

Some Estimates in Variational Quantum Monte Carlo

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- Variational Monte Carlo
- 'Standard' VMC is $P = \psi^2$, and failure of CLT is an artifact of this
- Conditions and expressions for Normally distributed estimates in general sampling
- **Efficient** sampling Monte Carlo implemented - not sampling from $P = \psi^2$
- Estimates for energies, excitation energies, ionisation energies, transition moments
- Results for first row atoms and some molecules

VMC and General Sampling

Sample with $P = \psi^2/w$ and construct an estimate:

$$\bar{\mu} = \frac{\sum w_i E_L(\mathbf{R}_i)}{\sum w_i}$$

- If both variances exist, Fieller's theorem tells us this is a sample from a Normal distribution with:

$$\mu = \frac{\int \psi^2 E_L d\mathbf{R}}{\int \psi^2 d\mathbf{R}} \quad , \quad \sigma^2 = \frac{1}{r} \frac{\int \psi^2/w d\mathbf{R} \int w\psi^2 (E_L - \mu)^2 d\mathbf{R}}{[\int \psi^2 d\mathbf{R}]^2}$$

- We can estimate the variance:

$$\bar{\sigma}^2 = \frac{r}{r-1} \frac{\sum w_i^2 (E_L(\mathbf{R}_i) - \bar{\mu})^2}{(\sum w_i)^2}$$

- $\bar{\sigma}^2 \neq (\text{sample variance})/r$
- These equations do not follow from the usual (univariate) Central Limit Theorem
- *Zero Variance Principle* is still valid - for exact $\psi \Rightarrow \bar{\sigma} = 0$

→ The error is **controlled** if the bivariate CLT is **valid** and $\langle w \rangle \neq 0$

Trail JR, Phys. Rev. E. **77**, 016703,016704 (2008)

When is it Normal ?

For standard sampling $P = \psi^2$ ($w = 1$)

- Normal for standard sampling and total energy ($P \propto 1/x^4$)
- Not Normal for standard sampling used with correlated sampling, forces, and many other estimates
- Different choices of P (equivalently w) are possible
- Changes computational cost: flops for evaluating P
- Changes distribution of random errors
- Failure of CLT from singularities in averaged quantities on the nodal surface

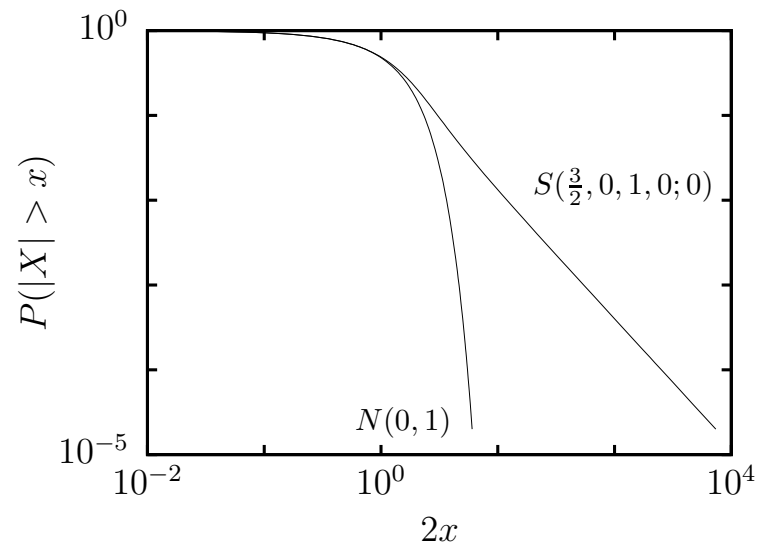
For some P (equivalently w) the variance is infinite, and CLT is invalid

→ Goal is to improve efficiency and reinstate the CLT where it is invalid for standard sampling

Why is Normality so important ?

