

A Full Configuration Interaction QMC Perspective on the Homogeneous Electron Gas

James Shepherd

Quantum Monte Carlo in the Apuan Alps
31st July 2012

Why?

- A real challenge for FCI(QMC)
- ...a new method in fully-periodic calculations
- Model systems can give physical insight without added complexity of real systems
- ...an application with development in mind
- Exact (finite/complete basis) energy benchmarks
- What does the wavefunction look like?

70 years of history...

- 1934 Wigner crystal
- 1956 Landau theory of Fermi liquids
- 1957 Gell-Mann & Brueckner RPA
- 1962 Overhauser spin/charge density waves
- 1980 Ceperley-Alder QMC

...and so much more.

DMC: very successful

Phase transitions, spectral functions, effective mass, QP renormalization factor

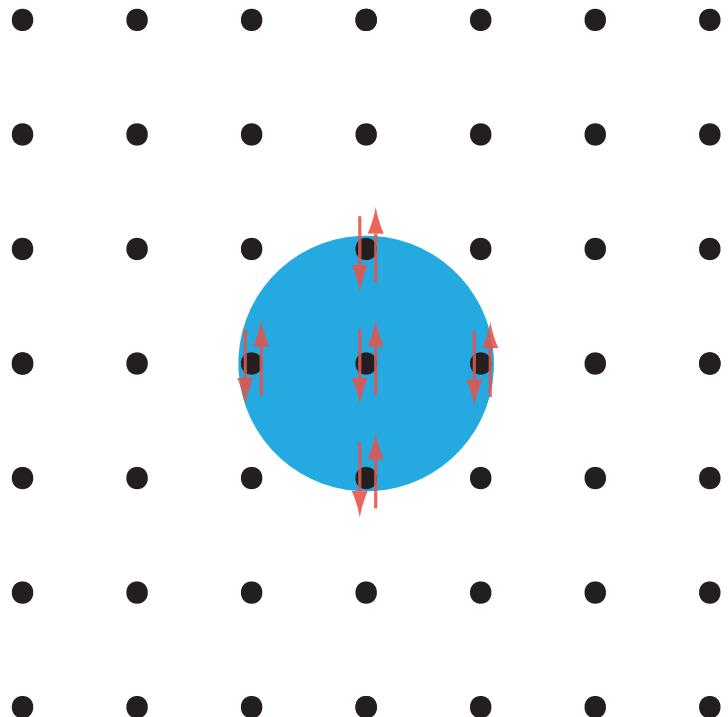
(Ceperley, Needs, Holzmann, Drummond, Foulkes, Ballone, Gurturbay etc. etc.)

Post-HF quantum chemistry: new studies of the solid state (e.g. LiH)

(Schutz, Manby, Kresse, Stoll, Fulde etc. etc.)

Open question: what does $|\Psi\rangle = \sum_i c_i |D_i\rangle$ look like?

$$\hat{H} = \sum_{\alpha} -\frac{1}{2} \nabla_{\alpha}^2 + \sum_{\alpha \neq \beta} \frac{1}{2} \hat{v}_{\alpha\beta} + \text{const.}$$



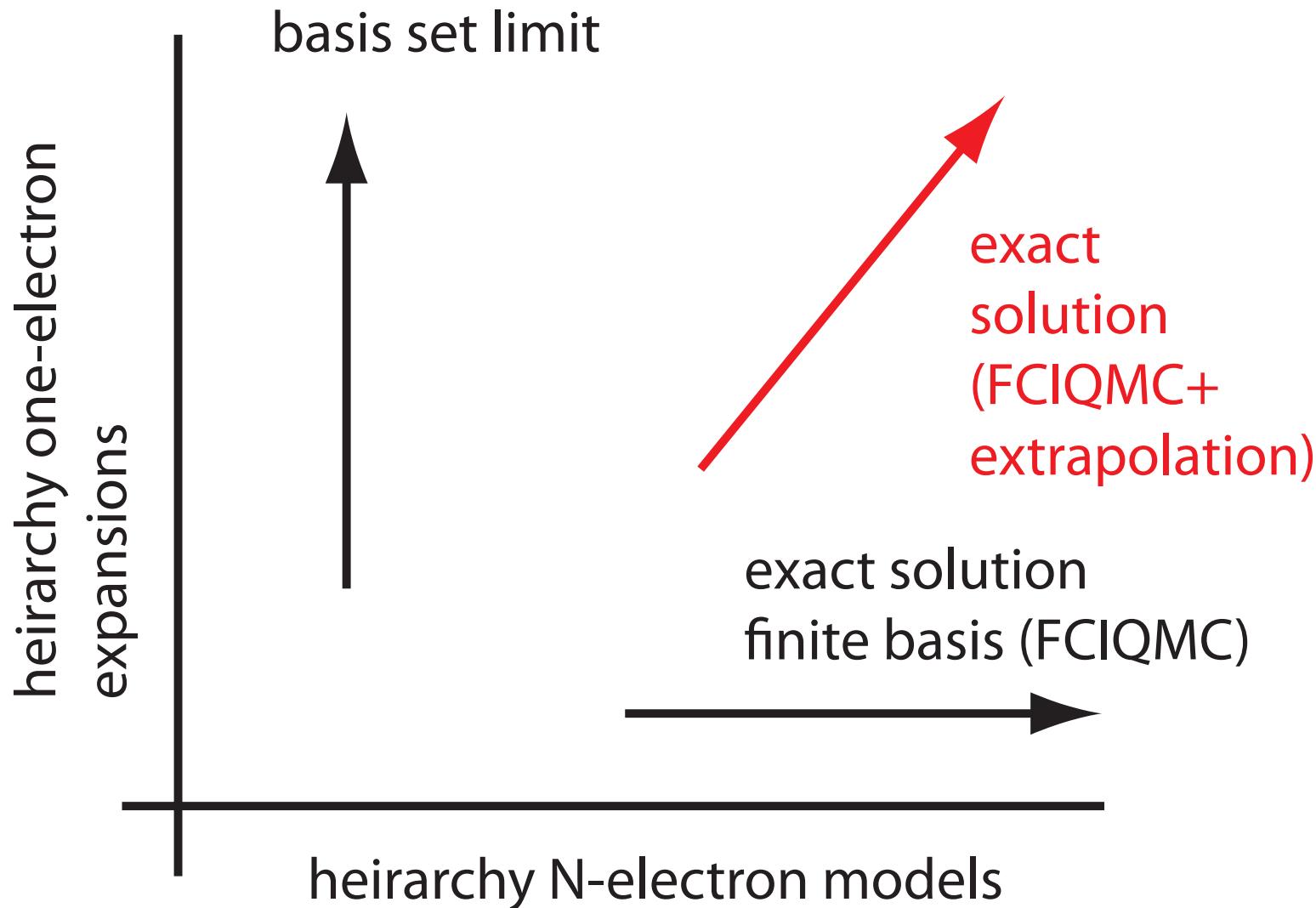
$$\psi_{\mathbf{k}} = \sqrt{\frac{1}{L^3}} e^{i\mathbf{k}\cdot\mathbf{r}} \delta_{s\sigma}$$

Finite particle number: N

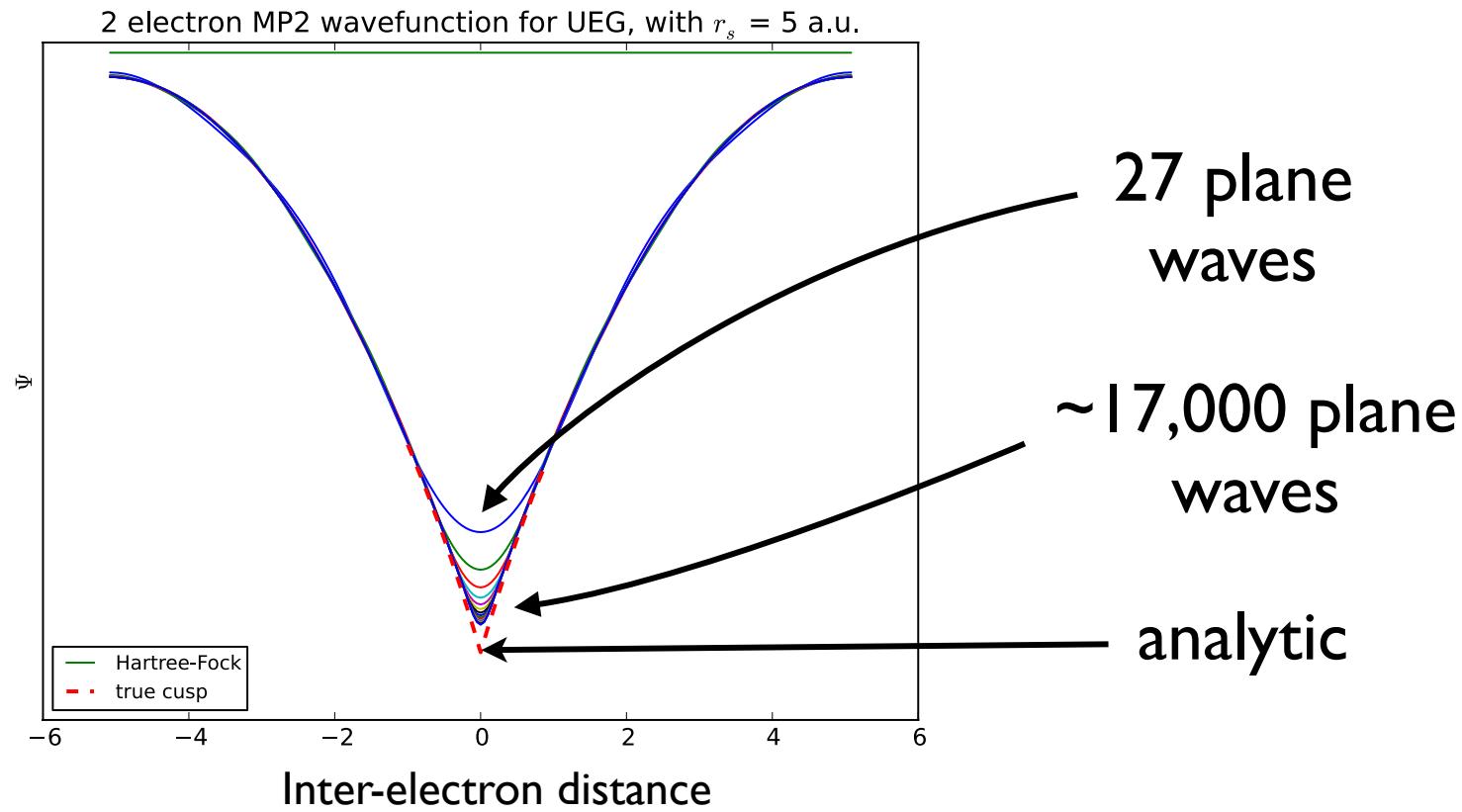
Uniform density: r_s

(I am discussing 3D but will show 2D schematics)

Using $|\Psi\rangle = \sum_i c_i |D_i\rangle$ to solve $\hat{H}\Psi = E\Psi$

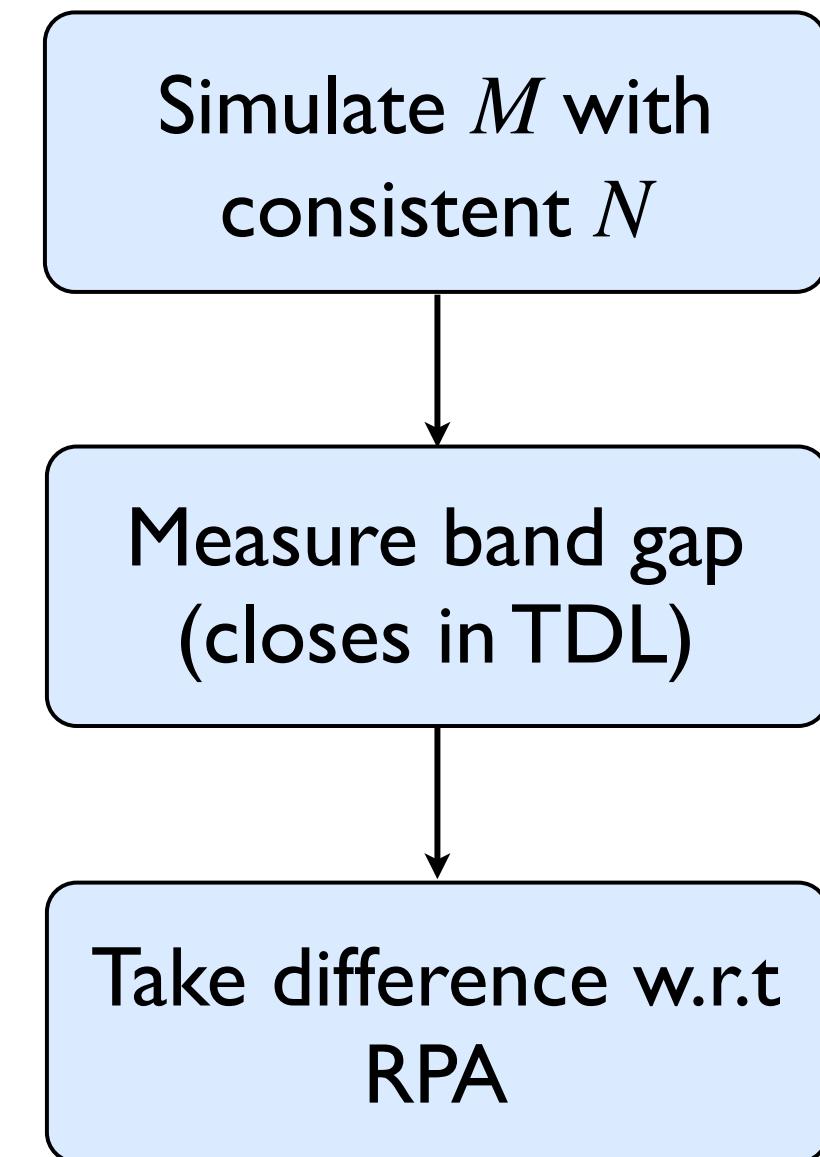
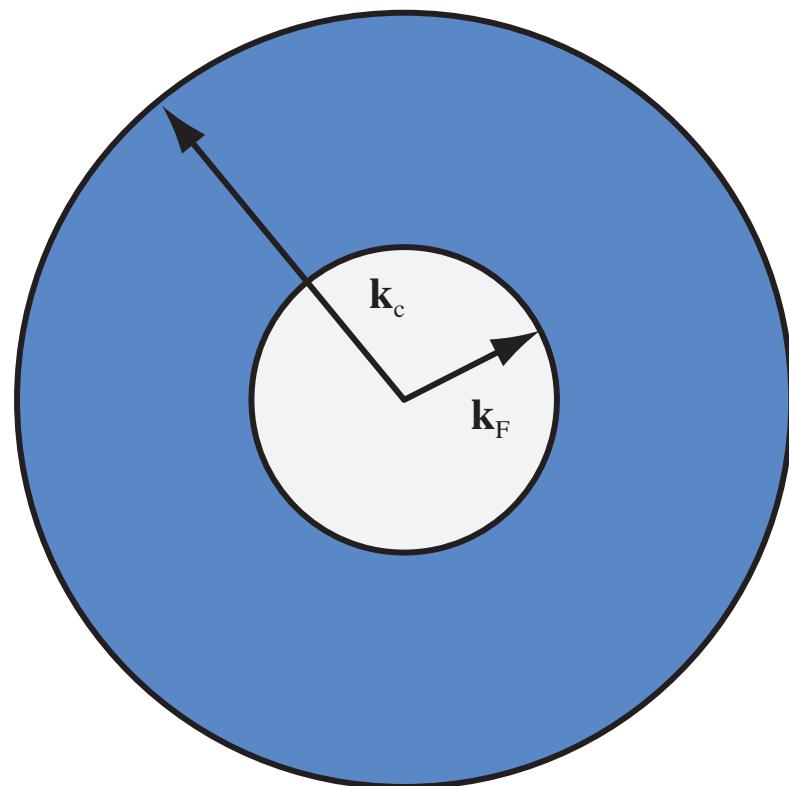


The problem with determinantal expansions...

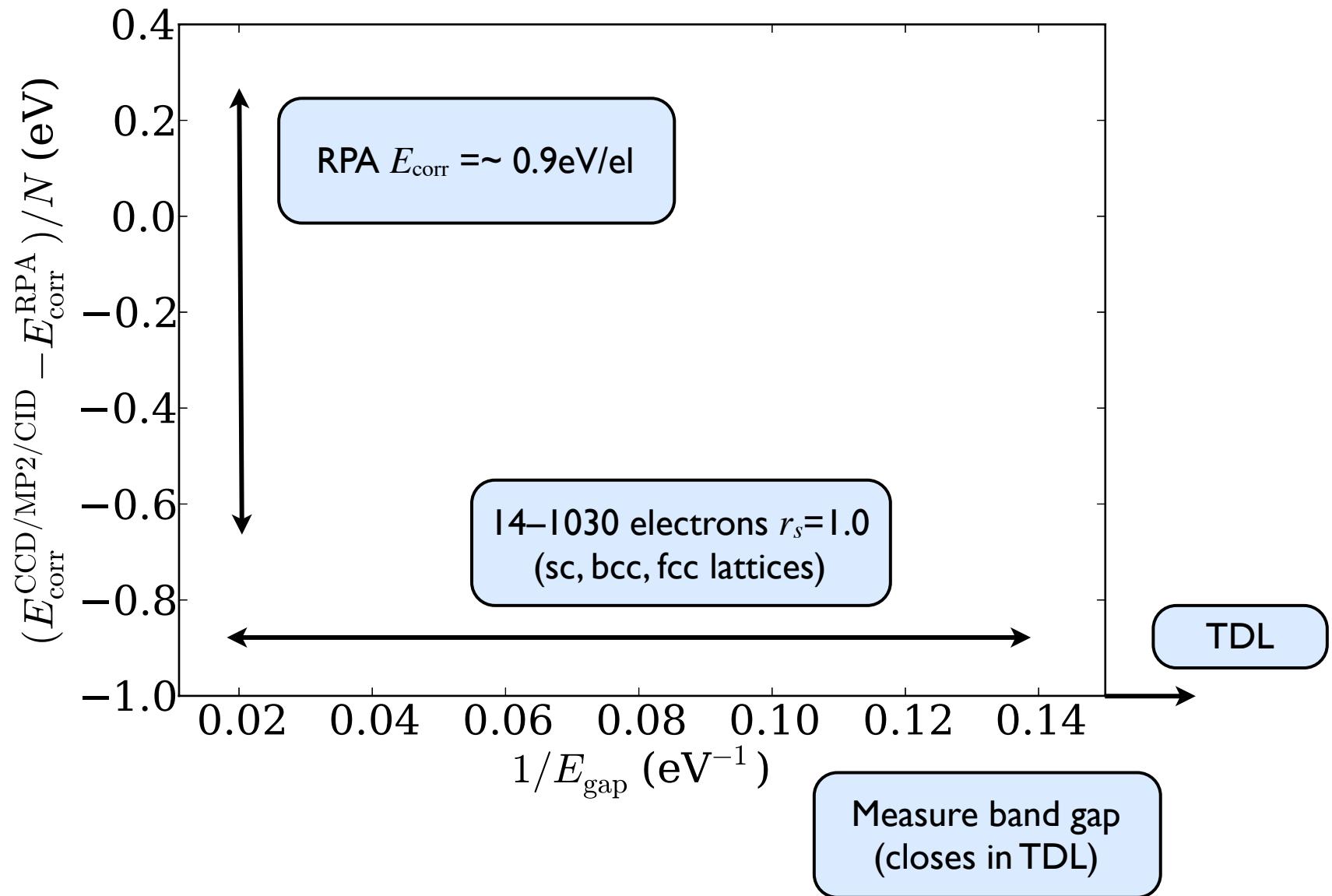


... simulation-cell HEG has large amounts of dynamic correlation at real metal r_s -values

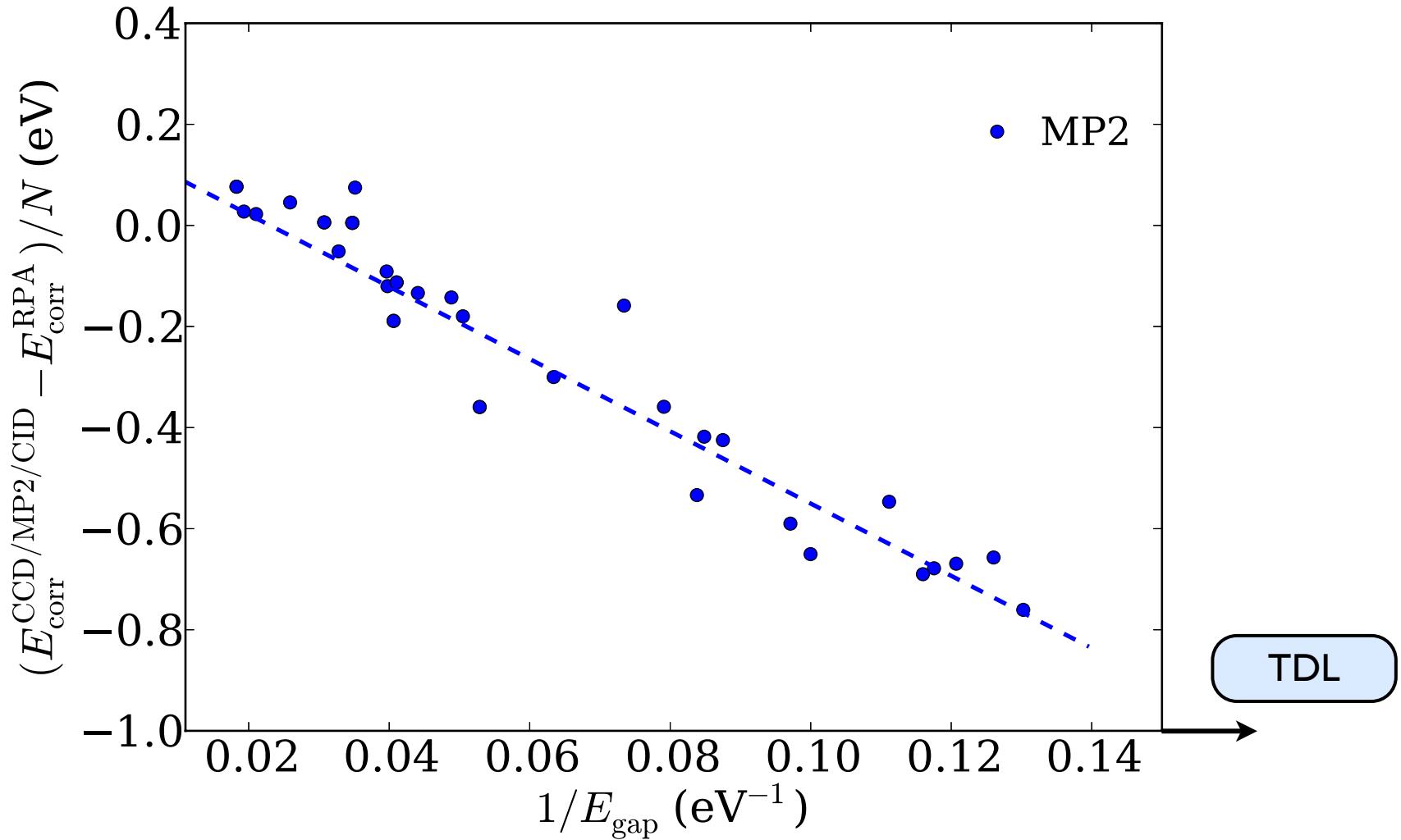
Perturbative approaches (MP2, CCD(T)) diverge...



