

# Generalized Local Density Approximation in One Dimension

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# My collaborator (and former boss)

## Prof. Peter Gill

- Professor at ANU since 2004
- Nottingham (1999-2004)
- Cambridge (1996-1999)
- Massey University (1993-1996)
- Postdoc with Pople (1989-1993)



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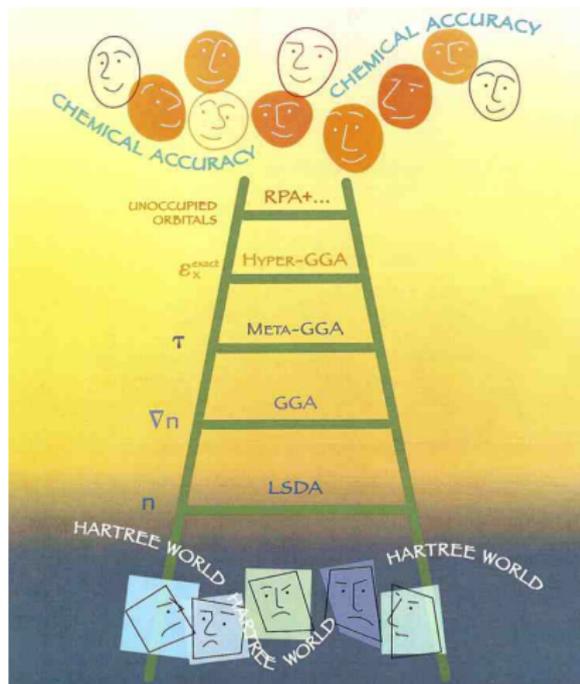
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- 😞😞 Not very accurate for correlation energy (**overestimated by roughly 200%**)

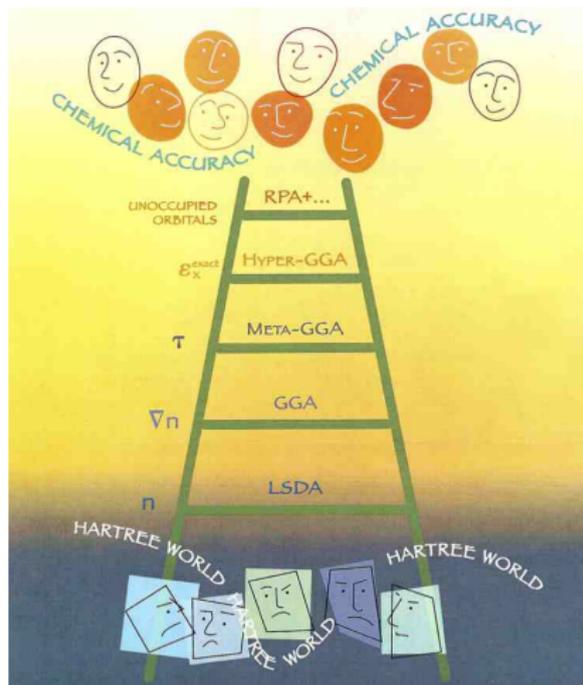
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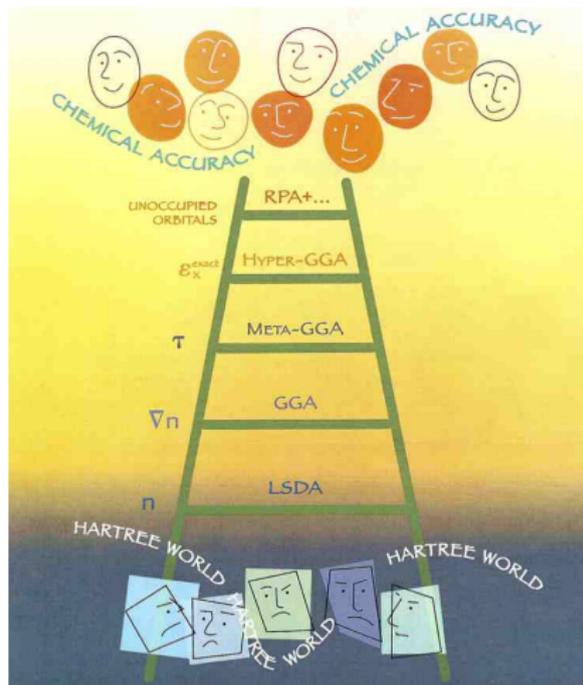
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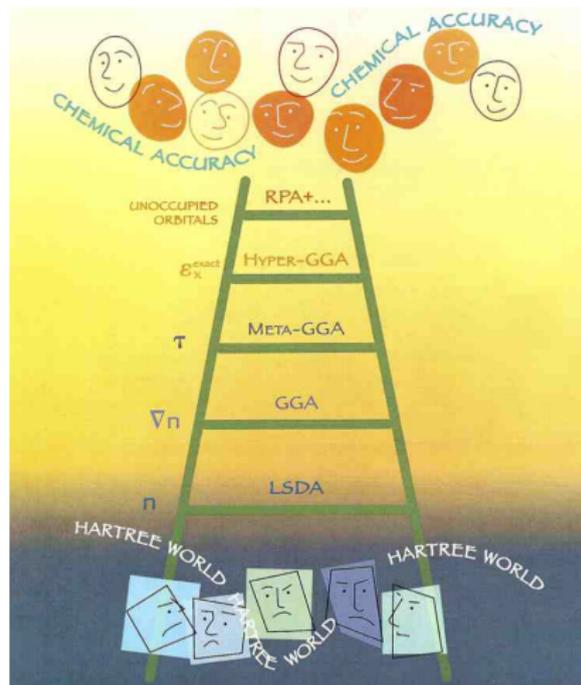
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- We propose to follow an alternative route to **heaven!**
- We add a new **local two-electron parameter  $\eta$**