We want Normally distributed estimates, from a general form of the Central Limit Theorem

If the distribution the estimates are drawn from is *not* normal then it is a **Stable Law** :



- Normal distribution and an example Stable law
- Probability that a sample fall outside of central interval size $2x$
- Width parameter is representative of error for Normal, not for Stable
- Width parameter is estimateable for Normal - sample standard error
- Width parameter is *not* estimateable for Stable - sample standard error is unrelated

Efficient sampling in VMC

- Draw position vectors from $P = |D_1|^2 + |D_2|^2$ ^a, no Jastrow, Backflow
→ Zero on coalescence planes only, non-zero on rest of nodal surface
- Perform Metropolis accept/reject with P , and use $w = \psi^2/P$
- Using r samples, the Bivariate CLT, and Fiellers theorem provides the Normal estimate

$$\bar{E}_{tot} = \frac{\sum w_i E_L(\mathbf{R}_i)}{\sum w_i}, \quad \bar{\sigma}^2 = \frac{r}{r-1} \frac{\sum w_i^2 (E_L(\mathbf{R}_i) - \bar{E}_{tot})^2}{(\sum w_i)^2}$$

- (wE_L, w) has no singularities and is bounded \Rightarrow all moments exist \Rightarrow distribution is Normal

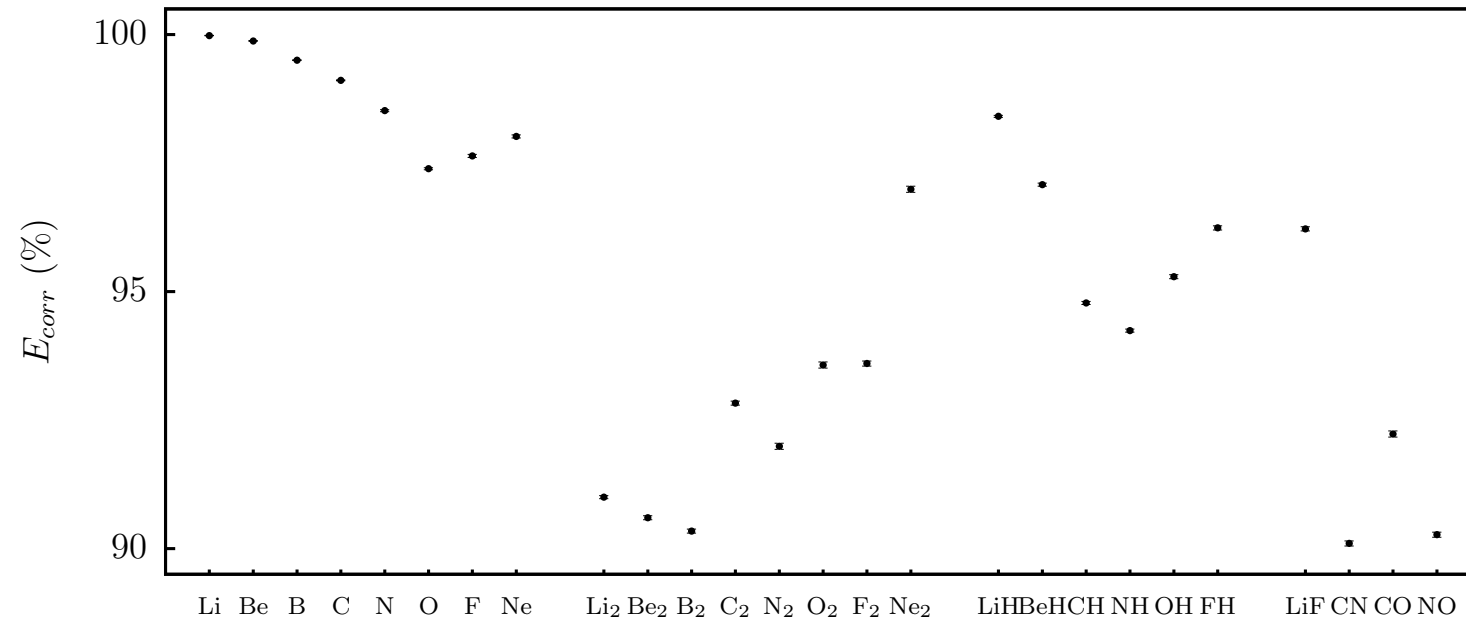
^aTrail JR and Maezono R, JCP (2010)

Efficient sampling in VMC

Optimisation and estimation of total energy:

- All-electron
- First row atoms + some diatomic molecules
- Numerical orbitals from ATSP2K and 2DHF (MCSCF and HF)
- Jastrow, Backflow, and 5 – 86 CSFs
- 48 h desktop time/system
- r for Final estimate:monitor estimate:optimisation $r = 9000 : 150 : 1$

Efficient sampling in VMC



- Improved results for less computational effort
- For a given error $5 - 35\times$ faster than standard sampling
- Surfaces for energy optimisation are Normal
- Do we need anisotropic Jastrow/Backflow ?