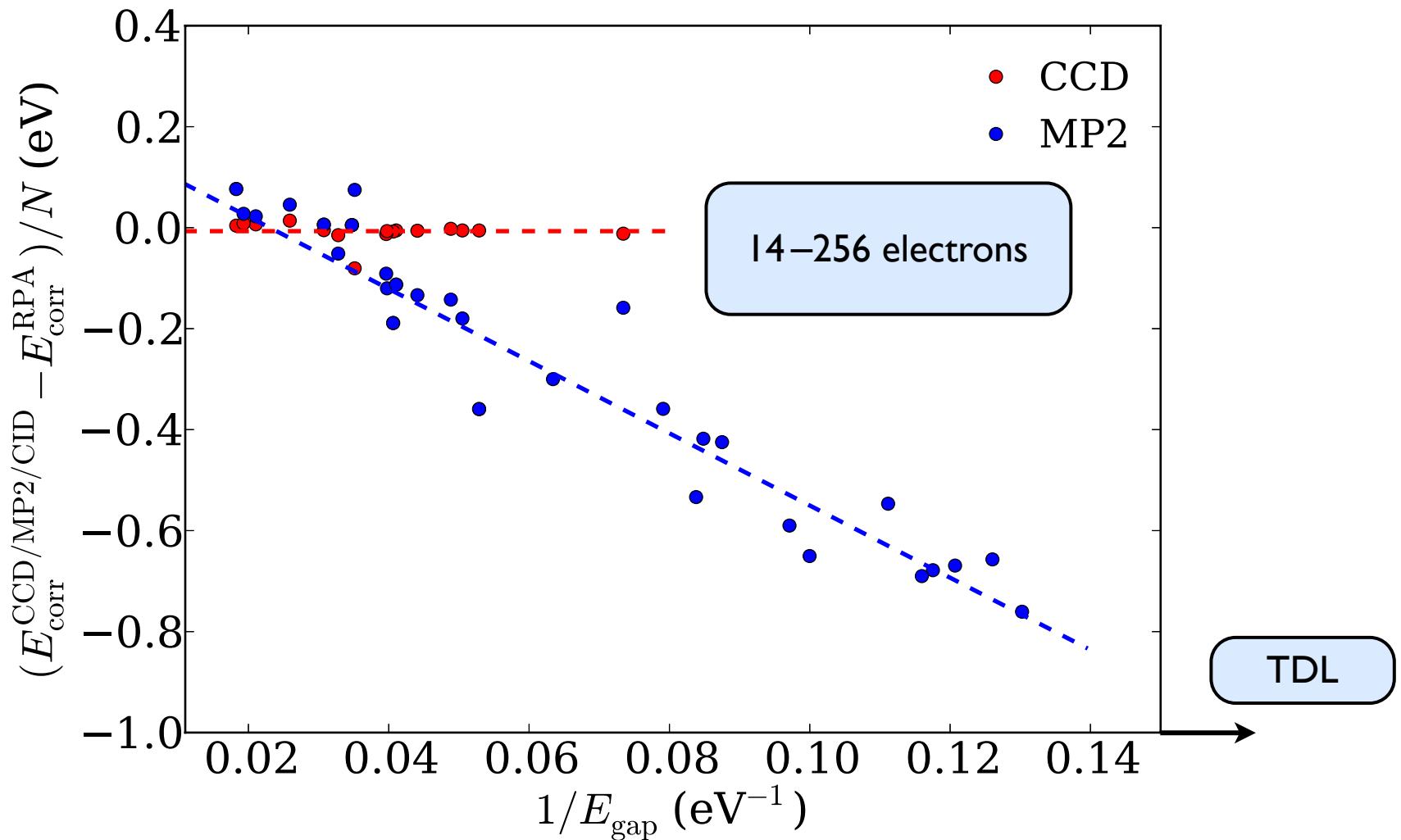
Take difference w.r.t
RPA



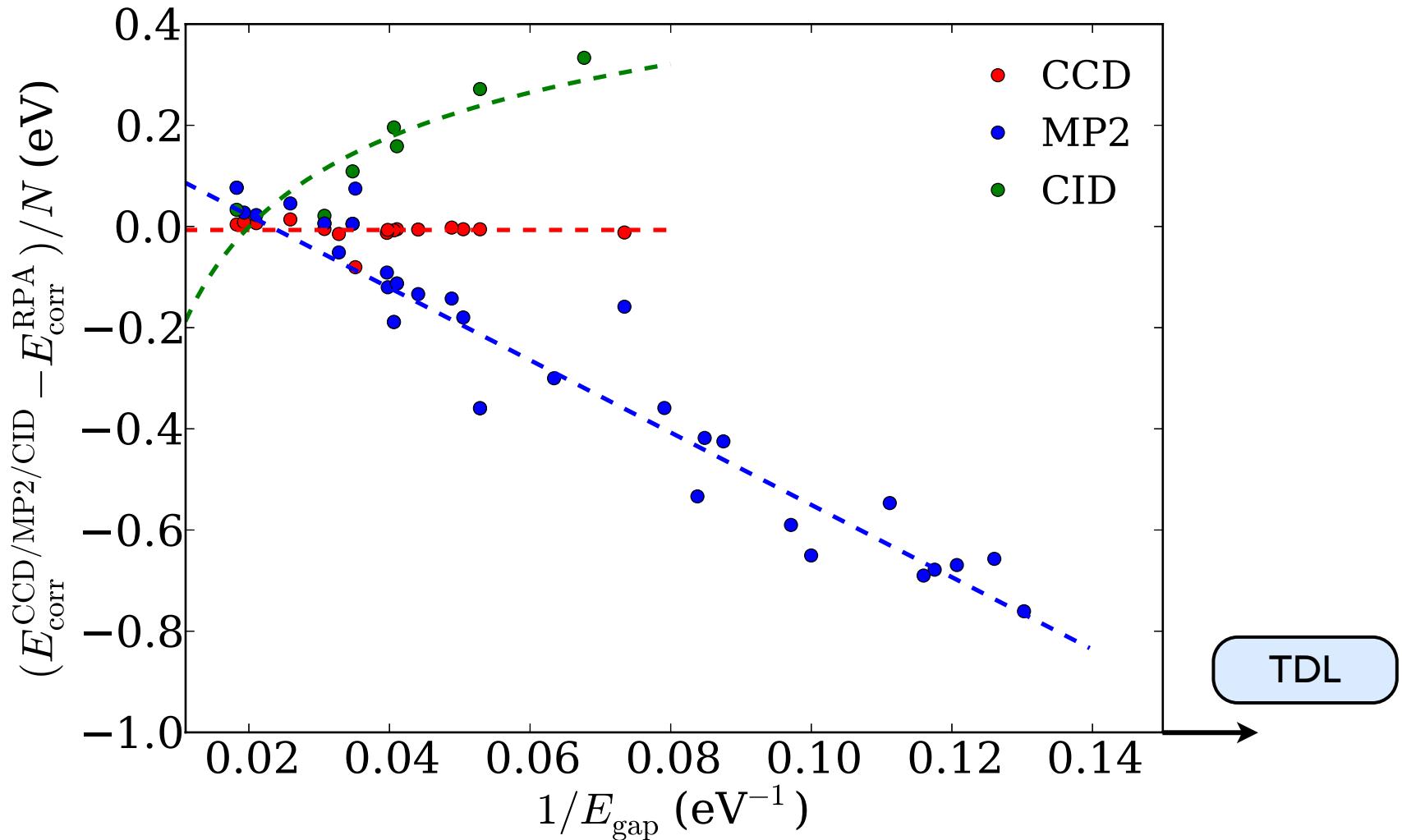
MP2 diverges...



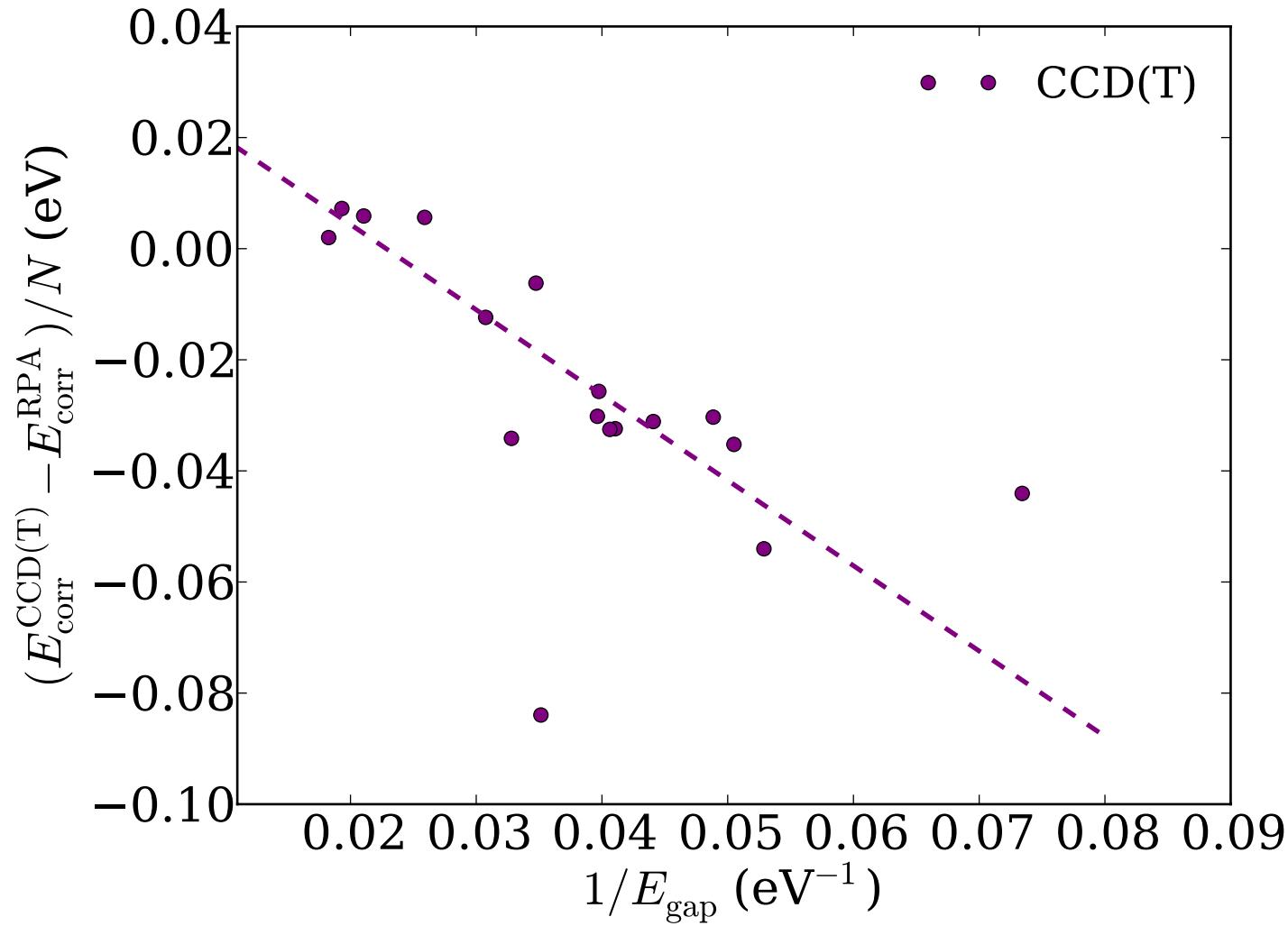
CCD converges...



CID goes to zero (C.E./elec)...



...and CCD(T) diverges.



FCIQMC

$$|\Psi\rangle = \sum_i c_i |D_i\rangle$$

+

$$\Psi_0 = \lim_{\tau \rightarrow \infty} e^{-\tau(\hat{H}-S)} \Psi_{\tau=0}$$



$$-\frac{dc_i}{d\tau} = (H_{ii} - S)c_i + \sum_{j \neq i} H_{ij}c_j$$

FCIQMC

Walker population

Spawning

$$-\overbrace{\frac{dc_i}{d\tau}}^{} = \underbrace{(H_{ii} - S)c_i}_{\text{Death/cloning}} + \overbrace{\sum_{j \neq i} H_{ij}c_j}^{\text{Spawning}}$$

Death/cloning

Instantaneous $\Psi(\tau), H_{ij}(\tau)$

Initiator FCIQMC (*i*-FCIQMC)

Coupling of $\Psi(\tau), H_{ij}(\tau)$

Only consider H_{ij} over those determinants with a certain population

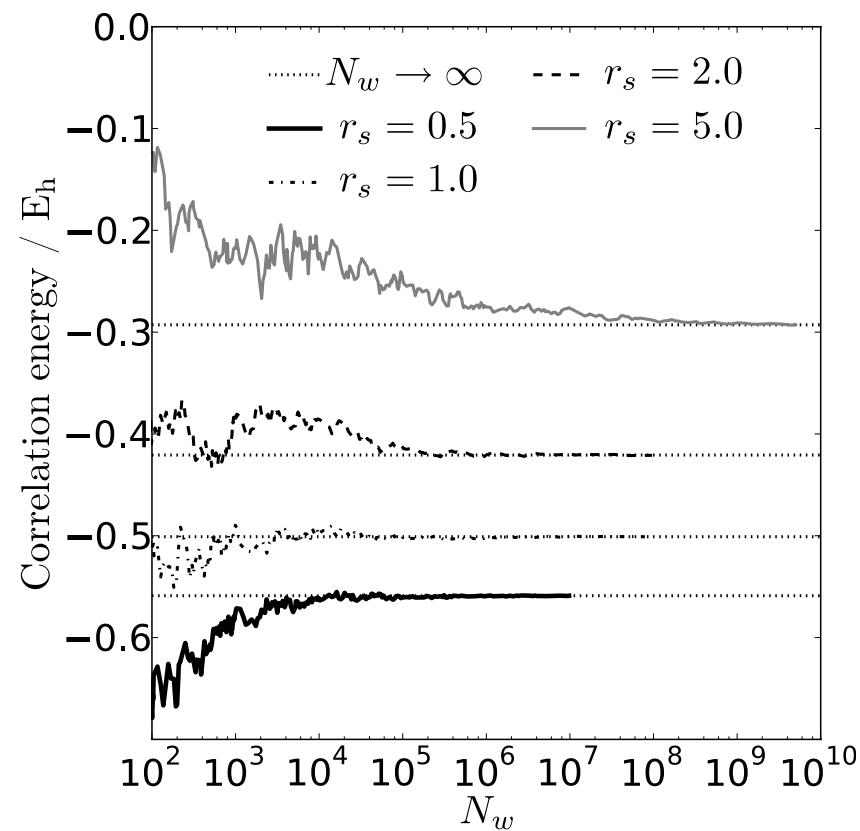
$$-\frac{dc_i}{d\tau} = (H_{ii} - S)c_i + \sum_{j \neq i} H_{ij}c_j$$

Stabilises S over a much larger range of walker populations

$N_w \rightarrow \infty$ limit must be found

Initiator FCIQMC (*i*-FCIQMC)

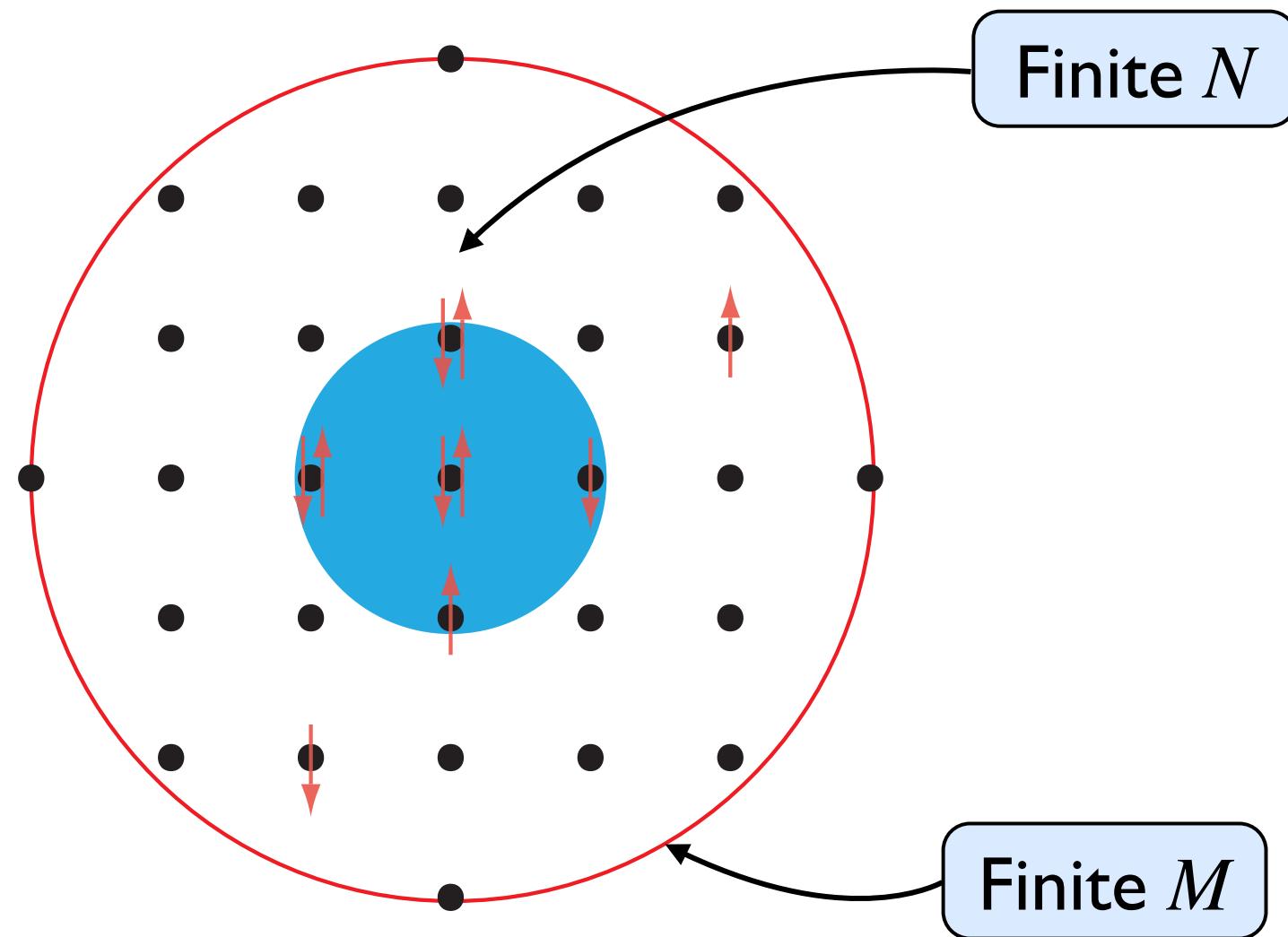
$N_w \rightarrow \infty$ limit must be found



