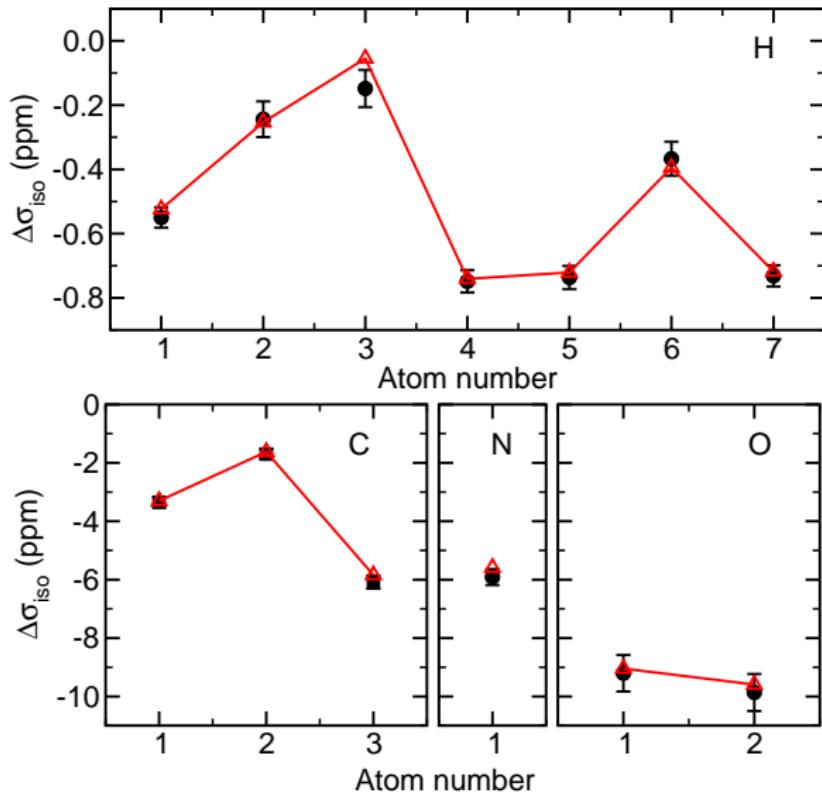
# L-alanine molecular crystal

L-alanine

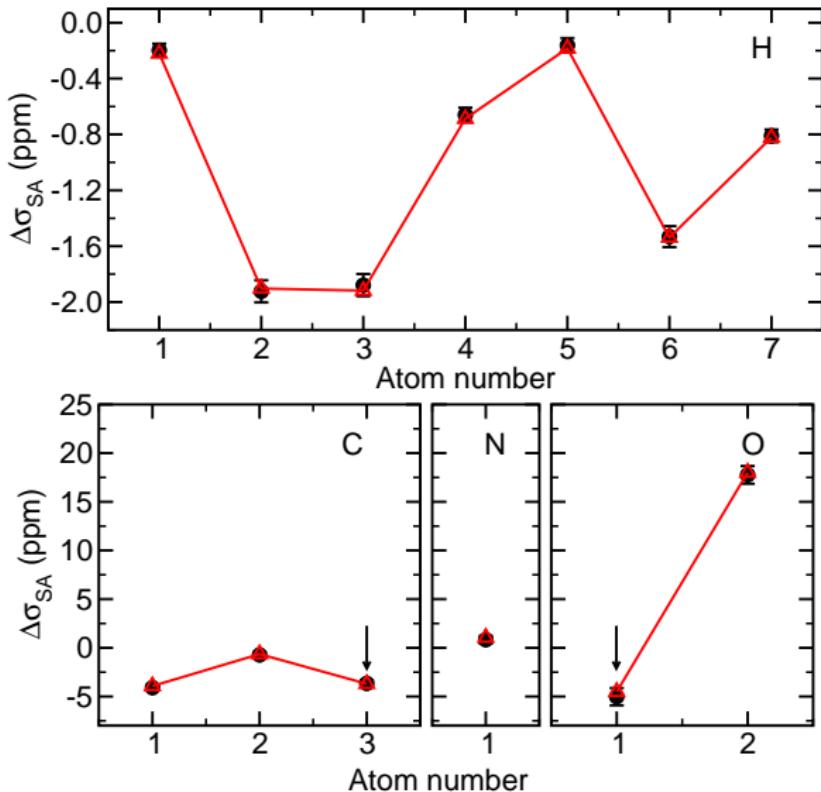


- ▶ Orthorhombic space group  $P2_12_12_1$  (4 molecules).
- ▶ L-alanine ( $C_3H_7NO_2$ ) has 52 atoms in the primitive cell.