Distribution of more general estimates

What is the distribution for more complex estimates ?

Example: Energy differences

$$\Delta E_{tot} = \frac{\int \psi_1^2 E_1 d\mathbf{R}}{\int \psi_1^2 d\mathbf{R}} - \frac{\int \psi_2^2 E_2 d\mathbf{R}}{\int \psi_2^2 d\mathbf{R}}$$

- Sample with $P = |D_1|^2 + |D_2|^2$ (no nodal surface, only coalescence planes)

$$\begin{aligned} \text{Est} [\Delta E_{tot}] &= \frac{\sum w_1 E_1}{\sum w_1} - \frac{\sum w_2 E_2}{\sum w_2} \\ &= \frac{S_2}{S_1} - \frac{S_4}{S_3} \end{aligned}$$

- Elements in sums S_n are correlated only when in same 'time slots'

⇒ Multivariate CLT is true

⇒ Each S_n is Normal

⇒ Correlation between S_n and S_m is linear

Algebra for sums of random variables

What is the distribution of

$$\Delta E_{tot} = \frac{S_2 S_3 - S_1 S_4}{S_1 S_3}, \quad S_n = \sum_i X_n(i)$$

with parameters expressed in terms of the estimateable

$$\mathbb{E} [X_n(i)] = \mu_n$$

$$Var [X_n(i)] = C_{nn}$$

$$Cov [X_m(i), X_n(j)] = C_{nm} \delta_{i,j}$$

?

Algebra for sums of random variables : Sums

Adding sums of the random variables, trivial to show that:

- Distribution is Normal

$$\mathbb{E} [S_1 + S_2] = r(\mu_1 + \mu_2)$$

$$Var [S_1 + S_2] = r(C_{11} + 2.C_{12} + C_{22})$$

$$Cov [S_1 + S_2, S_3 + S_4] = r(C_{13} + C_{14} + C_{23} + C_{24})$$

→ All the $S_1 + S_2$ are Normal

→ Correlation between any $S_j + S_k$ and $S_l + S_m$ is linear

Algebra for sums of random variables : Quotients

Quotients of sums of the random variables, Fiellers theorem provides:

- Distribution is Normal

$$\mathbb{E} \left[\frac{S_2}{S_1} \right] = \frac{\mu_2}{\mu_1}$$
$$Var \left[\frac{S_2}{S_1} \right] = \frac{1}{r} \frac{1}{\mu_1^2} \left[C_{22} - 2 \frac{\mu_2}{\mu_1} C_{12} + \left(\frac{\mu_2}{\mu_1} \right)^2 C_{11} \right]$$

- No expression for correlation of different quotients

So,

$$\frac{S_1 + S_2}{S_3 + S_4}$$

is Normal and parameters are estimateable

Algebra for sums of random variables : Products

Products of sums of the random variables:

- What is the bivariate distribution of $(S_1 S_2, S_3 S_4)$?
- Derive co-moments and compare with bivariate Normal....define $\Delta X_1(i) = X_1(i) - \mu_1$

$$\begin{aligned} \Delta(S_1 S_2) &= S_1 \cdot S_2 - \mathbb{E}[S_1 \cdot S_2] \\ &= \sum_{ij} \Delta X_1(i) \Delta X_2(j) + r \sum_i [\mu_1 \Delta X_2(i) + \mu_2 \Delta X_1(i)] - r C_{12} \end{aligned}$$

- Co-moments are defined by $\mu_{m,n} = \mathbb{E}[\Delta(S_1 S_2)^m \cdot \Delta(S_3 S_4)^n]$

$$\begin{aligned} \mu_{m,n} &= \mathbb{E} \left[\left(\sum_{ij} \Delta X_1(i) \Delta X_2(j) + r \sum_i [\mu_1 \Delta X_2(i) + \mu_2 \Delta X_1(i)] - r C_{12} \right)^m \right. \\ &\quad \left. \times \left(\sum_{ij} \Delta X_3(i) \Delta X_4(j) + r \sum_i [\mu_3 \Delta X_4(i) + \mu_4 \Delta X_3(i)] - r C_{34} \right)^n \right] \end{aligned}$$

- Multiply out and count equivalent terms
- Pick out the dominant r terms.