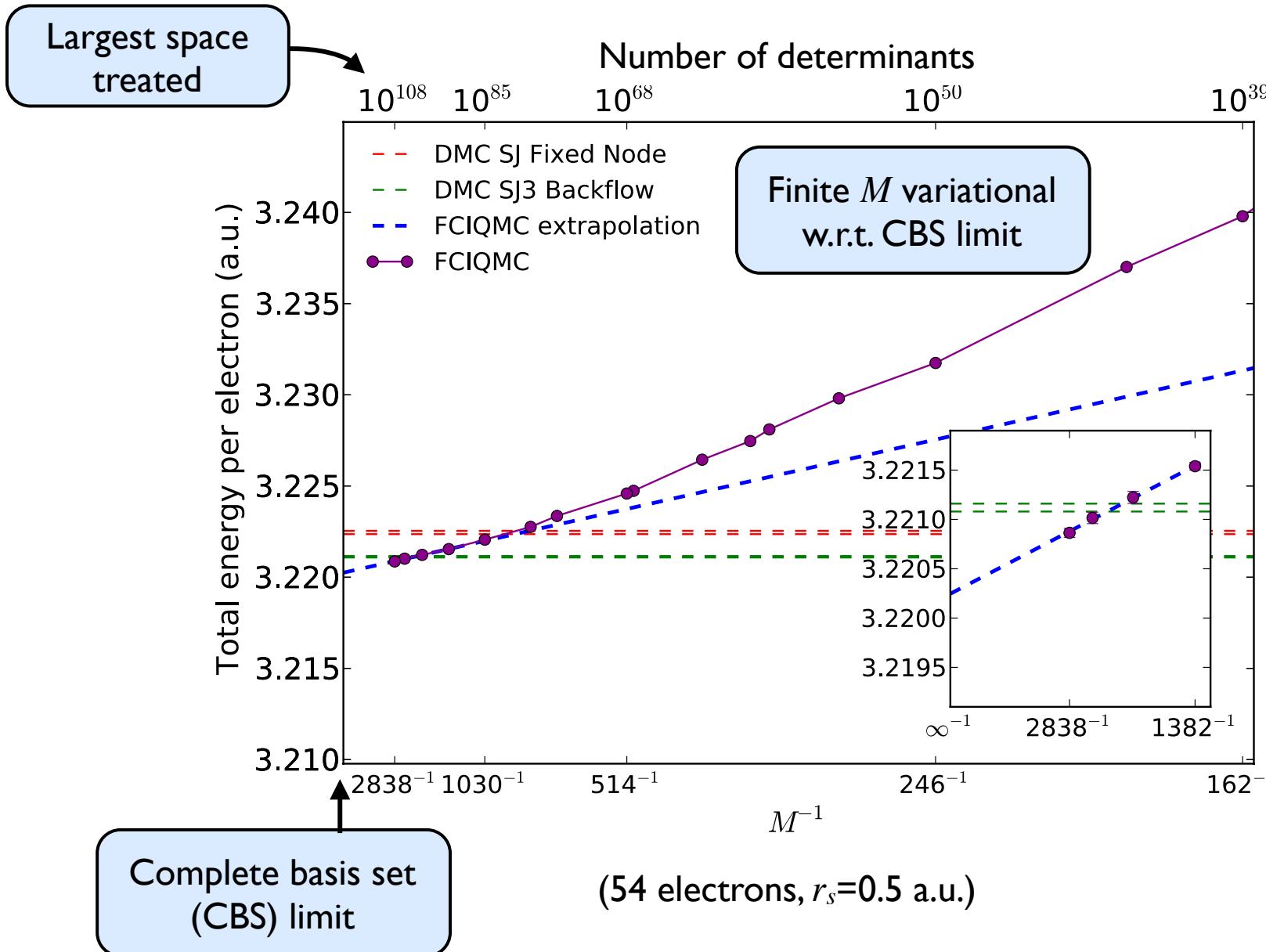
Overview

- Introduction
- Plane wave basis set incompleteness error
- Comparison with DMC and finite size effects
- Insights from the FCIQMC wavefunction

Plane wave basis sets

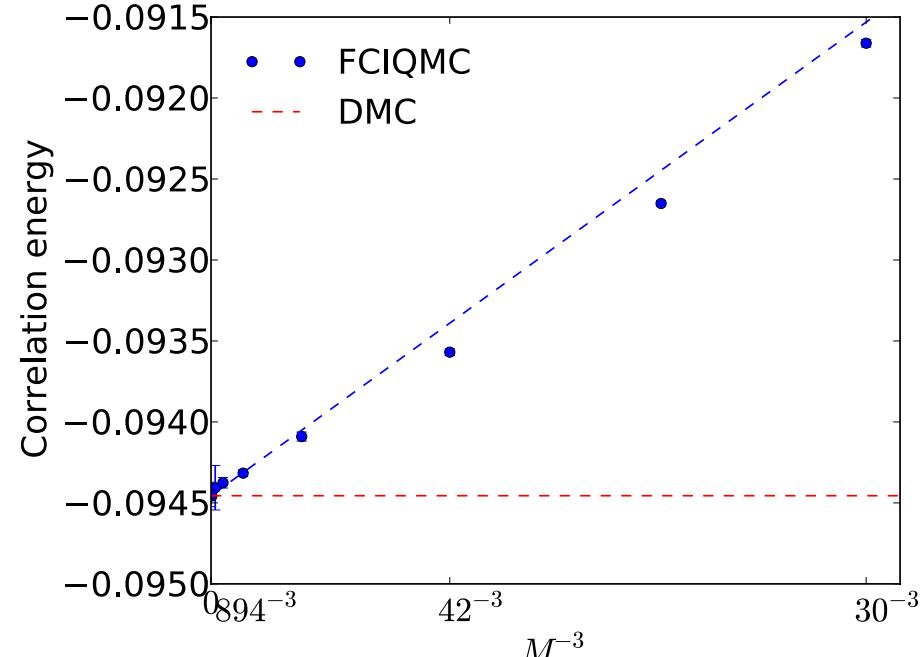
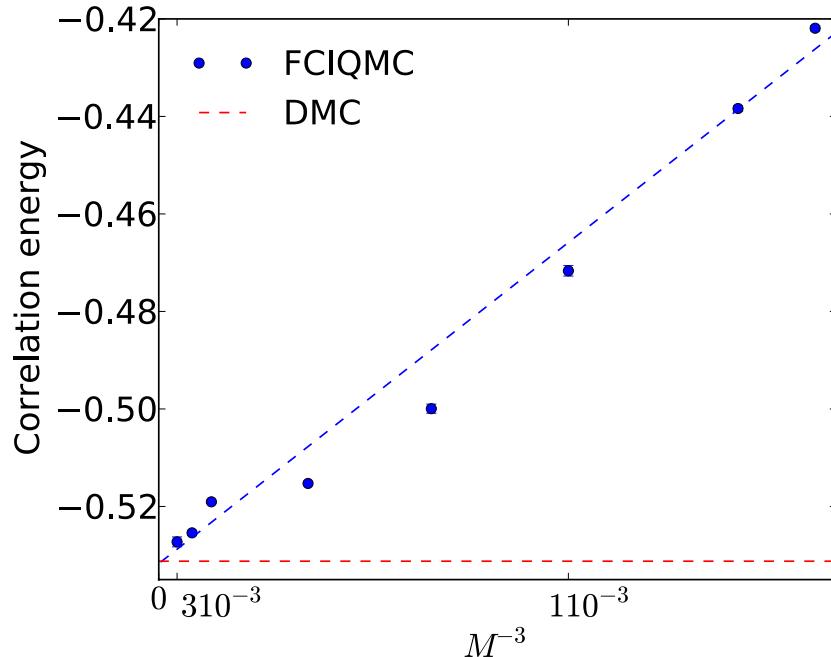


CBS limit is recovered as $1/M$...



Spin-polarization/dimension dependent...

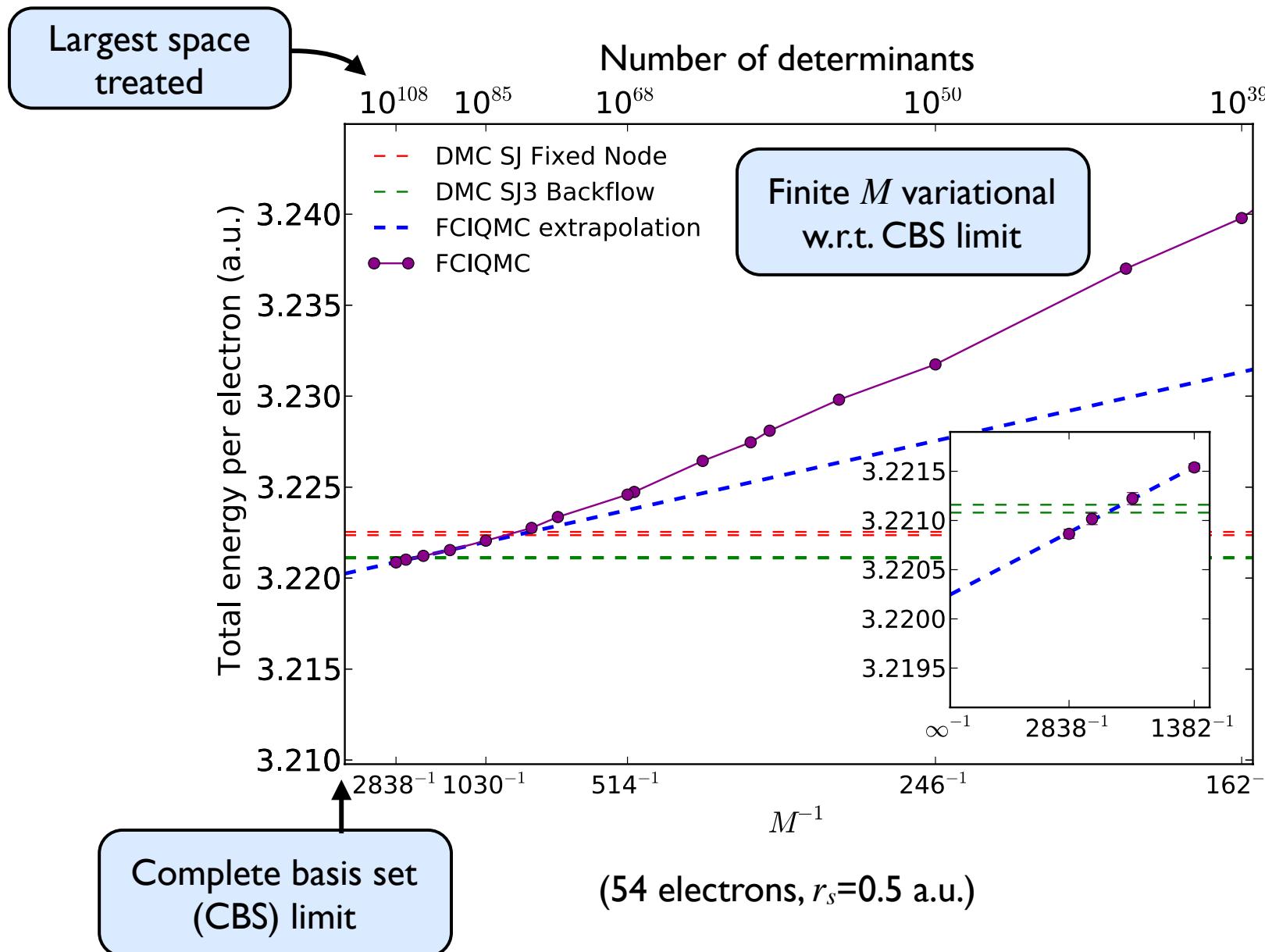
e.g. 1D spin-polarized, $1/M^3$



25 electrons, 1D
ferromagnetic

5 electrons, 1D
ferromagnetic

CBS limit is recovered as $1/M$...



$$E_{\text{corr}} = \sum_{j \in \{\text{doubles}\}} \langle D_j | H | D_0 \rangle \frac{c_j}{c_0}$$

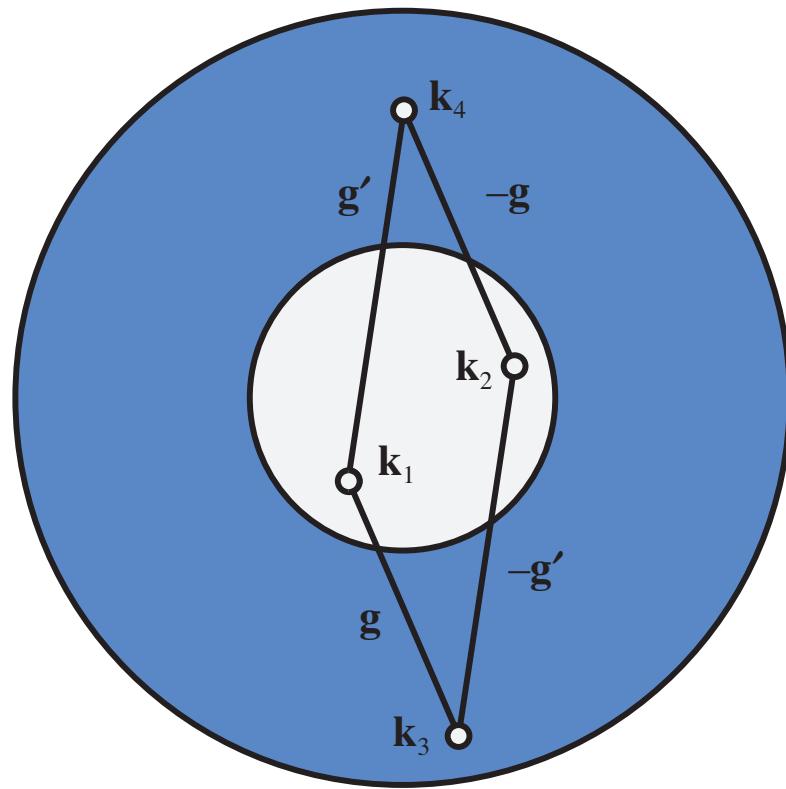


$$E_{\text{corr}} = \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \delta_{\mathbf{k}_i - \mathbf{k}_a, \mathbf{k}_j - \mathbf{k}_b} \left(v_{\mathbf{k}_i - \mathbf{k}_a} - v_{\mathbf{k}_j - \mathbf{k}_a} \right) \frac{c_{\mathbf{k}_i \mathbf{k}_j}^{\mathbf{k}_a \mathbf{k}_b}}{c_0}$$

Pairwise addition of
virtual contributions



Consider a different definition of basis set



- It is possible to define two momentum transfer vectors for an excitation.
- Show there are *three* sensible choices for basis sets based on this.
- Originates from basis set convergence work in solid state systems.

Consider a different definition of basis set

$$E_{\text{corr}} = \sum_{j \in \{\text{doubles}\}} \langle D_j | H | D_0 \rangle \frac{c_j}{c_0}$$



$$E_{\text{corr}} = \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \delta_{\mathbf{k}_i - \mathbf{k}_a, \mathbf{k}_j - \mathbf{k}_b} \left(v_{\mathbf{k}_i - \mathbf{k}_a} - v_{\mathbf{k}_j - \mathbf{k}_a} \right) \frac{c_{\mathbf{k}_a \mathbf{k}_b}}{c_0}$$



Pairwise addition of
virtual contributions

Assume fixed

Consider a different definition of basis set

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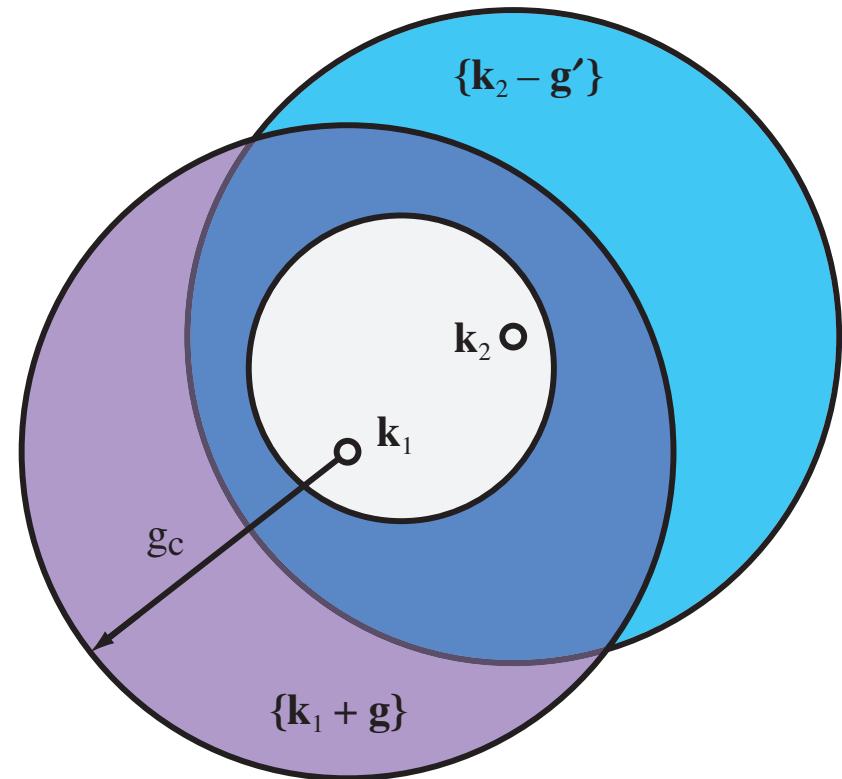
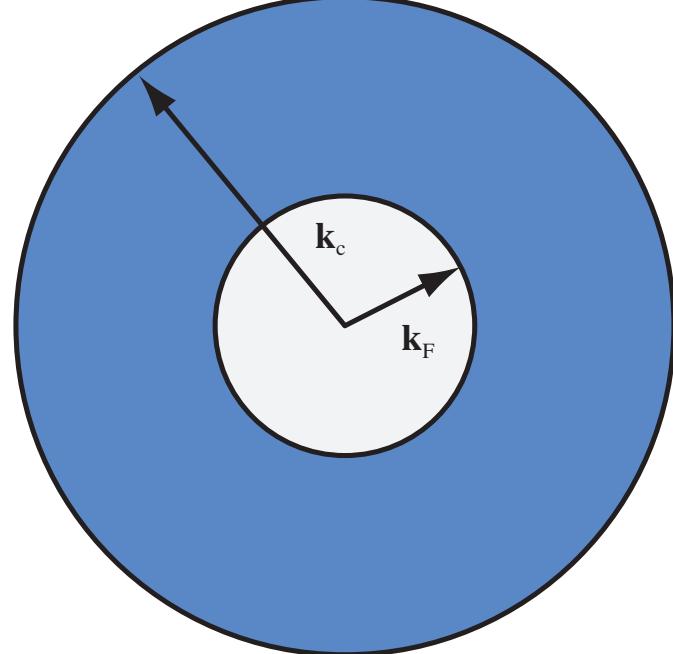


$$E_{\text{corr}} = \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \delta_{\mathbf{k}_i - \mathbf{k}_a, \mathbf{k}_j - \mathbf{k}_b} \left(v_{\mathbf{k}_i - \mathbf{k}_a} - v_{\mathbf{k}_j - \mathbf{k}_a} \right) \frac{c_{\mathbf{k}_i \mathbf{k}_j}^{\mathbf{k}_a \mathbf{k}_b}}{c_0}$$