# Isotropic shift

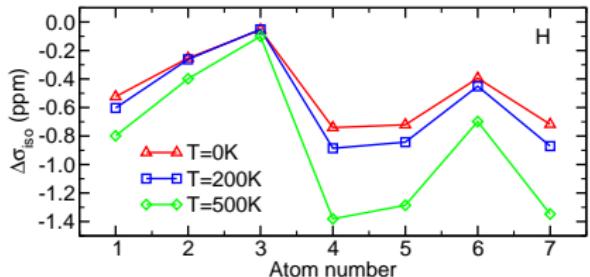


## Shielding anisotropy

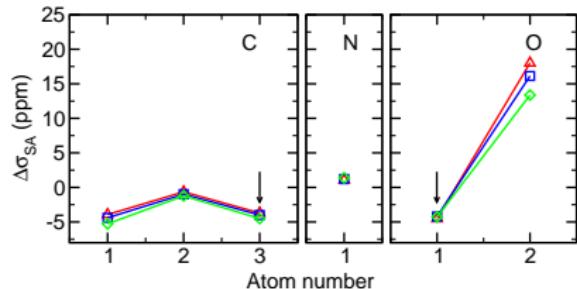
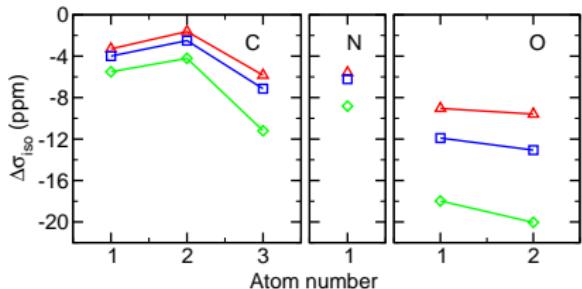
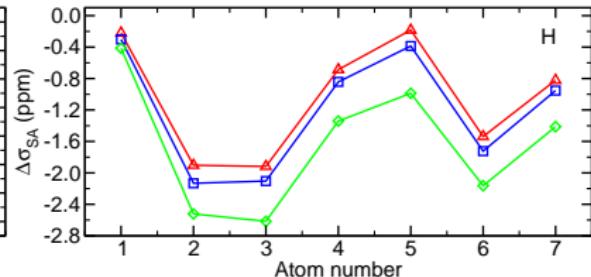