## Generalized LDA: Hole curvature

- Suppose that an electron is at the point  $\mathbf{r}$ .
- The probability that another electron lies at a distance  $u$  is

$$P(u|\mathbf{r}) = \frac{\int \rho_2(\mathbf{r}, \mathbf{r} + \mathbf{u}) d\Omega_u}{\rho(\mathbf{r})}$$

where  $\Omega_u$  is the angular part of  $u$  and  $\rho_2(\mathbf{r}_1, \mathbf{r}_2)$  is the reduced 2nd-order density matrix

- $P''(0|\mathbf{r})$  indicates the **width of the hole** around the electron at  $\mathbf{r}$ .
- Therefore, the **dimensionless curvature**

$$\eta(\mathbf{r}) = r_s(\mathbf{r})^3 P''(0|\mathbf{r})$$

measures of the proximity of other electrons to an electron at  $\mathbf{r}$ .

# Uniform electron gases

## Infinite UEG

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  - 4 **In the limit** as  $n \rightarrow \infty$  and  $V \rightarrow \infty$ , one obtains an infinite UEG

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## Finite UEGs

- One can also construct UEGs using a **finite number of electrons**

Loos & Gill, *J Chem Phys* 135 (2011) 214111  
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😊😊😊 For  $n \rightarrow \infty$ , we get the **infinite** UEG!!

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## A Few Finite UEGs

$\mathcal{D}$	System	Name
1	$n$ electrons on a ring	$n$ -ringium
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# Ringium: “— One Ring to Rule Them All —”

## Electrons on a Ring



## Wavefunctions & Energies

$$\hat{H} = -\frac{1}{2R^2} \sum_{i=1}^n \frac{\partial^2}{\partial \theta_i^2} + \sum_{i < j}^n \frac{1}{r_{ij}}$$

$$\epsilon = ? \quad \Psi = ?$$

$$r_s = \frac{1}{2\rho} = \frac{\pi R}{n}$$

Remarque: the electrons interact “through” the ring

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When  $r_{12} \rightarrow 0$ ,

$$\Psi(r_{12}) = r_{12} \left( 1 + \frac{r_{12}}{2} \right) + O(r_{12}^3)$$

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- The system is **spin-blind**:
  - The divergence of the Coulomb operator mimics the **Pauli principle**
  - **Ferromagnetic** and **paramagnetic** states are degenerate
  - The **fermonic** and **bosonic** states are degenerate (**Bose-Fermi mapping**)

Lee & Drummond, *Phys Rev B* 83 (2011) 245114

Astrakharchik & Girardeau, *Phys Rev B* 83 (2011) 153303

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- The HF energy is

$$\epsilon_{\text{HF}} = \frac{n^2 - 1}{n^2} \frac{\pi^2}{24 r_s^2} + \frac{1}{4 r_s} \left( \sum_{k=1}^n \frac{4 - 1/n^2}{2k - 1} - 3 \right)$$

Loos & Gill, *J Chem Phys* 138 (2013) 164124

## 1st weapon: Perturbation expansions

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- We can also find the **low-density expansion** coefficients (valid for  $r_s \gg 1$ )

$$\varepsilon_c = \frac{\gamma_2}{r_s} + \frac{\gamma_3}{r_s^{3/2}} + \dots$$

Loos, *J Chem Phys* 138 (2013) 064108

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- 😊 Hylleraas-type calculations can be done for few electrons
- 😊 It works well for **intermediate**  $r_s$
- 😞 **However**, the many-electron integrals are too numerous and too difficult for larger number of electrons
  - So how can we calculate accurate energies for **intermediate**  $r_s$ ?