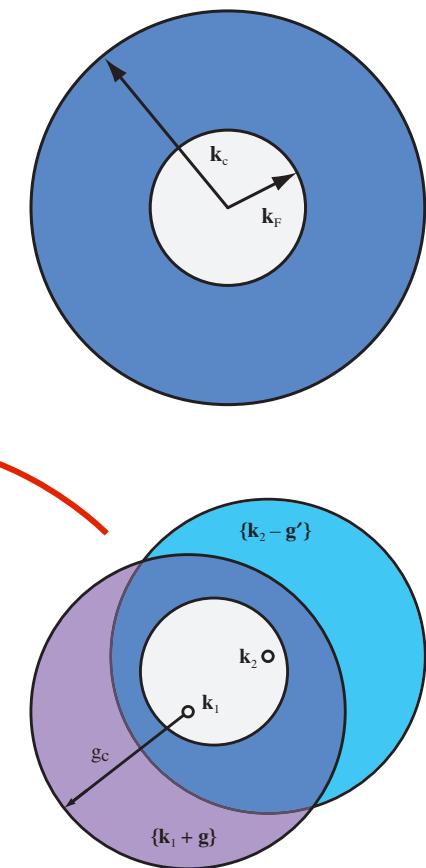
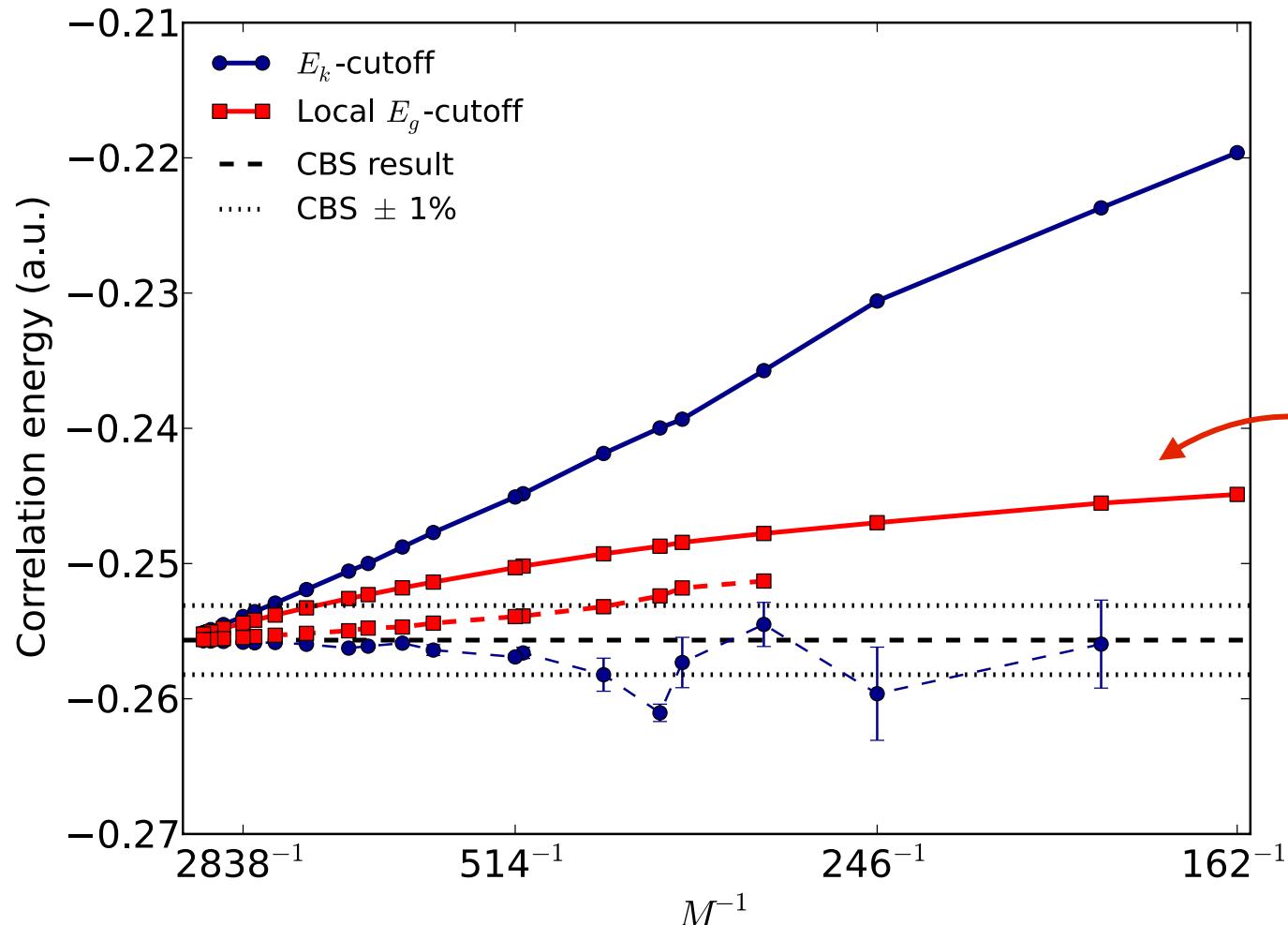


$$\sum_{ab}^{k_c} F(k) \longrightarrow \sum_{ab}^{g_c} F(g)$$

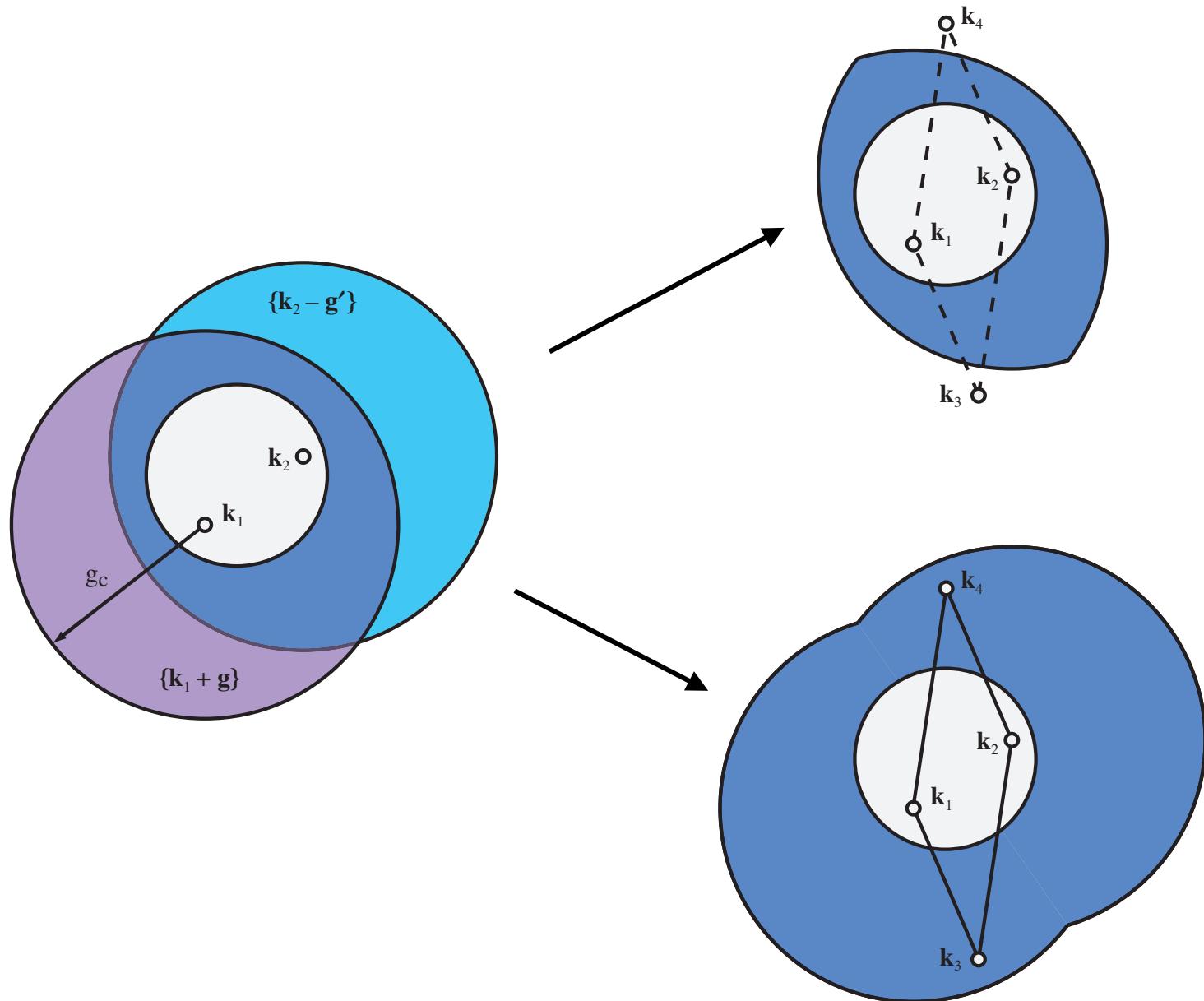
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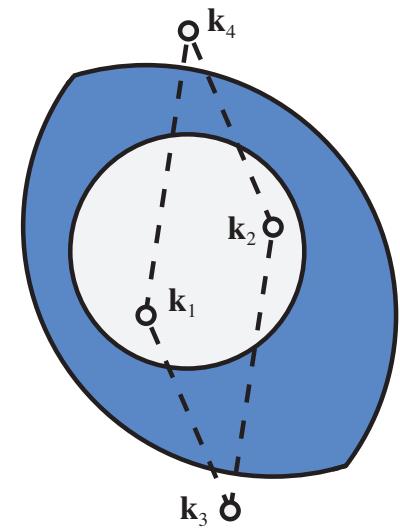
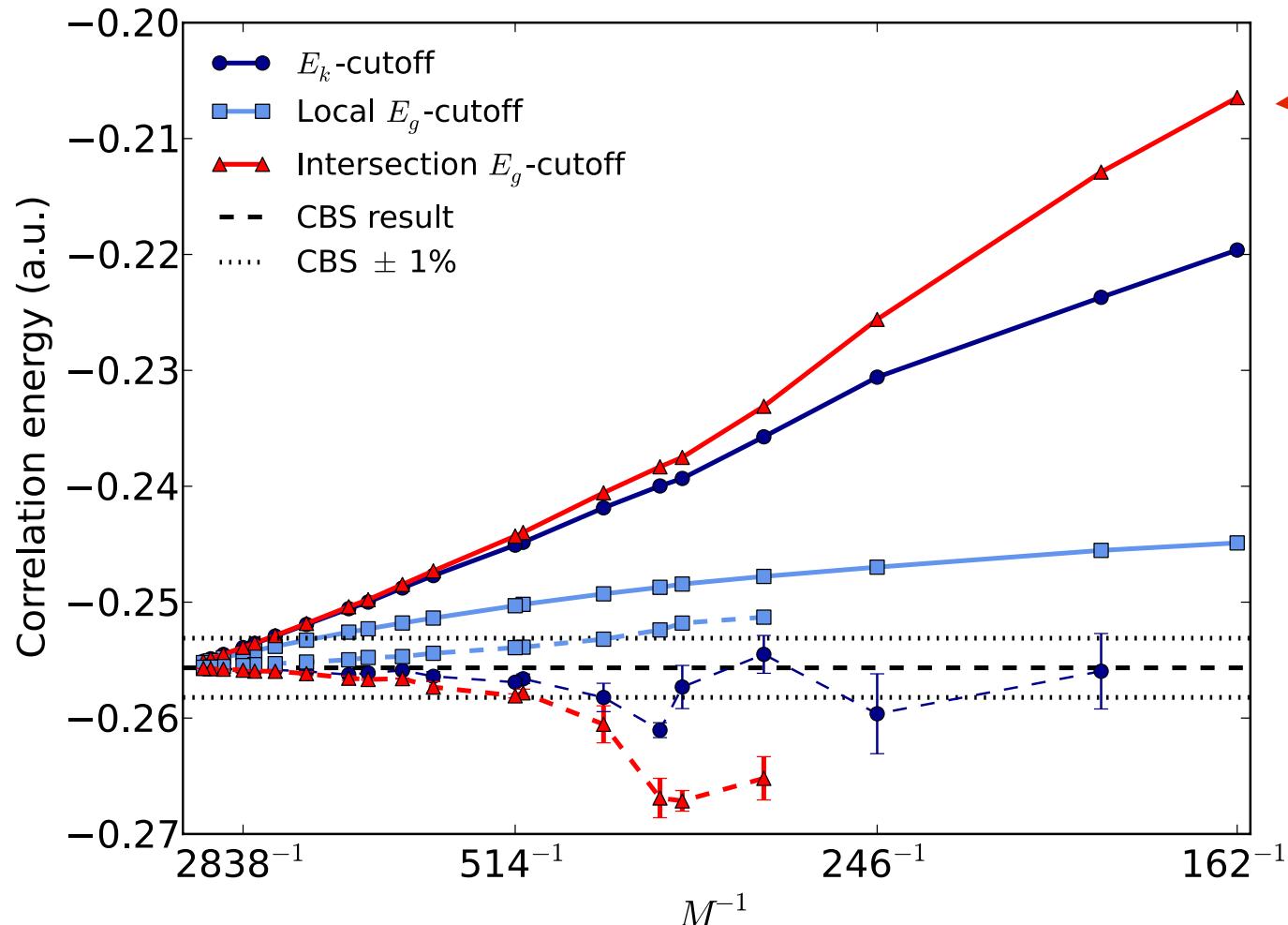
HEG/extrapolation: JJS, A Gruneis, GH Booth, G Kresse, A Alavi, PRB, 86, 035111 (2012).



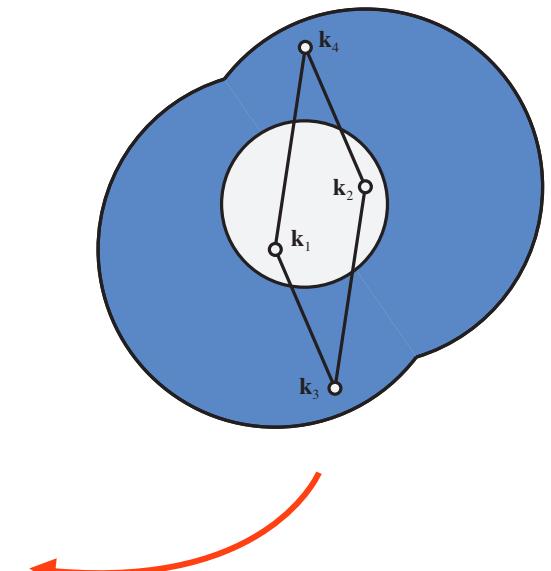
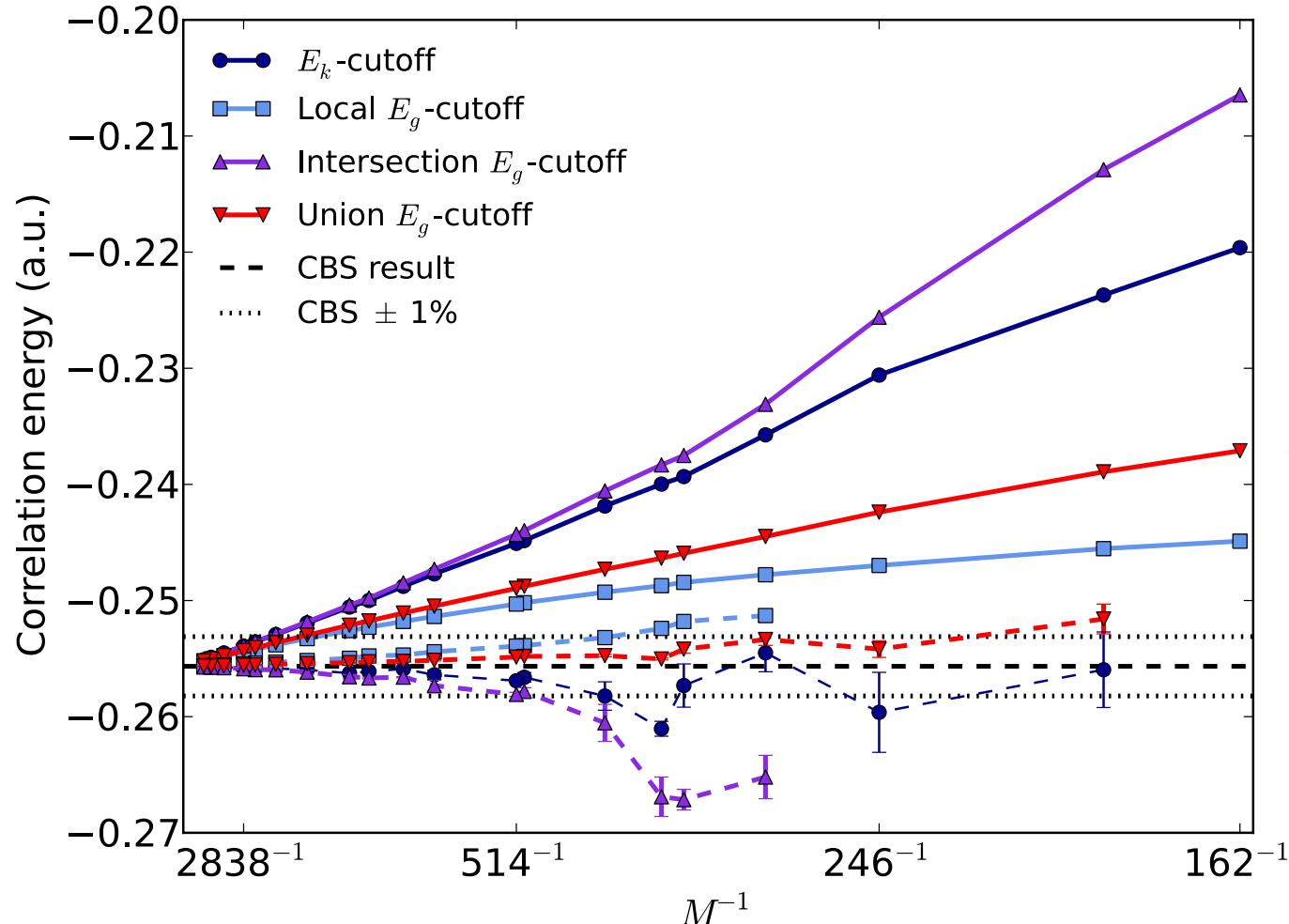
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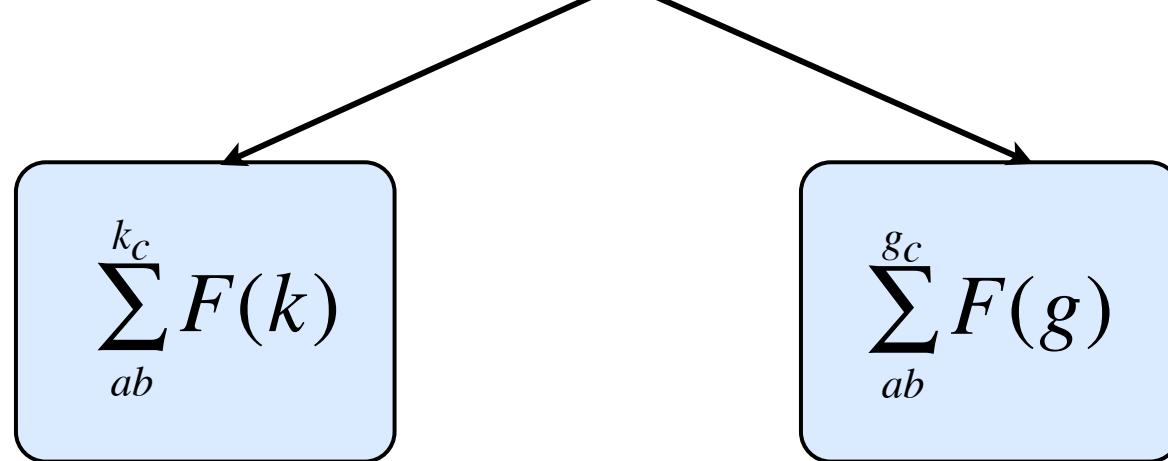
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For multi-reference calculations, there is a trade-off...

$$E_{\text{corr}} = \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \delta_{\mathbf{k}_i - \mathbf{k}_a, \mathbf{k}_j - \mathbf{k}_b} \left(v_{\mathbf{k}_i - \mathbf{k}_a} - v_{\mathbf{k}_j - \mathbf{k}_a} \right) \frac{c_{\mathbf{k}_i \mathbf{k}_j}^{\mathbf{k}_a \mathbf{k}_b}}{c_0}$$



Slower/less accurate
extrapolation

Neglects coefficient
relaxation

r_s	M	k extrapolation	g extrapolation
0.5	1850	-0.5969(3)	-0.5959(7)
1.0	1850	-0.5325(4)	-0.5316(4)
2.0	1850	-0.4447(4)	-0.444(1)
5.0	778	-0.306(1)	-0.307(1)

For multi-reference calculations, there is a trade-off...

$$E_{\text{corr}} = \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \delta_{\mathbf{k}_i - \mathbf{k}_a, \mathbf{k}_j - \mathbf{k}_b} \left(v_{\mathbf{k}_i - \mathbf{k}_a} - v_{\mathbf{k}_j - \mathbf{k}_a} \right) \frac{c_{\mathbf{k}_i \mathbf{k}_j}^{\mathbf{k}_a \mathbf{k}_b}}{c_0}$$

- Single (multi-reference) calculation
- Immediate CBS estimate
- Systematically improvable

$$\sum_{ab}^{g_c} F(g)$$

Overview

- Introduction
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How good are the HF nodes?

- Single-determinant Slater-Jastrow wavefunction
- Real-space propagation using a Green's function propagator
- Makes the Schrödinger equation diffusive
- Must enforce anti-symmetry explicitly:

$$f(\mathbf{x}, \tau) = \Psi(\mathbf{x}, \tau) \phi_T(\mathbf{x})$$

- How good are the trial wavefunction's nodes?
(infer from energy)

i-FCIQMC, exact except for:

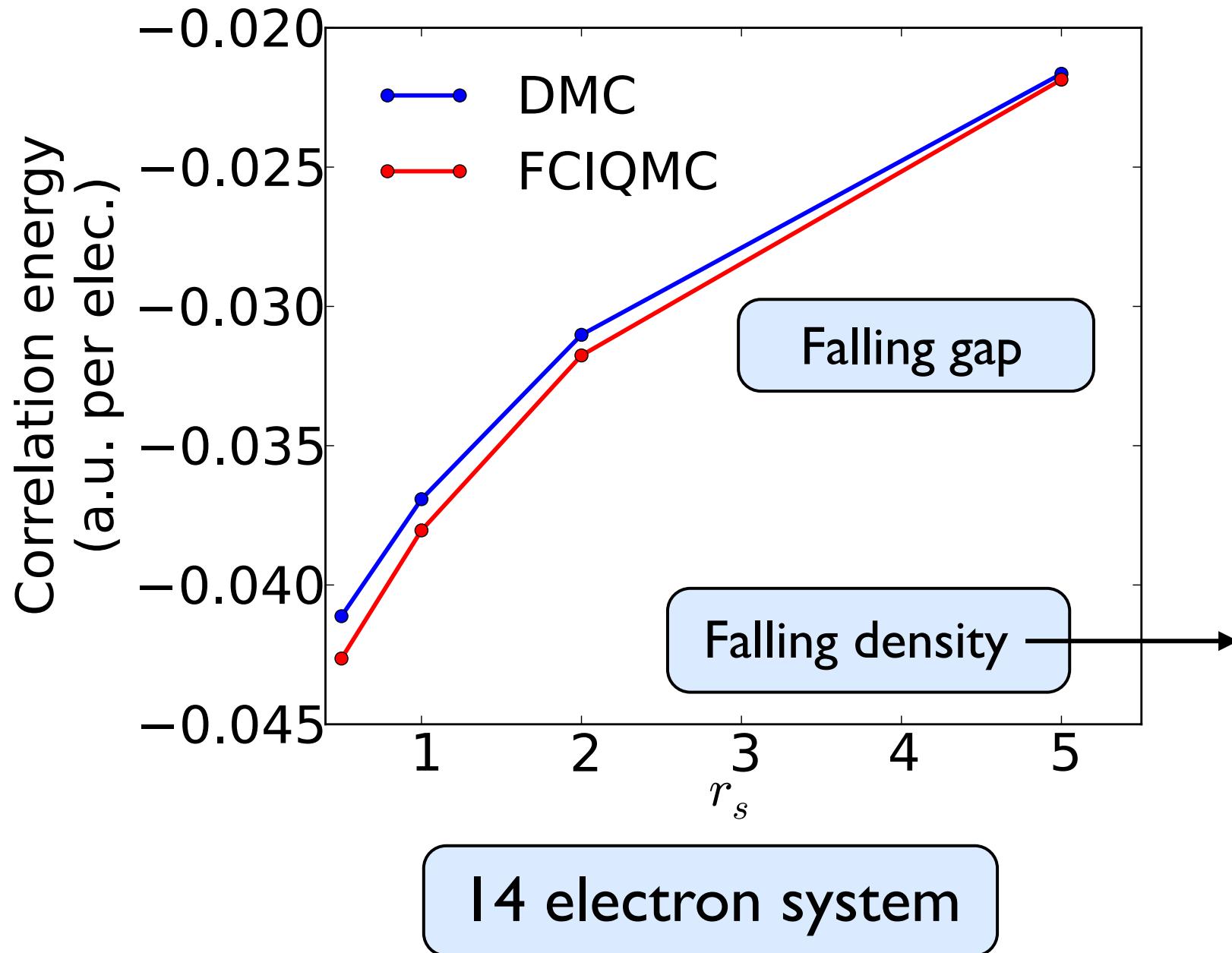
- Basis set incompleteness error
- Initiator error

DMC, exact except for:

- Fixed-node error (after backflow)
- (Time-step error)

CASINO: RJ Needs, MD Towler, ND Drummond and P Lopez Rios, J. Phys.: Condensed Matter 22, 023201 (2010).