# Temperature dependence

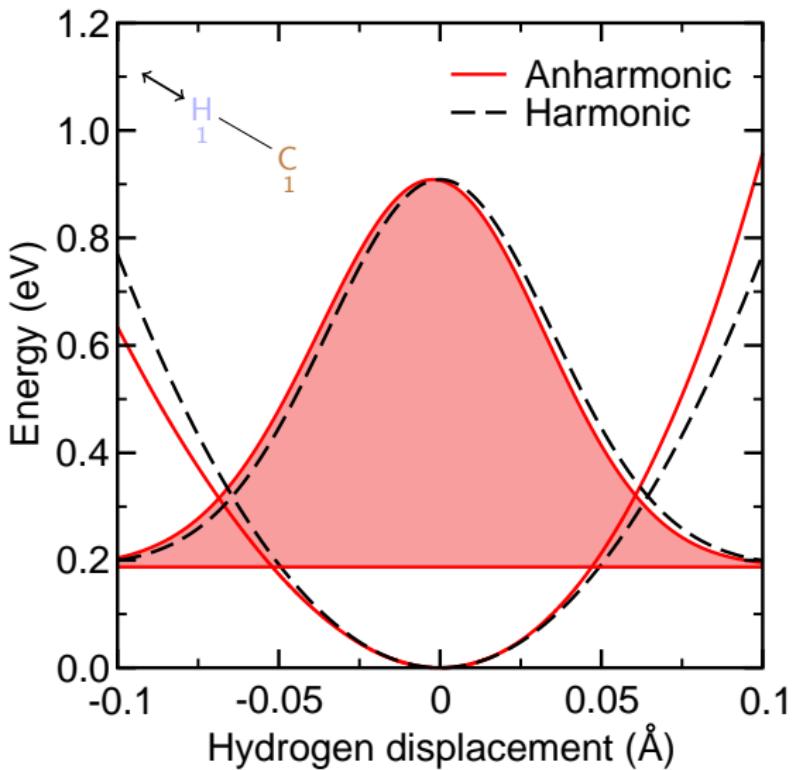
Isotropic shift



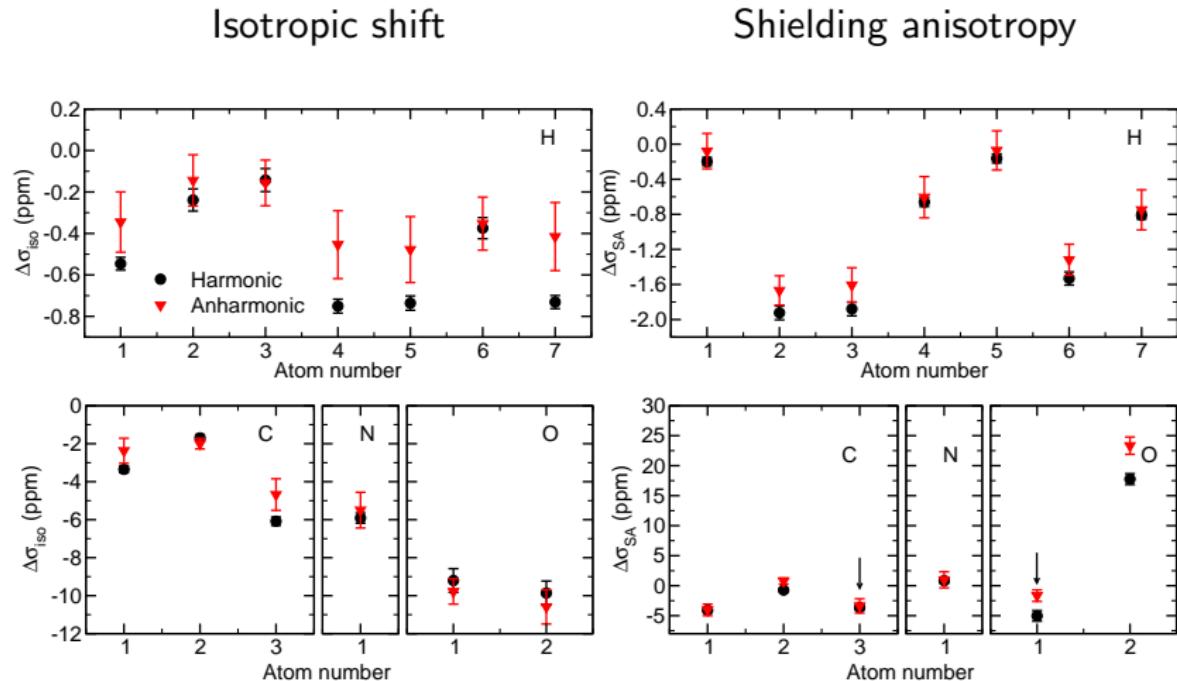
Shielding anisotropy



## Anharmonic vibrations



# Anharmonic coupling to the shielding tensor



# Summary

- ▶ Anharmonic energy:
  - ▶ Vibrational self-consistent field method for solids.
  - ▶ Phase diagram of high-pressure solid molecular hydrogen.
  - ▶ Relative stability of hexagonal and cubic ice.
- ▶ Vibrational coupling:
  - ▶ Quadratic expansion and Monte Carlo.
  - ▶ Electron-phonon coupling in semiconductors.
  - ▶ Vibrational effects on the chemical shielding tensor.

# Outlook

- ▶ Anharmonic vibrations can be important:
  - ▶ Small enthalpy differences between competing phases.
  - ▶ Harmonic instabilities.
  - ▶ High temperatures.
  - ▶ ...
- ▶ Vibrational coupling can be important:
  - ▶ Zero-point correction and finite temperature.
  - ▶ Any quantity available at the static level can be coupled to vibrations.
  - ▶ Electronic band gaps, NMR, ...

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