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😊😊😊 Fortunately, the **HF nodes** (*i.e.*  $r_{ij} = 0$ ) are **exact!** **Youpi!**

$$\Psi_{\text{Trial}}(\{r_{ij}\}) = \Psi_{\text{HF}}(\{r_{ij}\}) \prod_{i < j}^n \left( \sum_k c_k r_{ij}^k \right)$$

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- The Schrödinger eqn is separable in **extracule** & **intracule** coordinates
- The **extracule equation** is trivial to solve
- The **intracule equation** is a Heun-type differential equation
- For certain “**eigenradi**”  $R$ , both  $\epsilon$  and  $\Psi$  can be obtained in closed form
- There are a **countably infinite** number of these closed-form solutions

Loos & Gill, *Phys Rev Lett* 103 (2009) 123008; *ibid* 108 (2012) 083002

## Some exact solutions

State	$R$	$\epsilon$	$\Psi(r_{12})$	$x = r_{12}/(2R)$
Ground	1/2	9/4	$r_{12}\sqrt{1+x}$	
	$\sqrt{3}/2$	2/3	$r_{12} [1 + \frac{1}{2}r_{12}]$	
	$\frac{1}{4}(\sqrt{33} + 3)$	$\frac{25}{96}(7 - \sqrt{33})$	$r_{12}\sqrt{1+x} [1 + (R - \frac{1}{2})x]$	
	$\sqrt{23}/2$	9/46	$r_{12} [1 + \frac{1}{2}r_{12} + \frac{5}{2}x^2]$	
⋮	⋮	⋮	⋮	
1st excited	$\frac{1}{4}(\sqrt{33} - 3)$	$\frac{25}{96}(7 + \sqrt{33})$	$r_{12}\sqrt{1-x} [1 + (R + \frac{1}{2})x]$	
	$\sqrt{5}/2$	9/10	$r_{12}\sqrt{1-x}\sqrt{1+x} [1 + \frac{1}{2}r_{12}]$	
	$\sqrt{33}/2$	8/33	$r_{12}\sqrt{1-x}\sqrt{1+x} [1 + \frac{1}{2}r_{12} + \frac{7}{2}x^2]$	
⋮	⋮	⋮	⋮	

Loos & Gill *Phys Rev Lett* 108 (2012) 083002

## Reduced correlation energies ( $mE_h$ ) for $n$ -ringium

Combining perturbation expansions, Hylleraas and DMC calculations leads to

$n$	$6\eta/\pi^2$	Seitz radius $r_S$								
		0	0.1	0.2	0.5	1	2	5	10	20
2	3/4	13.212	12.985	12.766	12.152	11.250	9.802	7.111	4.938	3.122
3	8/9	18.484	18.107	17.747	16.755	15.346	13.179	9.369	6.427	4.030
4	15/16	21.174	20.698	20.249	19.027	17.324	14.762	10.390(0)	7.085(0)	4.425(0)
5	24/25	22.756	22.213	21.66(2)	20.33(1)	18.439(1)	15.644(2)	10.946(0)	7.439(0)	4.636(0)
6	35/36	23.775	23.184	22.63(2)	21.14(1)	19.137(1)	16.192(2)	11.285(0)	7.653(0)	4.762(0)
7	48/49	24.476	23.850	23.24(2)	21.70(1)	19.607(1)	16.554(2)	11.509(0)	7.795(0)	4.844(0)
8	63/64	24.981	24.328	23.69(3)	22.11(1)	19.940(1)	16.808(2)	11.664(0)	7.890(0)	4.901(0)
9	80/81	25.360	24.686	24.04(2)	22.39(1)	20.186(1)	16.995(3)	11.777(0)	7.960(0)	4.941(0)
10	99/100	25.651	24.960	24.25(4)	22.62(1)	20.373(1)	17.134(2)	11.857(0)	8.013(0)	4.973(0)
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
∞	1	27.416	26.597	25.91(1)	23.962(1)	21.444(0)	17.922(0)	12.318(0)	8.292(0)	5.133(0)

Lee & Drummond, *Phys Rev B* 83 (2011) 245114

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[Cioslowski, \*J Chem Phys\* 136 \(2012\) 044109](#)
  - 2 Fitting based on hypergeometric functions  
(maybe related to the ISI functional of Seidl and Perdew)  
[Seidl, Perdew & Kurth, \*Phys Rev Lett\* 84 \(2000\) 5070](#)

# Take-home messages

## How can we use these new UEG results?

- $n$  electrons on a ring gives UEGs of any desired density  $\rho$
- We have calculated their correlation energies **very accurately**
- Our results permit a **generalization of the LDA** for finite systems
- This improves the accuracy of the **lowest rung** of Jacob's Ladder
- **Next**, we will extend this approach to electrons in 2D and 3D