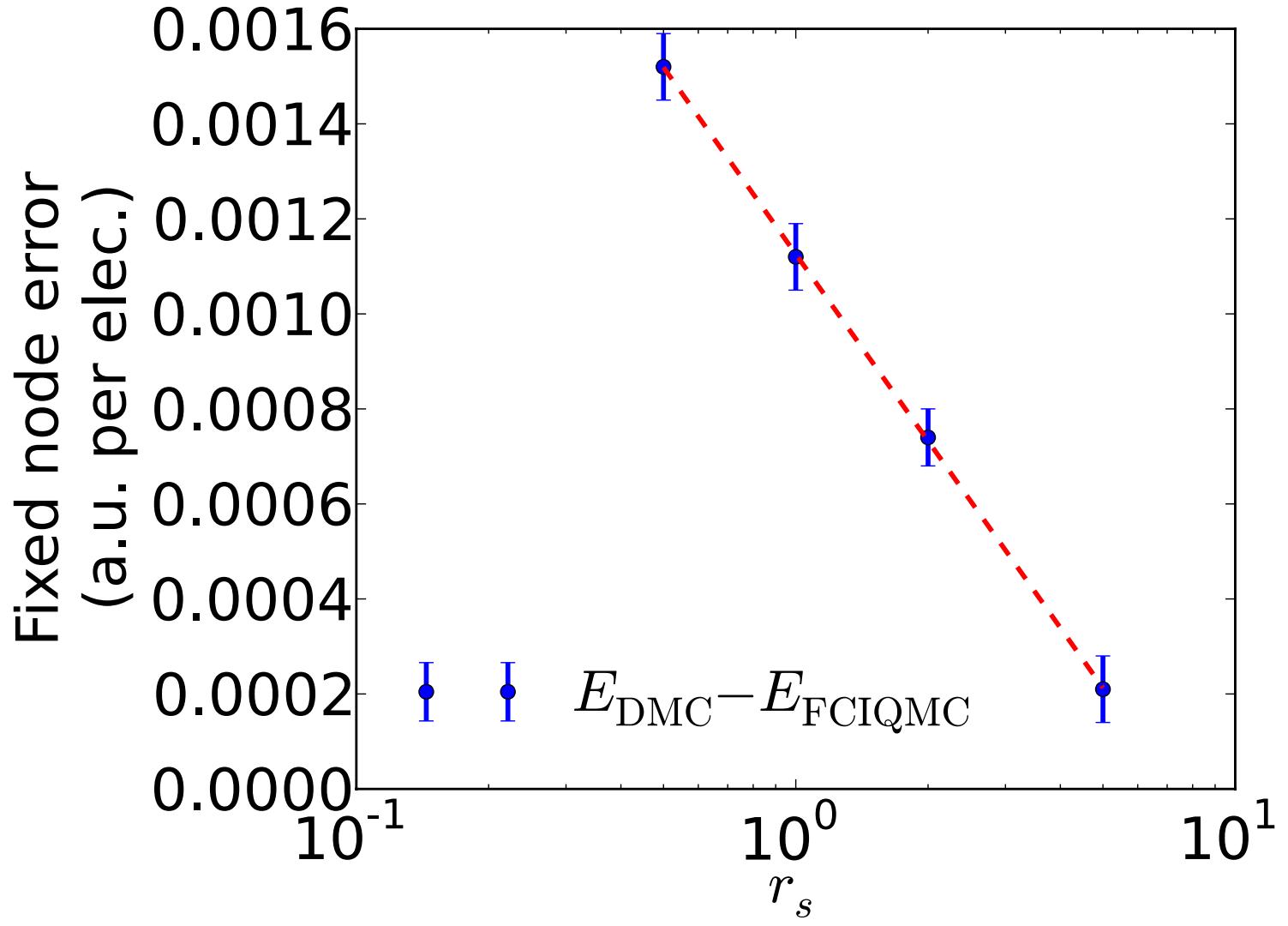
Backflow: P. Lopez Rios, A. Ma, ND Drummond, MD Towler and RJ Needs, Phys. Rev. E 74, 066701 (2006)



Consistent with:

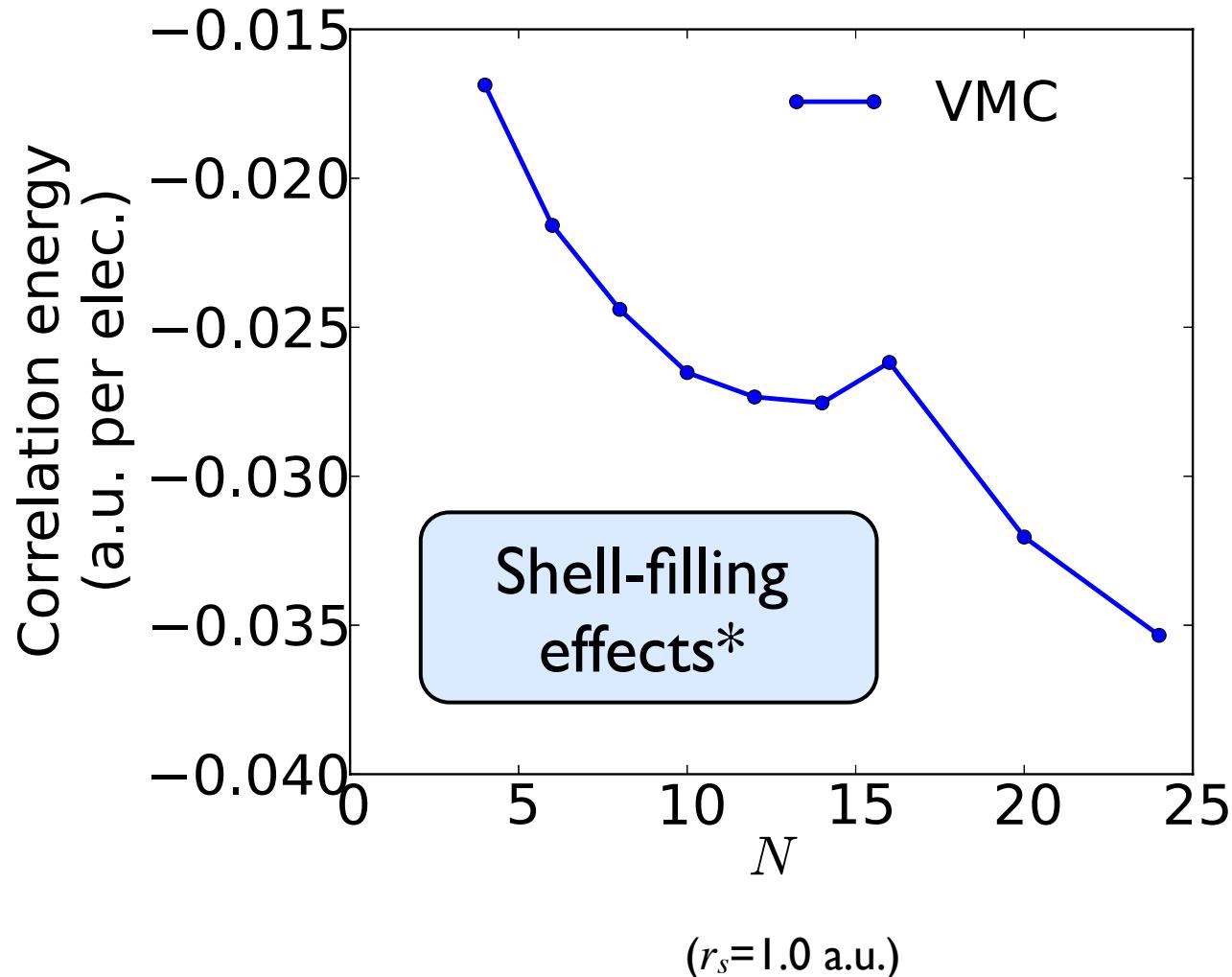
K.M. Rasch, L. Mitas, Chem. Phys. Lett. 528, 59 (2012)

Y Kwon, D. M. Ceperley and R. M. Martin, Phys. Rev. B 58, 6800 (1998).



14 electron system

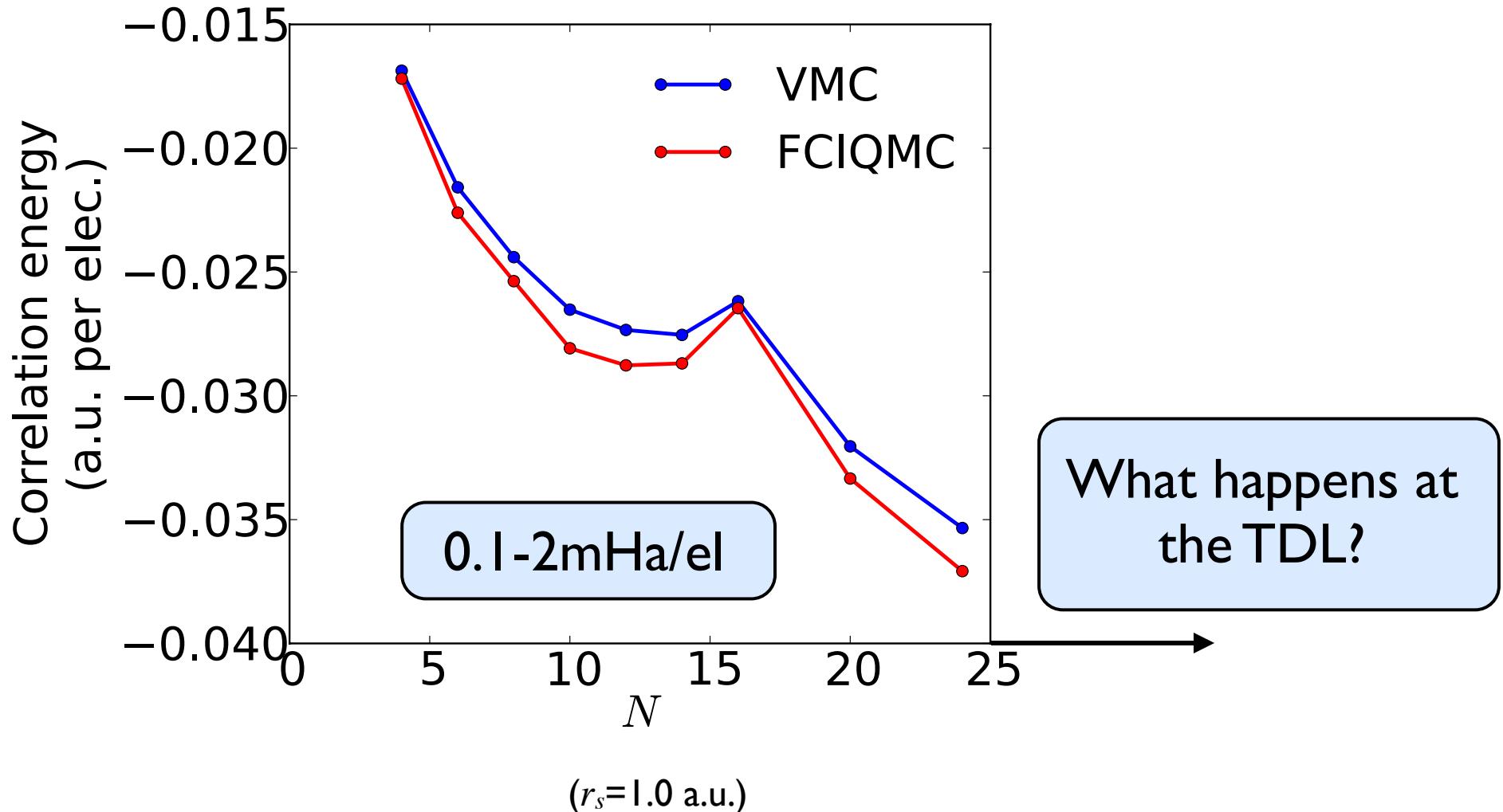
The correlation energy undulates with changing N ...



VMC results due to Pablo López-Ríos

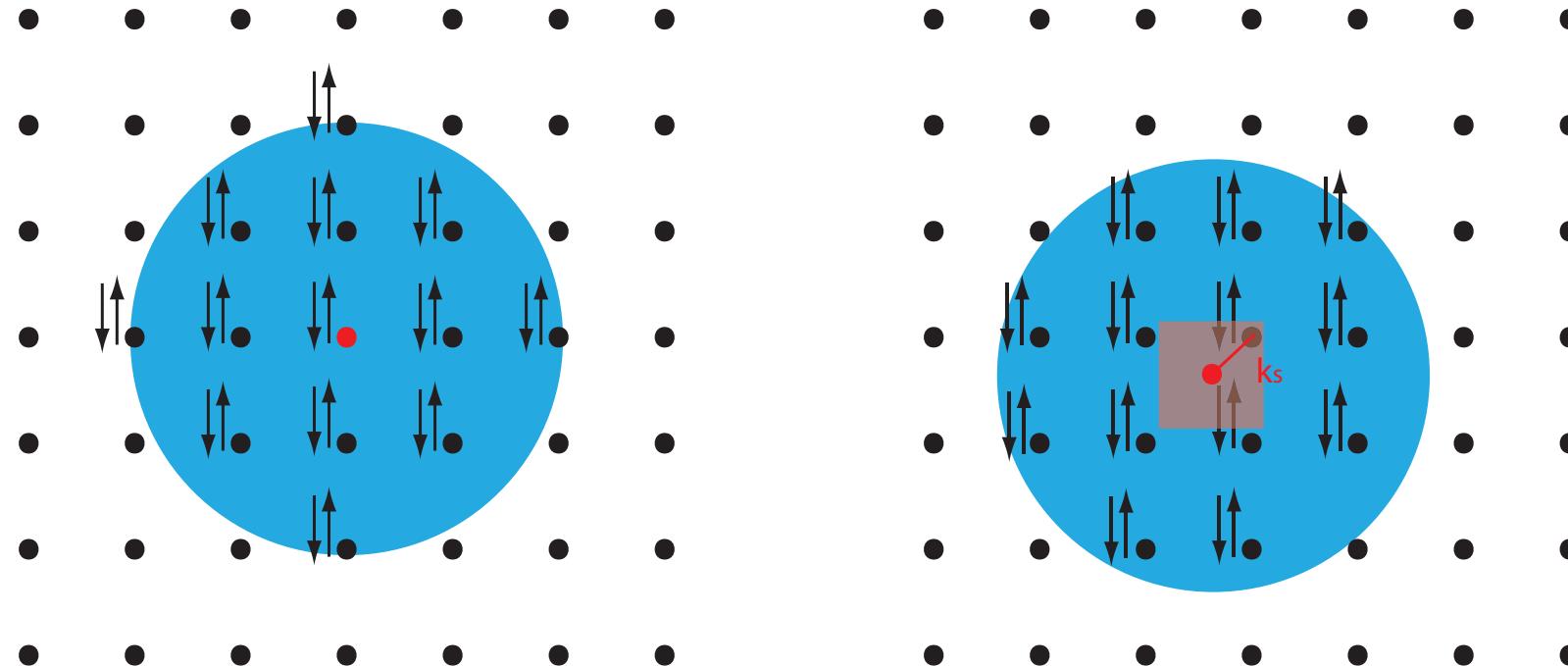
*Non-zero twist angle

...but so does the gap.



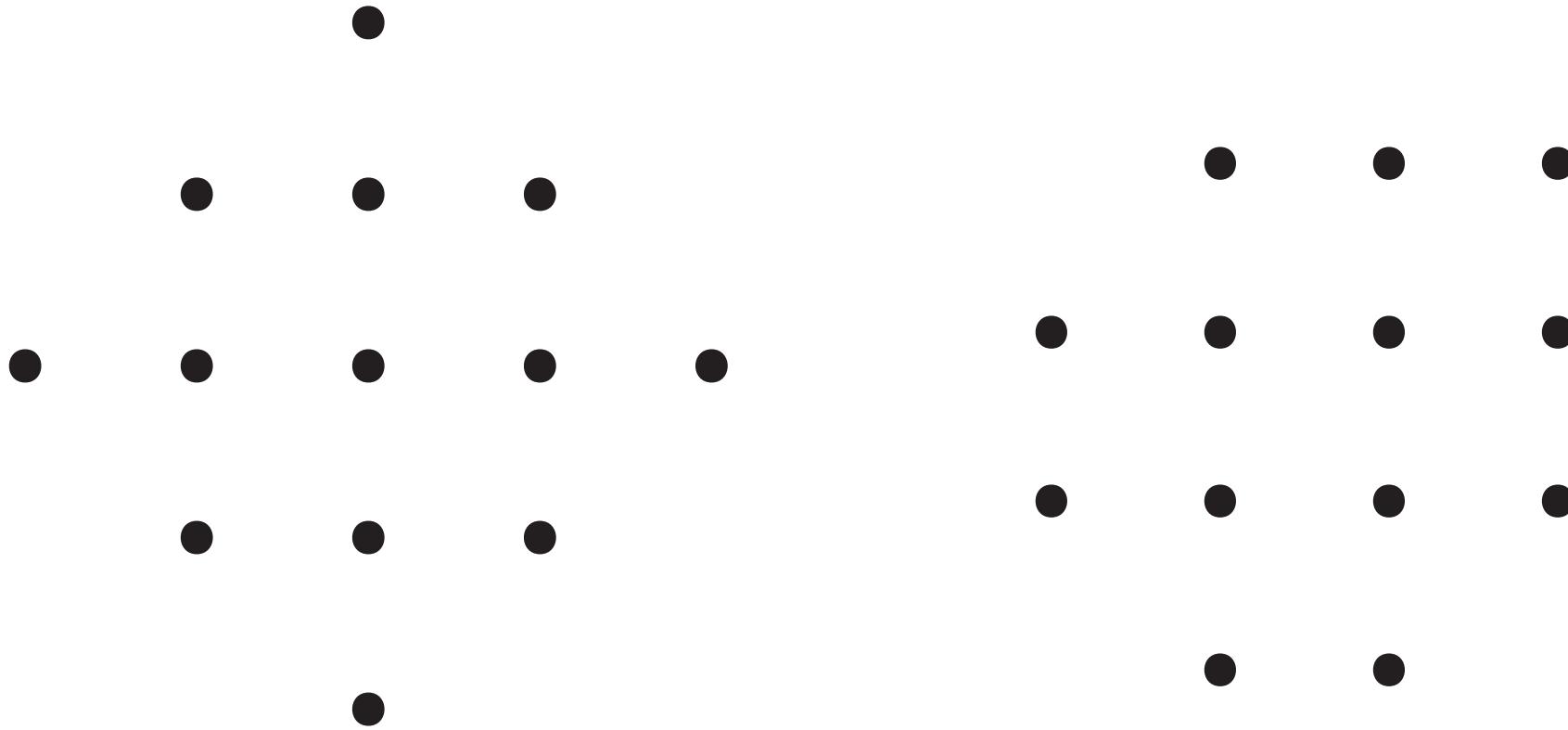
VMC results due to Pablo López-Ríos

Simulation cell HEG systems can resemble clusters of electrons in k -space more than the continuum picture of a sphere



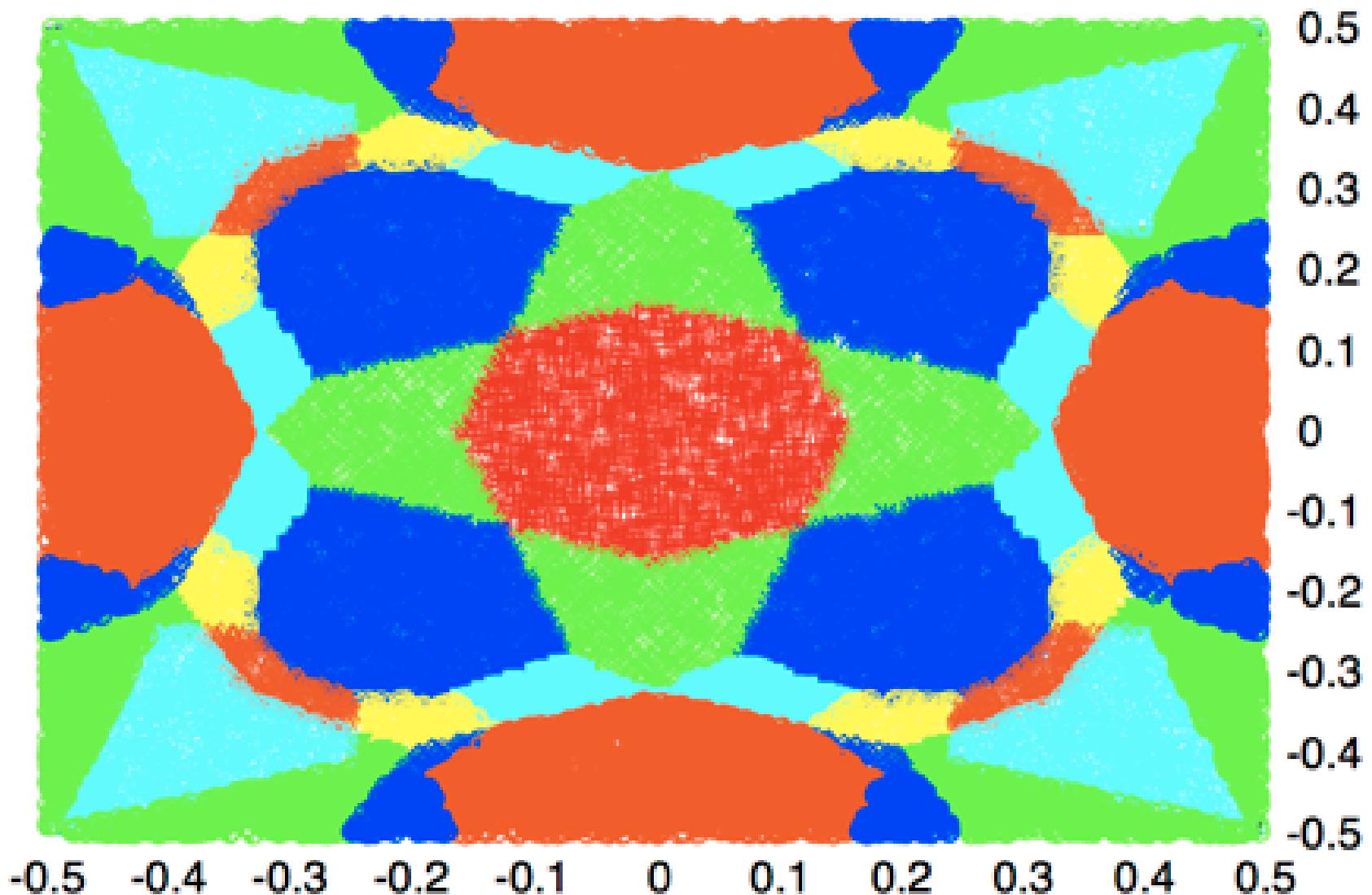
Offset can cause ‘repopulation’

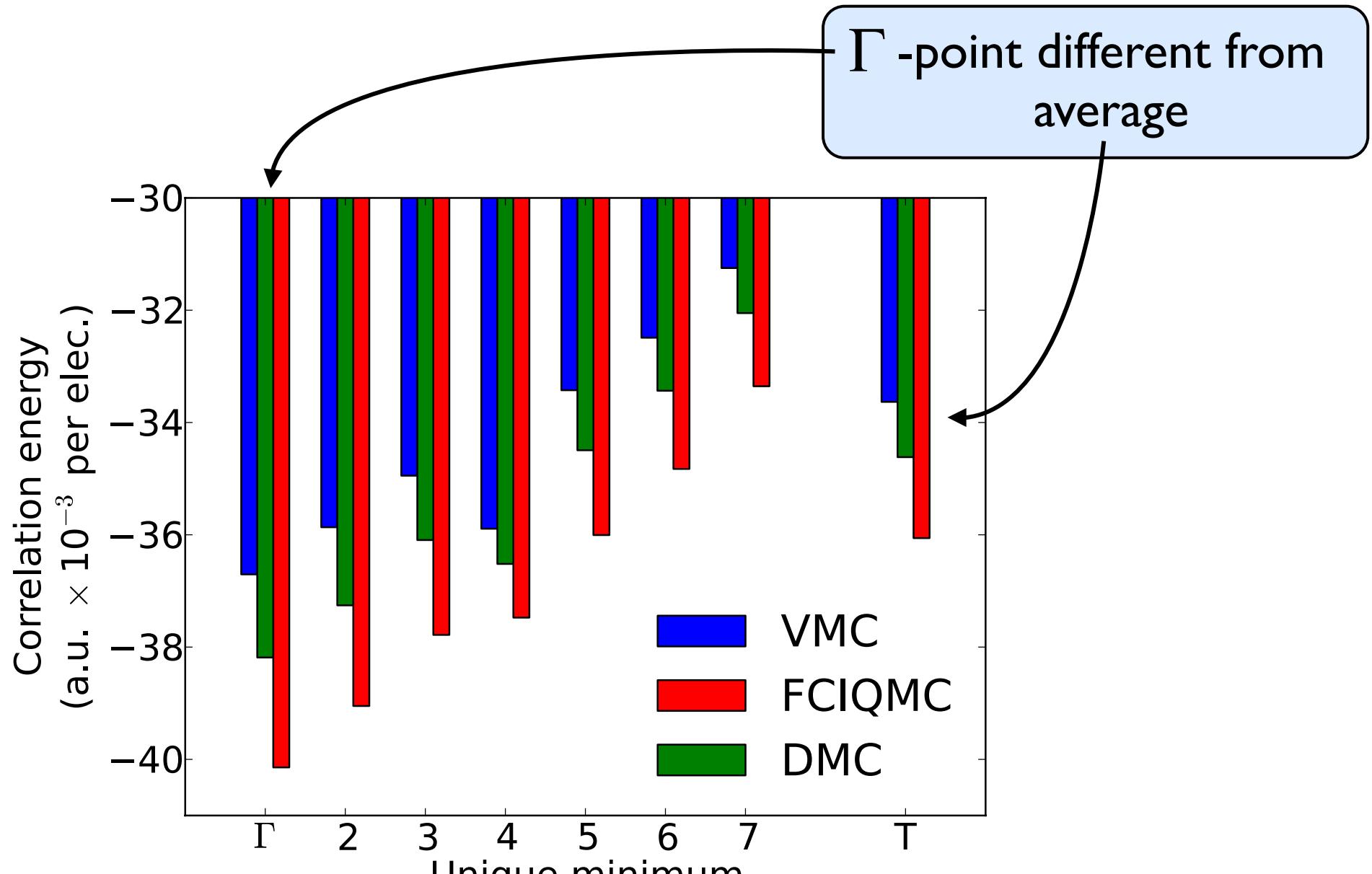
C Lin, F-H Zong, D M Ceperley, Phys Rev E 64, 016702 (2001)
N Drummond, R Needs, A Sorouri, W Foulkes, Phys Rev B, 78, 125106 (2008)



These now have different correlation energies...
(HF changes regardless)

54 electron system

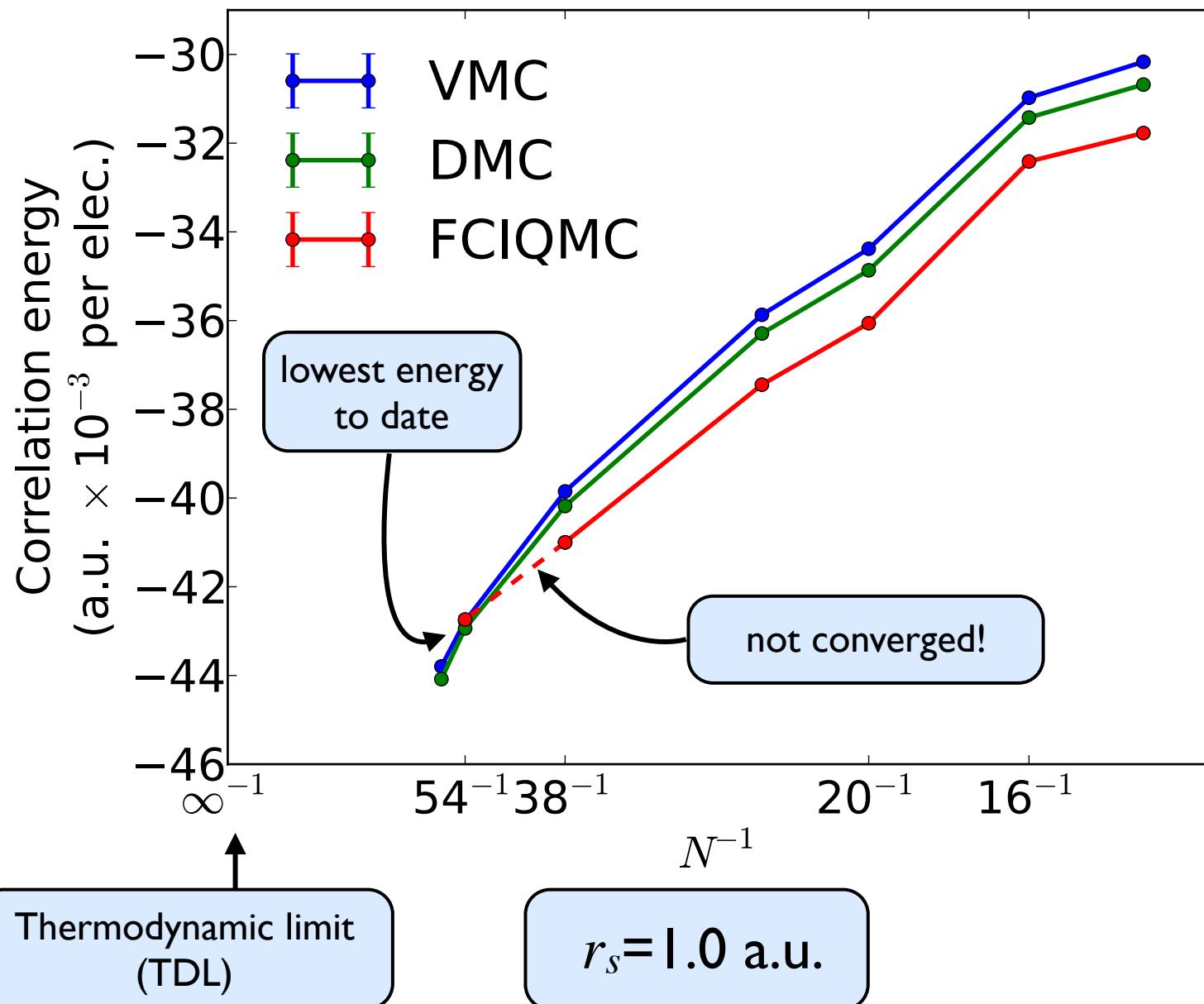




DMC/VMC results due to Pablo López-Ríos

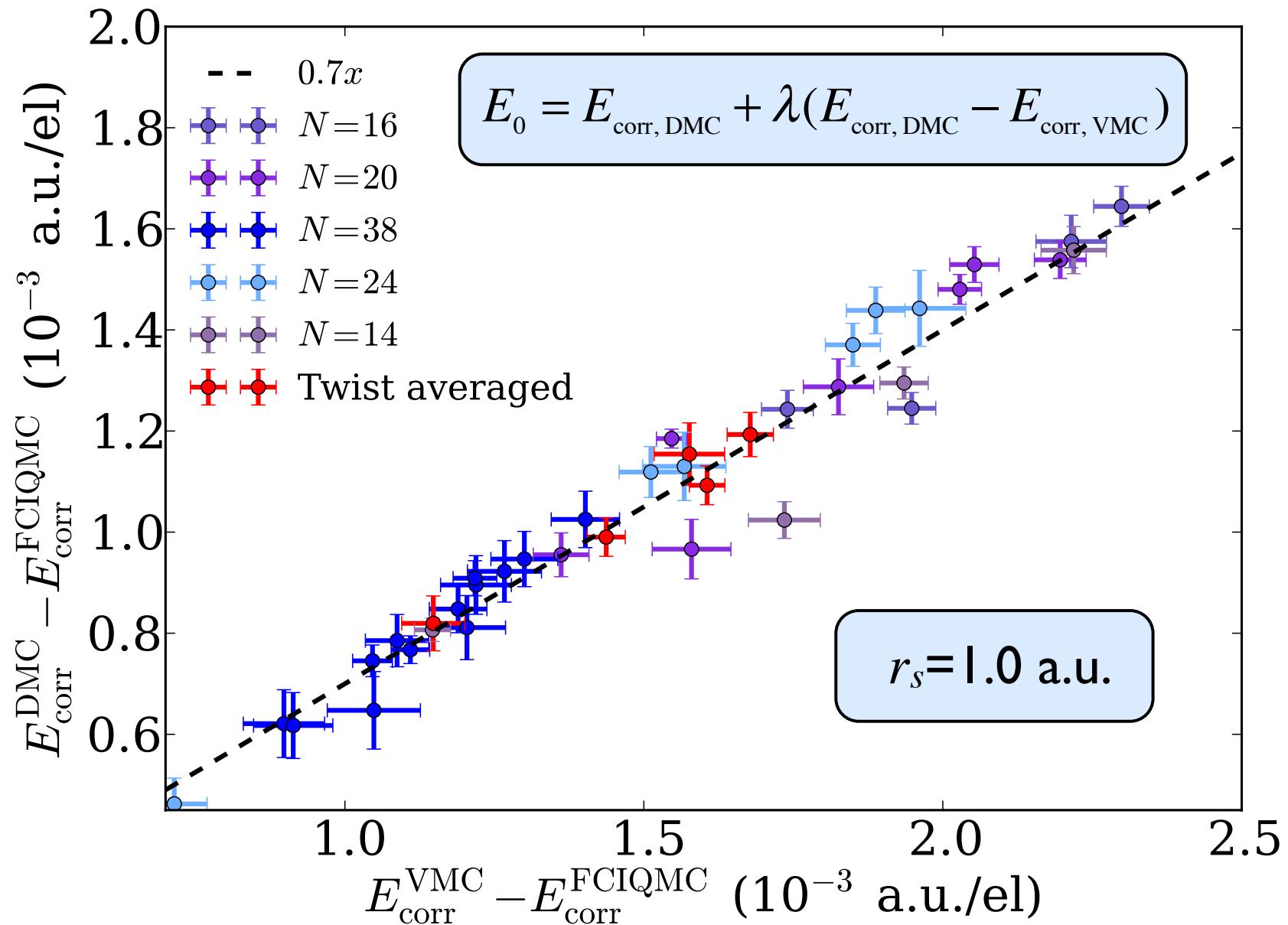
20 electron system

Twist-averaging yields a smooth pattern with $N...$



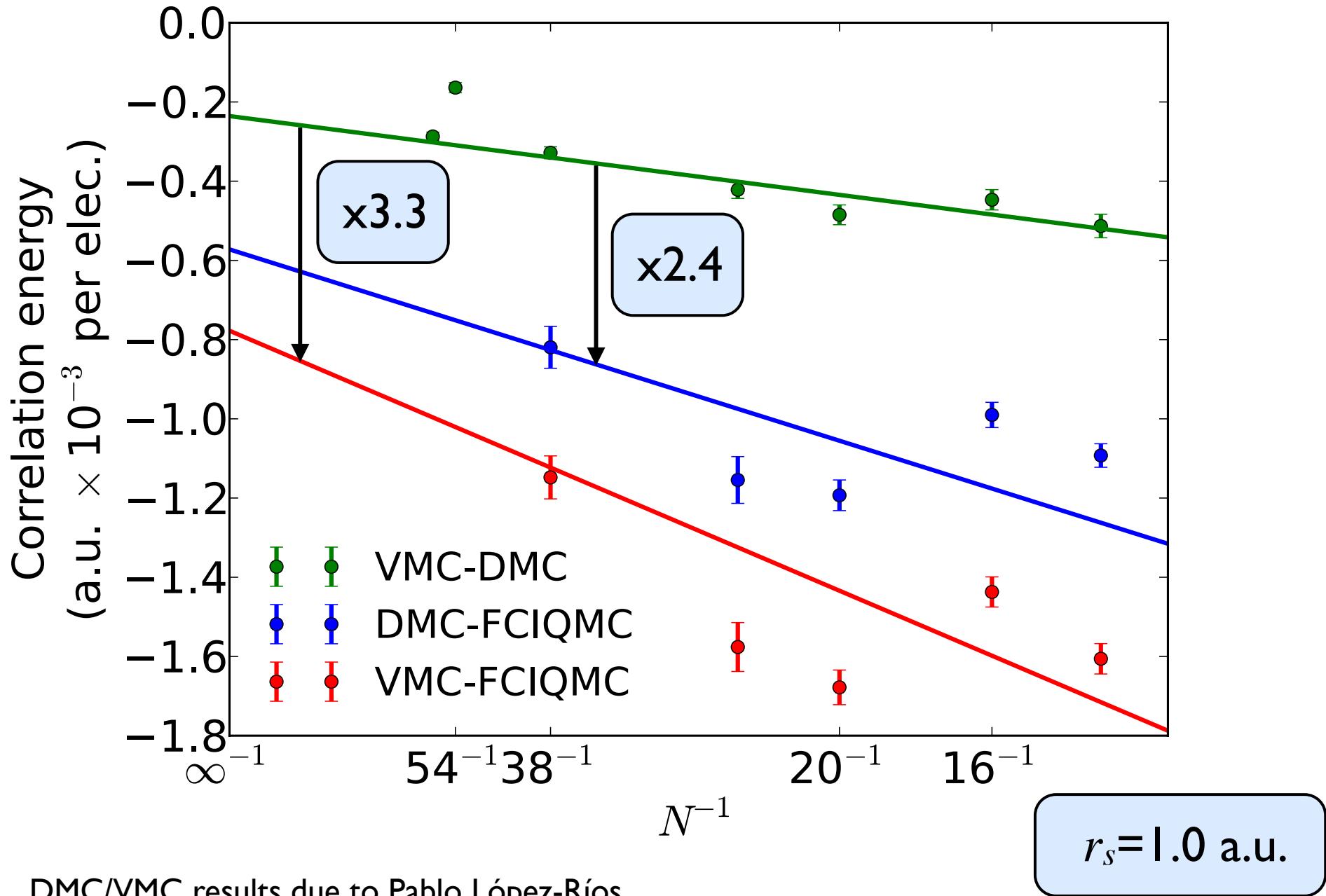
DMC/VMC results due to Pablo López-Ríos

With an underlying relationship

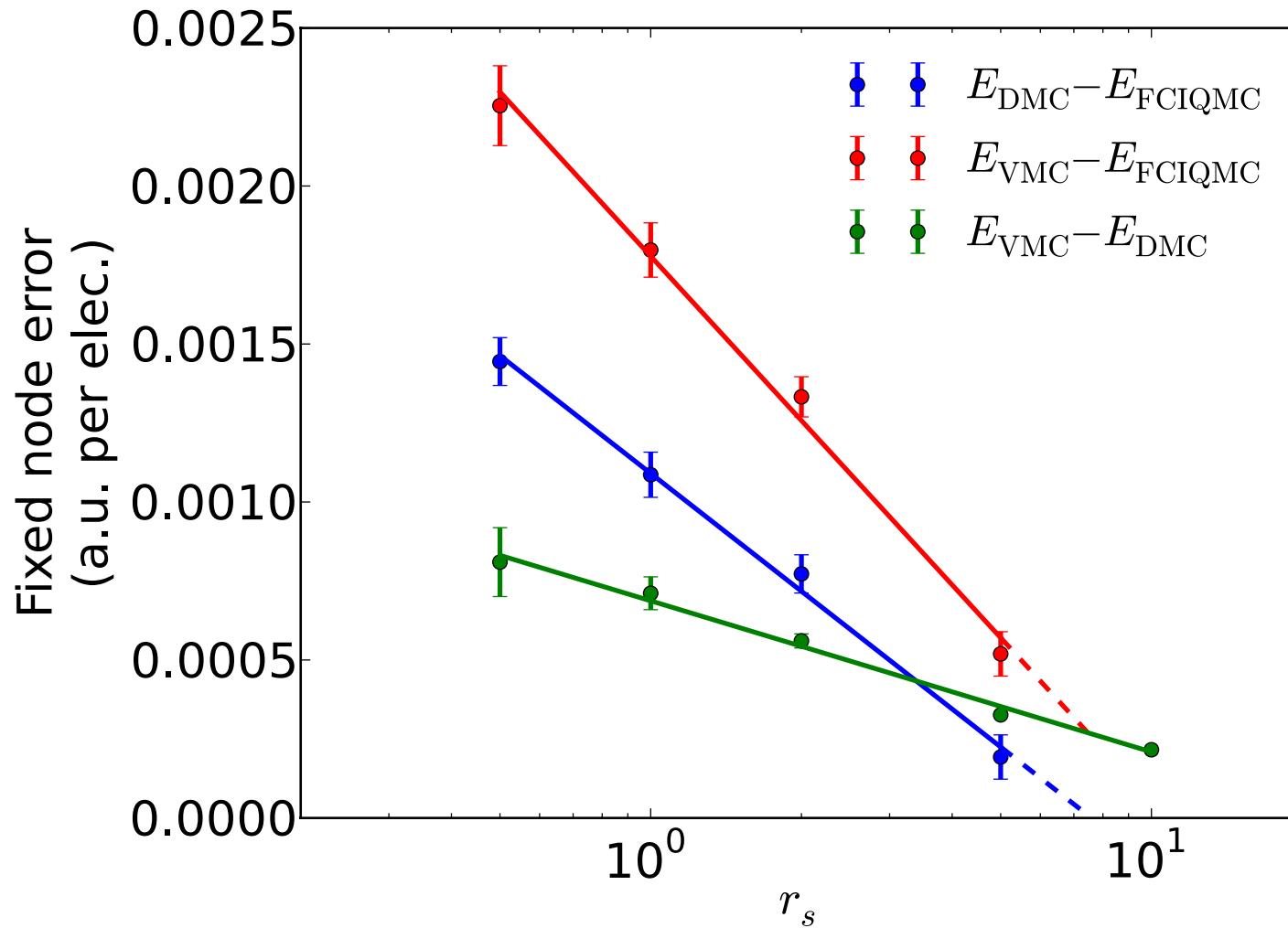


DMC/VMC results due to Pablo López-Ríos

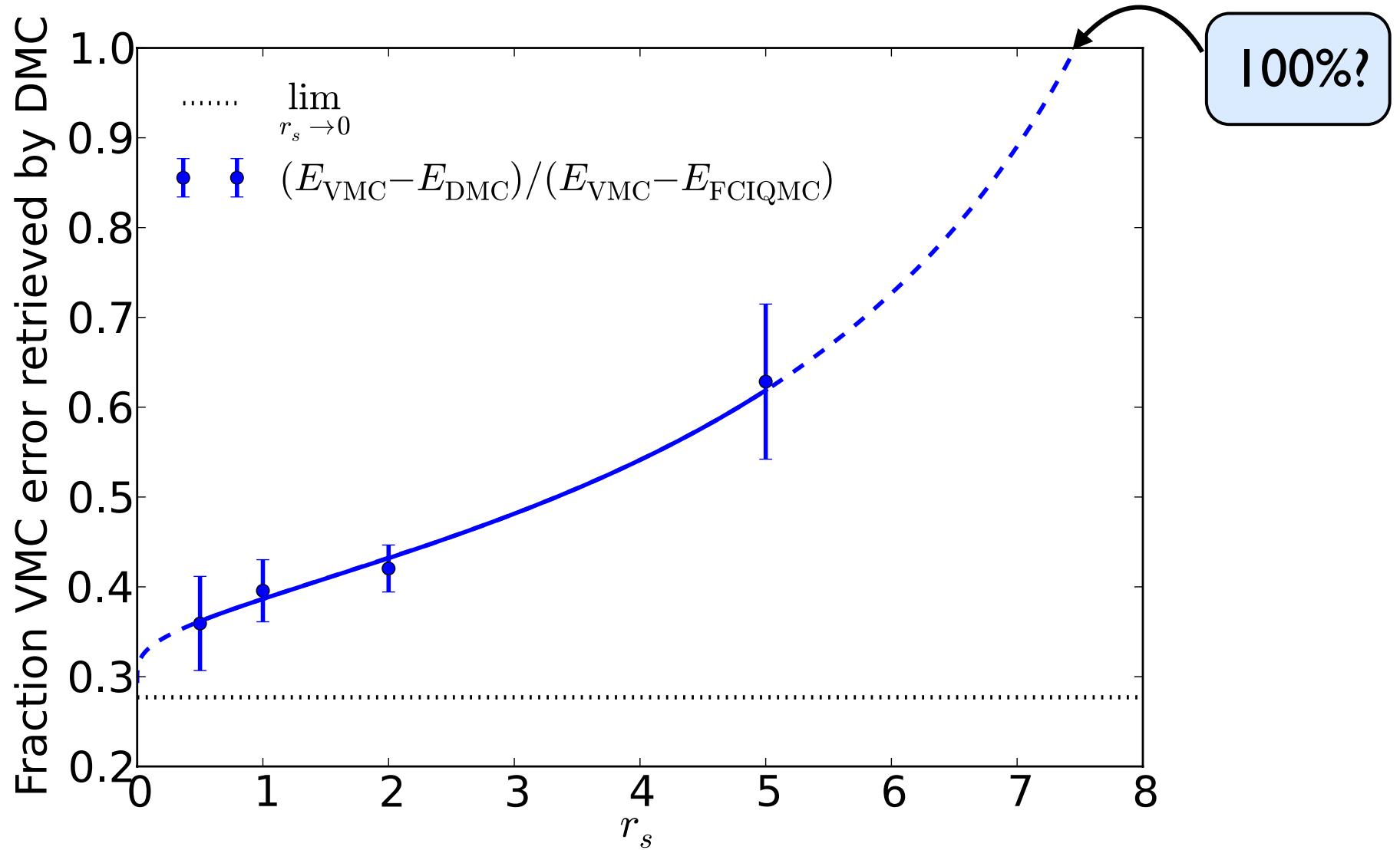
Using this relationship...



Ratios are hopefully fairly constant with N ,
so return to 14 electrons...



How well does DMC do in retrieving the missing correlation energy in VMC?



Overview

- Introduction
- Plane wave basis set incompleteness error
- Comparison with DMC and finite size effects
- Insights from the FCIQMC wavefunction

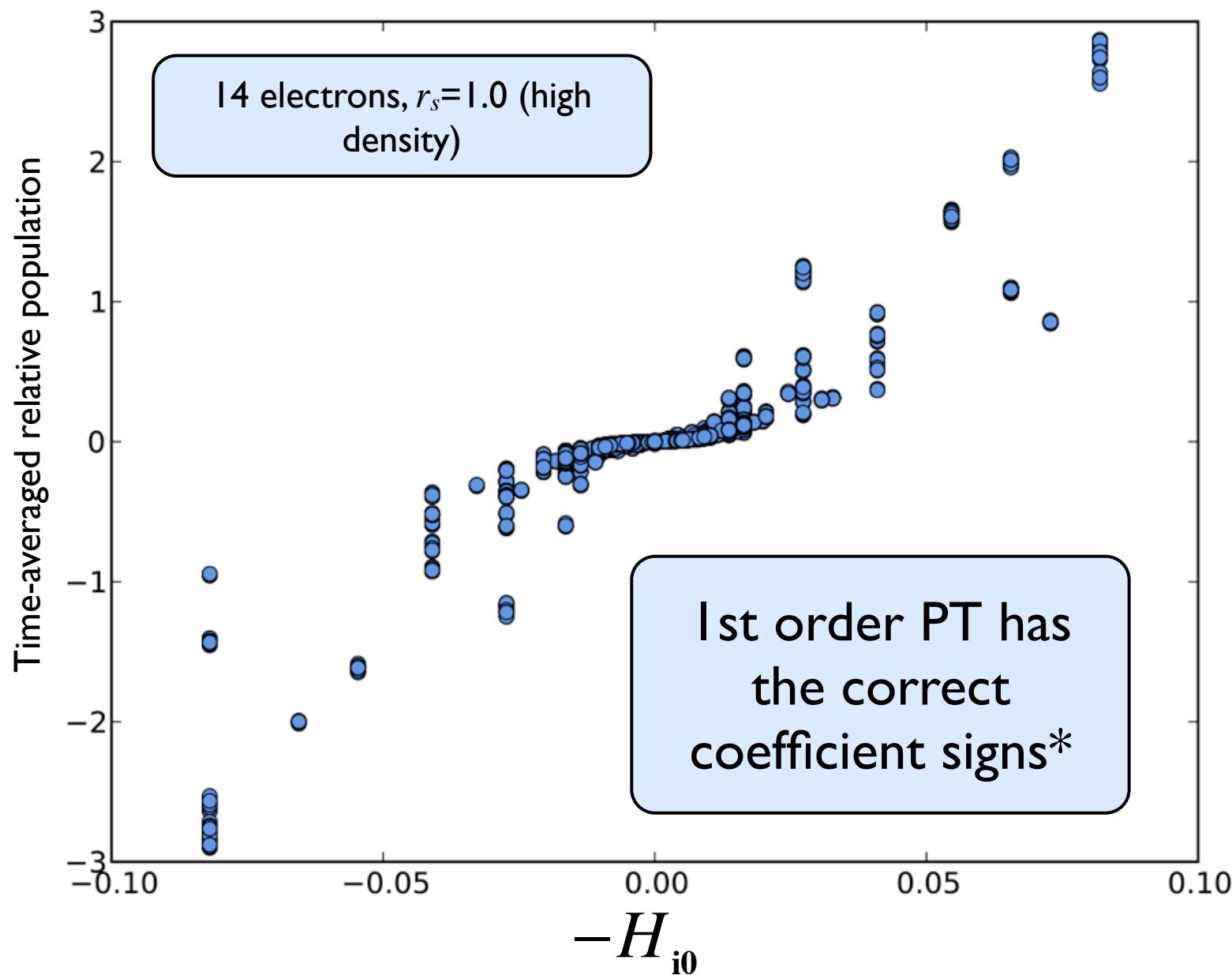
What about the wavefunction?

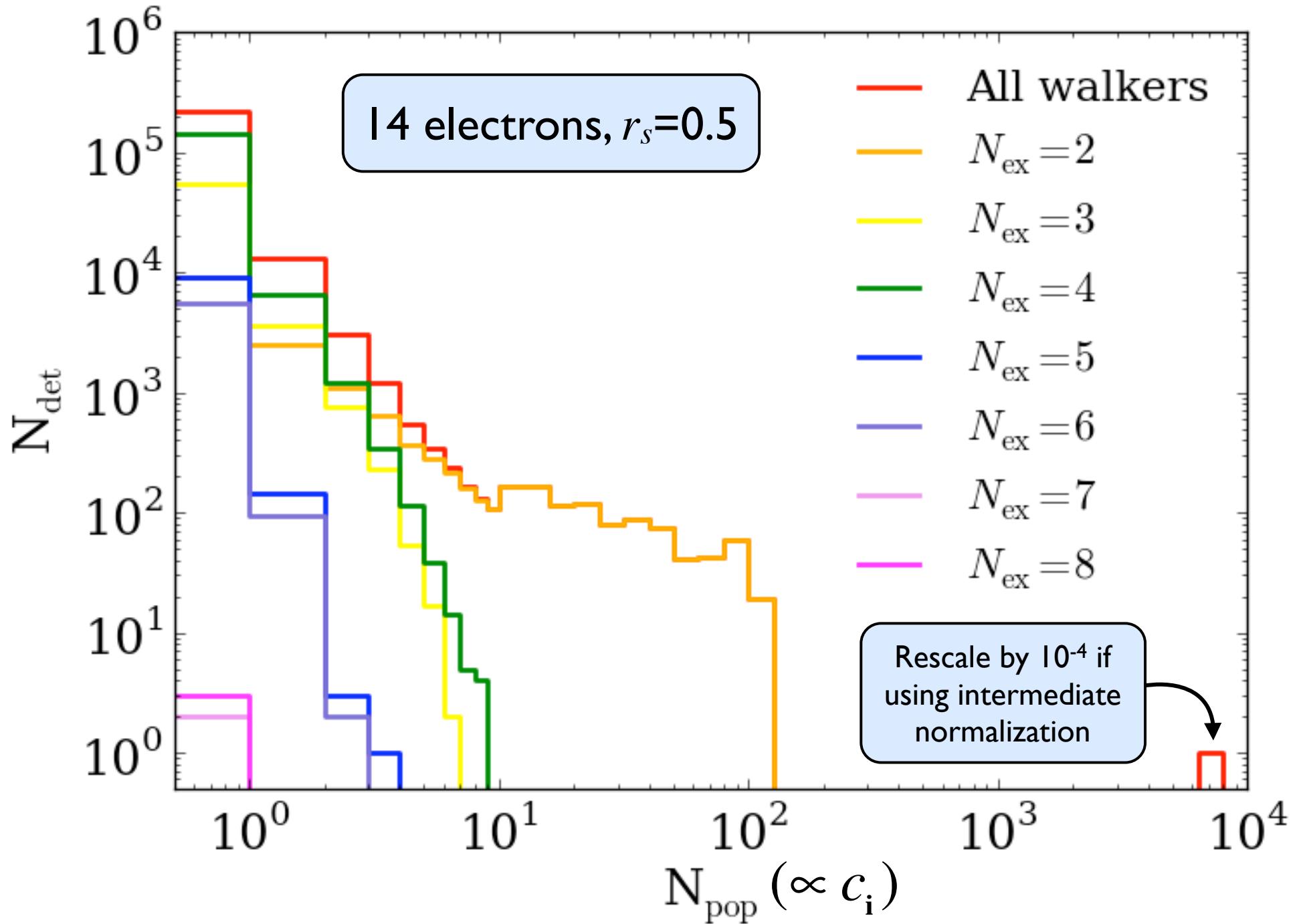
i-FCIQMC directly simulates a representation of the wavefunction (walkers)

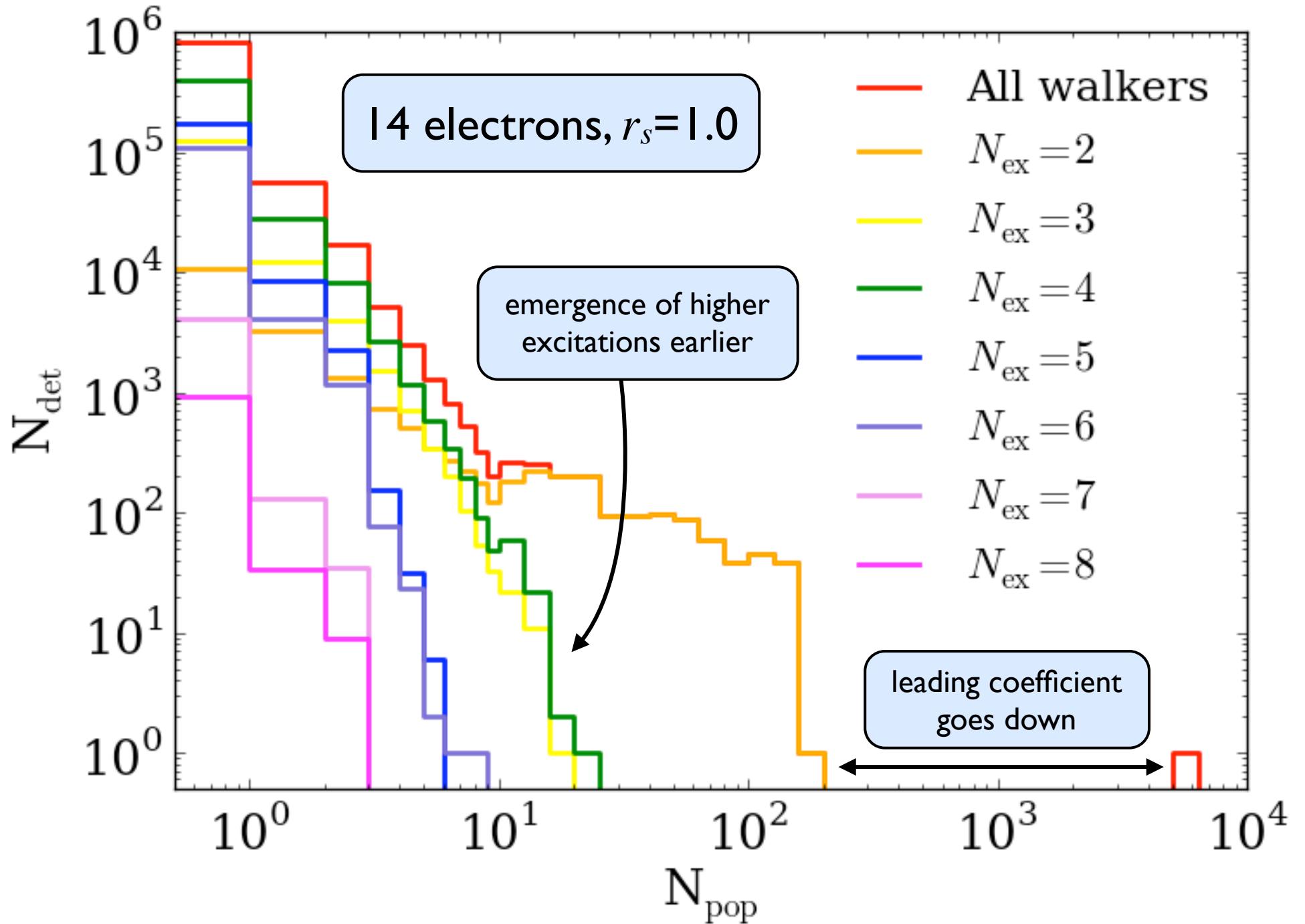
- Is the wavefunction recognisably multi-reference?
- What are the most important excitations?
- How does this compare with molecules?
- ...or other theories?

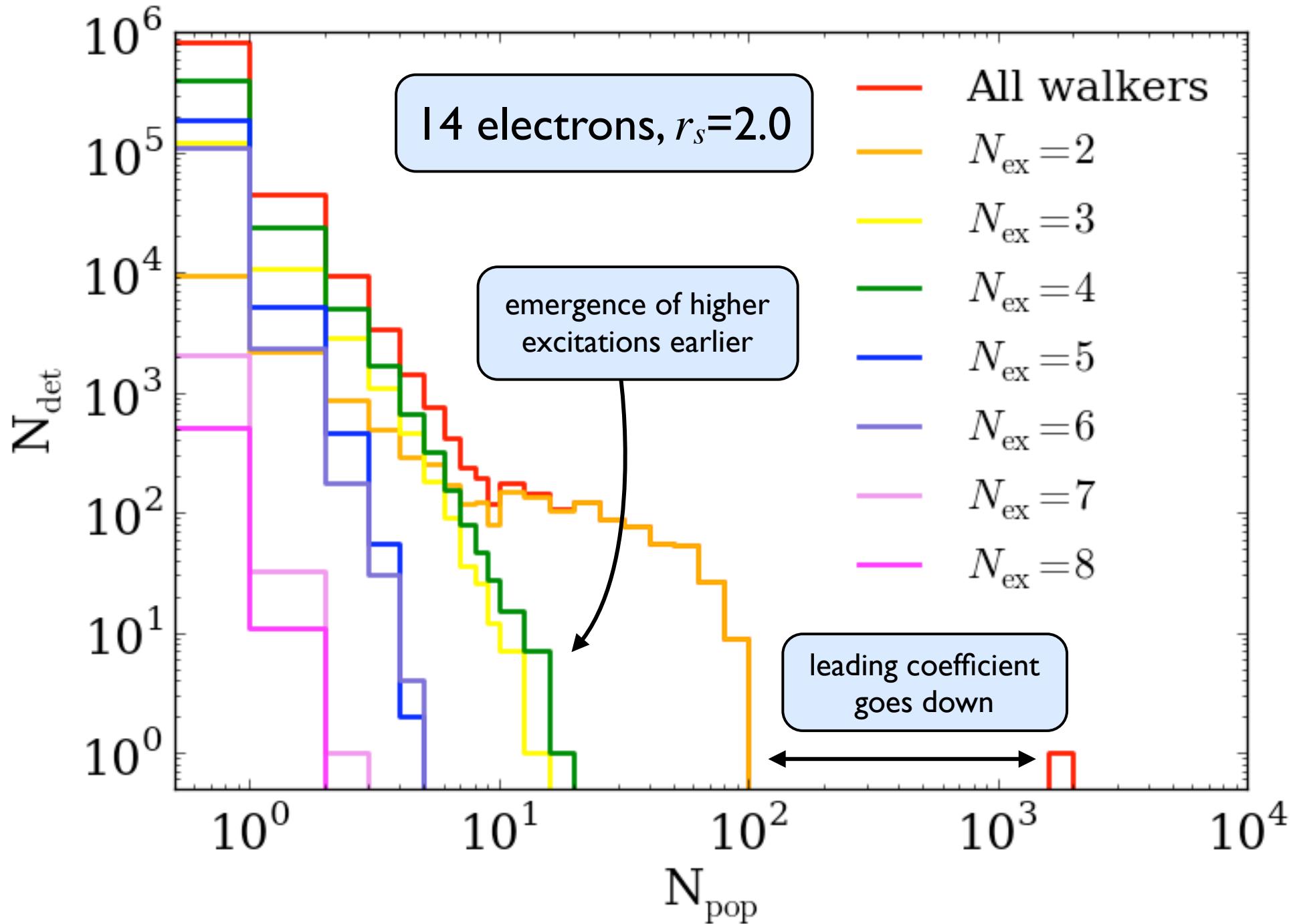
(Speculative/preliminary)

A remarkably simple ‘sign structure’?

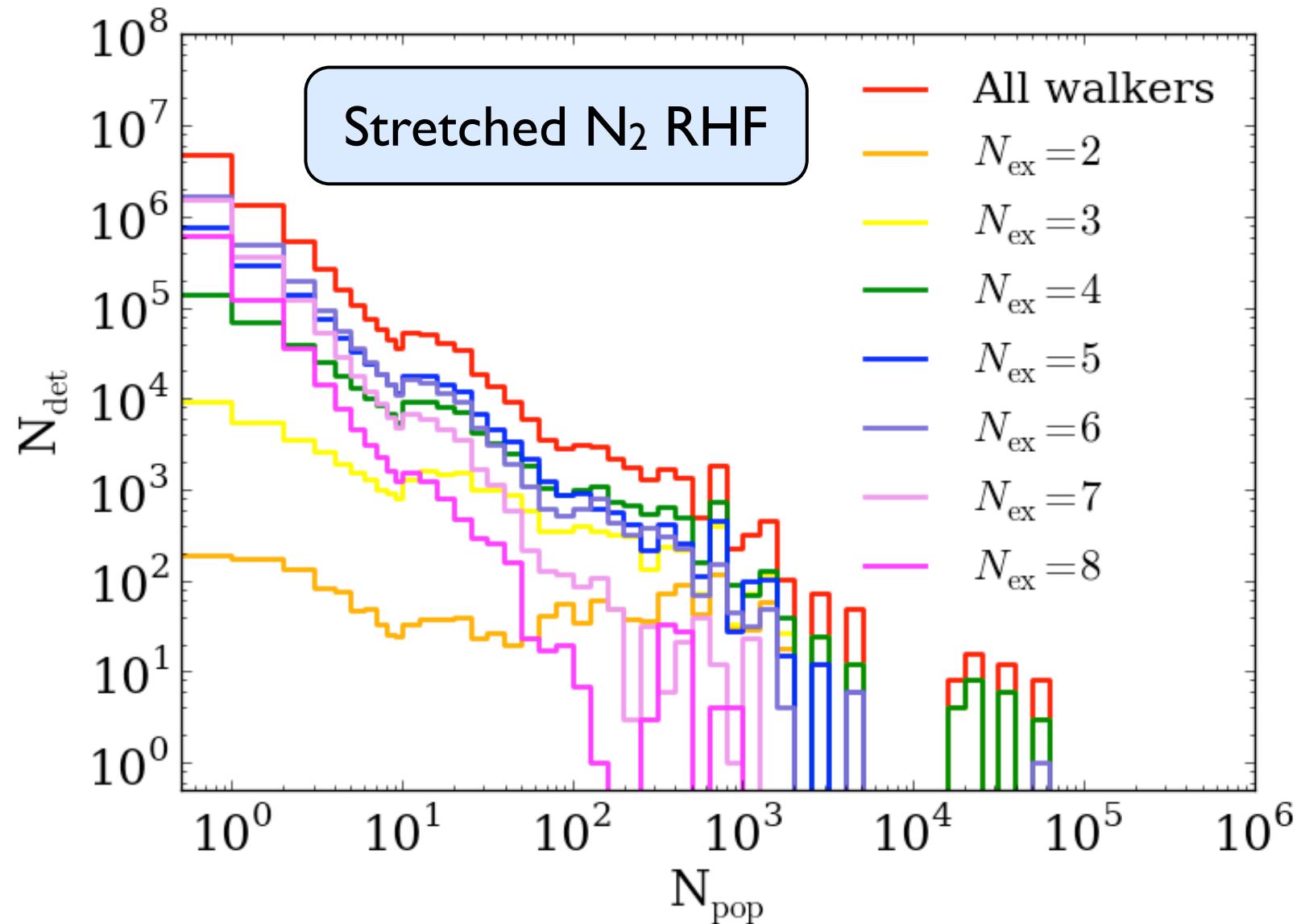




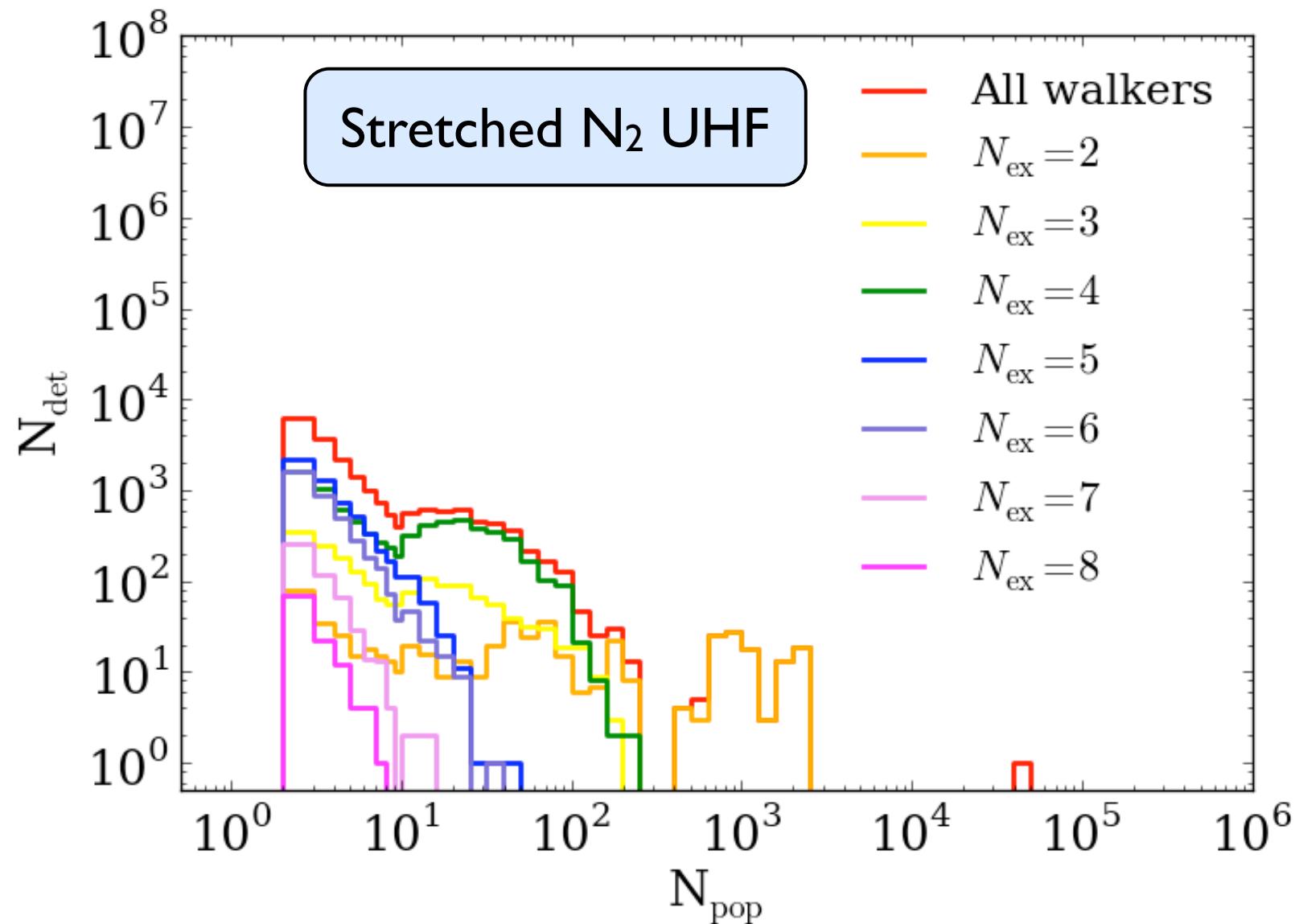


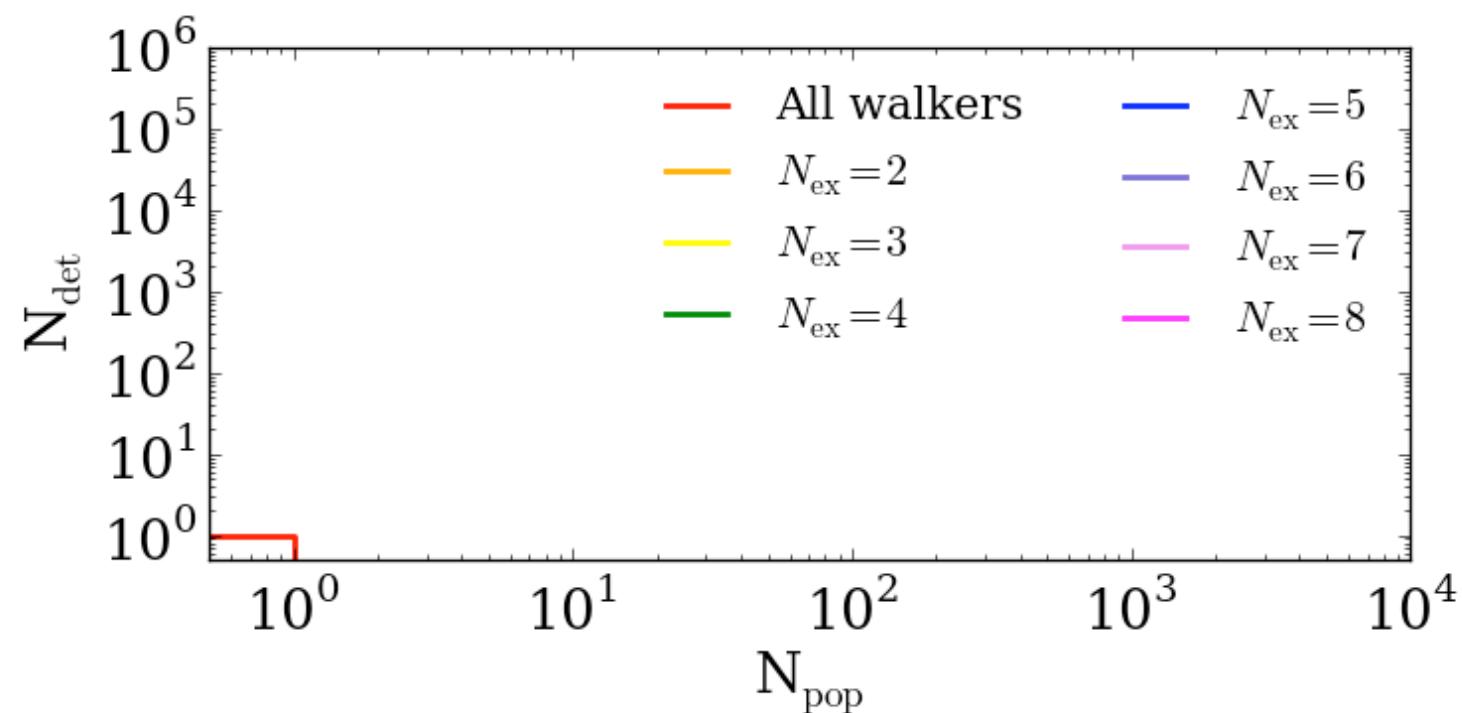
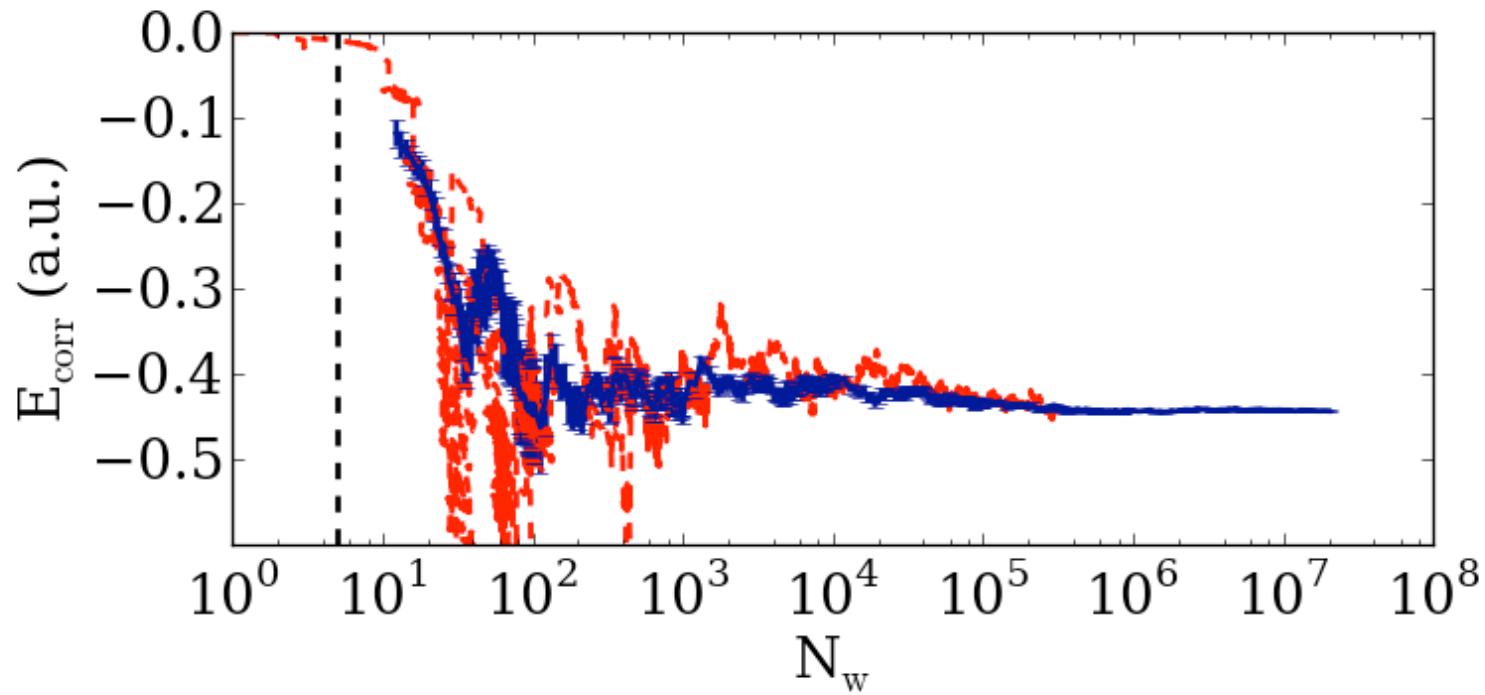


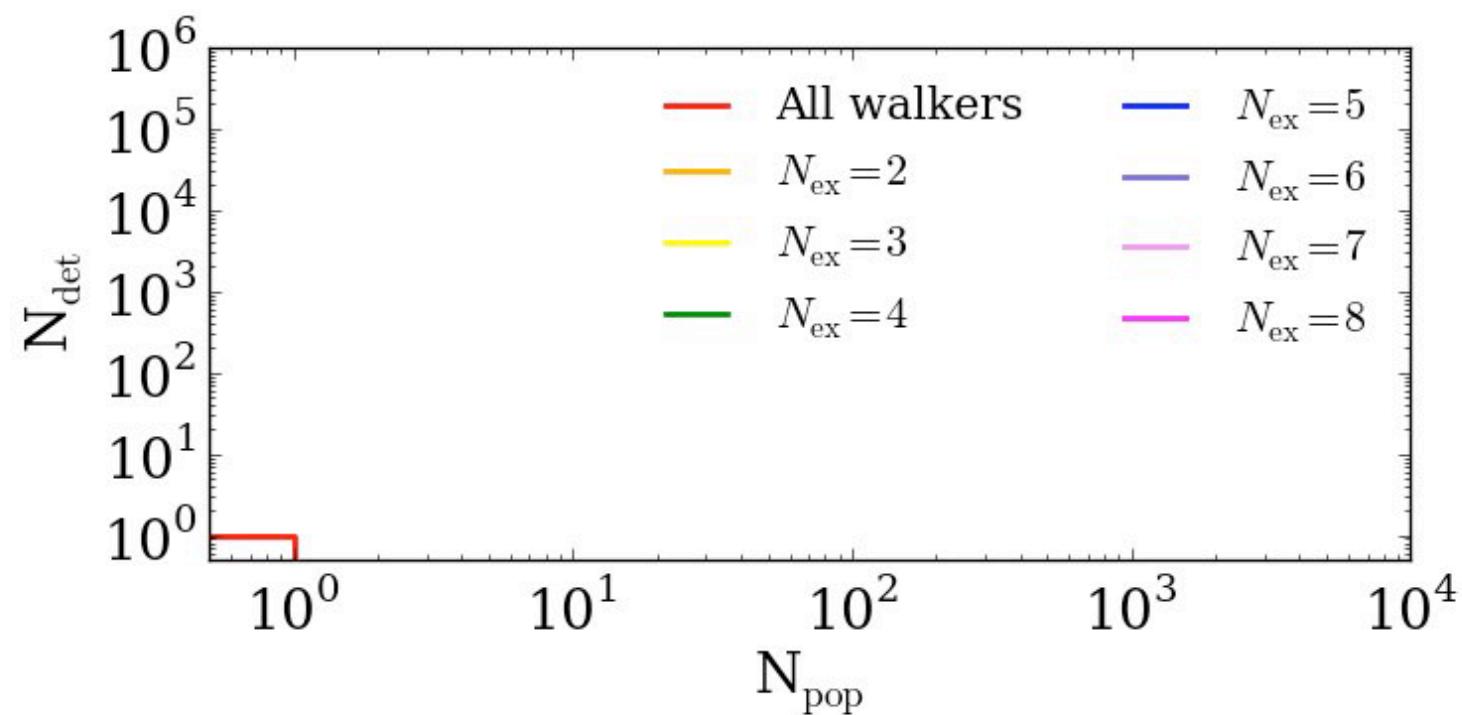
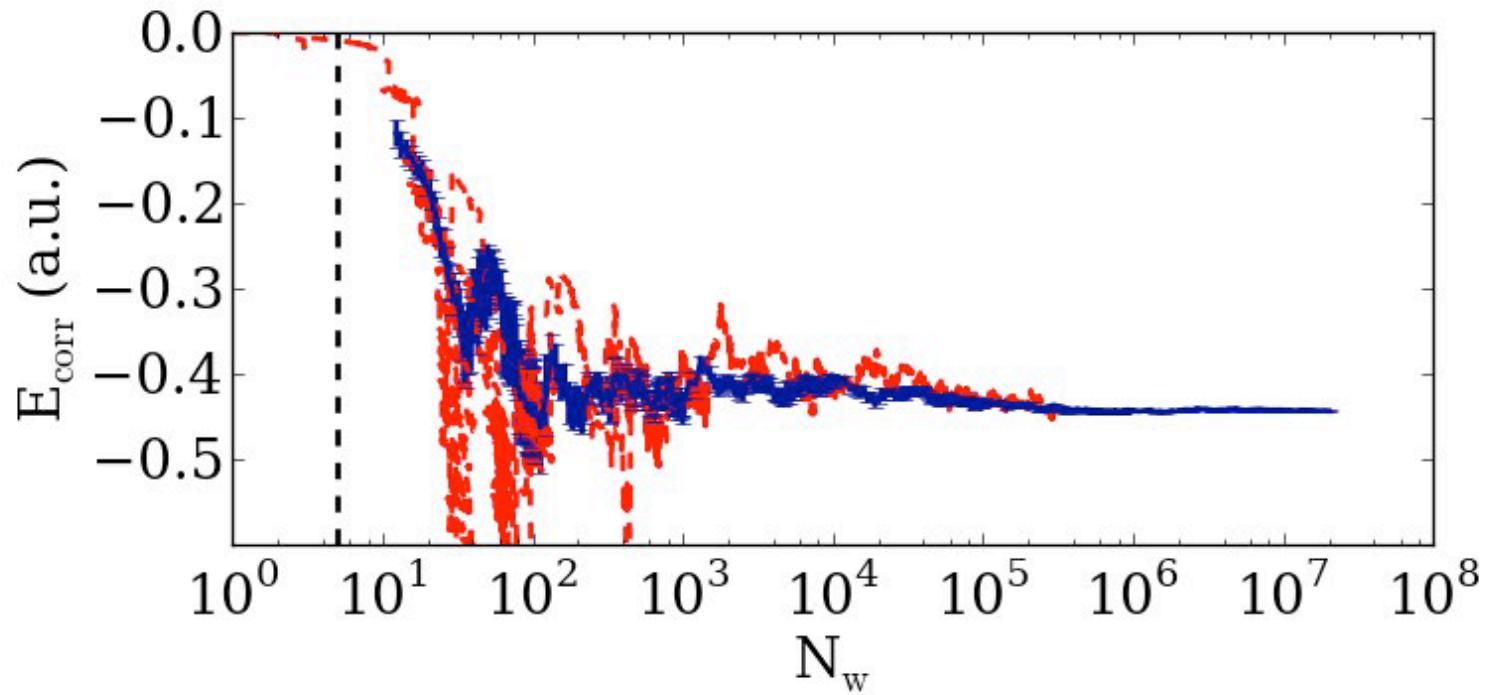
How does this relate to molecules?



How does this relate to molecules?







Acknowledgements

Ali Alavi, George Booth (NECI, Cambridge)

Georg Kresse, Andreas Grueneis (VASP, Vienna)

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Jennifer Mohr (1D/2D HEG)

Robert Thomas (N_2 RHF/UHF results)

Catherine Overy (discussions)

Special thanks to ‘Quantum Monte Carlo and the CASINO program IV’ 2009 mentors

References

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- i-FCIQMC: D Cleland, GH Booth, A Alavi, JCP, 132, 4, 041103 (2010)
- i-FCIQMC/HEG: JJS, GH Booth, A Gruneis, A Alavi, PRB, 85, 081103(R) (2012)
- i-FCIQMC/HEG/extrap/initiator error: JJS, GH Booth, A. Alavi, JCP, 136, 244101 (2012)
- HEG/extrapolation: JJS, A. Gruneis, GH Booth, G Kresse, A Alavi, PRB, 86, 035111 (2012)
- HEG/cusps: A Gruneis *et al.* in preparation
- HEG/divergences: JJS *et al.* in preparation