Algebra for sums of random variables : Products

Powers of r in each term after multiplication:

- 2^{nd} order part: Powers p_1, q_1 from equivalence, none from prefactor
- 1^{st} order part: Powers p_2, q_2 from equivalence, p_2, q_2 from prefactor
- 0^{th} order part: Powers p_3, q_3 from equivalence, p_3, q_3 from prefactor
- Multiplication gives $p_1 + p_2 + p_3 = m, q_1 + q_2 + q_3 = n$
- Count powers of r from prefactors $[r]$, and from sums (r)

⇒ Terms have powers of r :

$$[1]^{p_1} (r)^{p_1} [r]^{p_2} (r)^{p_2} [r]^{p_3} (1)^{p_3} \cdot [1]^{q_1} (r)^{q_1} [r]^{q_2} (r)^{q_2} [r]^{q_3} (1)^{q_3} \mathbb{E} [\dots]$$

But Expectations are zero if any index is unique:

$$\begin{aligned} \mathbb{E} [\Delta X_1(i) \cdot \Delta X_2(j) \cdot \Delta X_3(k) \cdots] &= \mathbb{E} [\Delta X_1(i)] \cdot \mathbb{E} [\Delta X_2(j) \cdot \Delta X_3(k) \cdots] \\ &= 0 \end{aligned}$$

- Overcounting due to including zero expectations values

Algebra for sums of random variables : Products

Count only terms with no unique indices:

$$\begin{aligned}
 & r^{(m+n)+(p_2+q_2)/2} \mathbb{E} [\dots] \quad \text{for } (p_2 + q_2) \text{ even} \\
 & r^{(m+n)+(p_2+q_2+1)/2} \mathbb{E} [\dots] \quad \text{for } (p_2 + q_2) \text{ odd}
 \end{aligned}$$

⇒ In large r limit $p_2 + q_2 = m + n$ terms dominate ($p_1 = p_3 = q_1 = q_3 = 0$)

⇒ In large r limit the co-moments are same as those for *sums*:

$$\begin{aligned}
 \mathbb{E} [\Delta(S_1 S_2)^m \cdot \Delta(S_3 S_4)^n] &= r^{m+n} \mathbb{E} \left[\left(\sum_i [\mu_1 \Delta X_2(i) + \mu_2 \Delta X_1(i)] \right)^m \right. \\
 &\quad \left. \times \left(\sum_i [\mu_3 \Delta X_4(i) + \mu_4 \Delta X_3(i)] \right)^n \right]
 \end{aligned}$$

Algebra for sums of random variables : Products

- In large r limit (S_1S_2, S_3S_4) is bivariate Normal with

$$\mathbb{E}[S_1S_2] = r^2\mu_1\mu_2$$

$$Var[S_1S_2] = r^3(\mu_1^2C_{22} + 2\mu_1\mu_2C_{12} + \mu_2^2C_{11})$$

$$Cov[S_1S_2, S_3S_4] = r^3(\mu_1\mu_3C_{24} + \mu_1\mu_4C_{23} + \mu_2\mu_3C_{14} + \mu_2\mu_4C_{13})$$

→ We have rules for obtaining the distribution of combinations of sums of random variables (that are Normal)

Energy differences

$$\begin{aligned}\text{Est} [\Delta E_{tot}] &= \frac{\sum w_1 E_1}{\sum w_1} - \frac{\sum w_2 E_2}{\sum w_2} \\ &= \frac{S_2 S_3 - S_1 S_4}{S_1 S_2}\end{aligned}$$

$$\begin{aligned}\overline{\Delta E}_{tot} &= \frac{\bar{\mu}_2}{\bar{\mu}_1} - \frac{\bar{\mu}_4}{\bar{\mu}_3} \\ \frac{r-1}{r} \overline{\sigma}^2 &= \frac{\sum w_1^2 (E_1 - \bar{\mu}_1)^2}{[\sum w_1]^2} - 2 \frac{\sum w_1 w_2 (E_1 - \bar{\mu}_1)(E_2 - \bar{\mu}_2)}{[\sum w_1][\sum w_2]} + \frac{\sum w_2^2 (E_2 - \bar{\mu}_2)^2}{[\sum w_2]^2}\end{aligned}$$

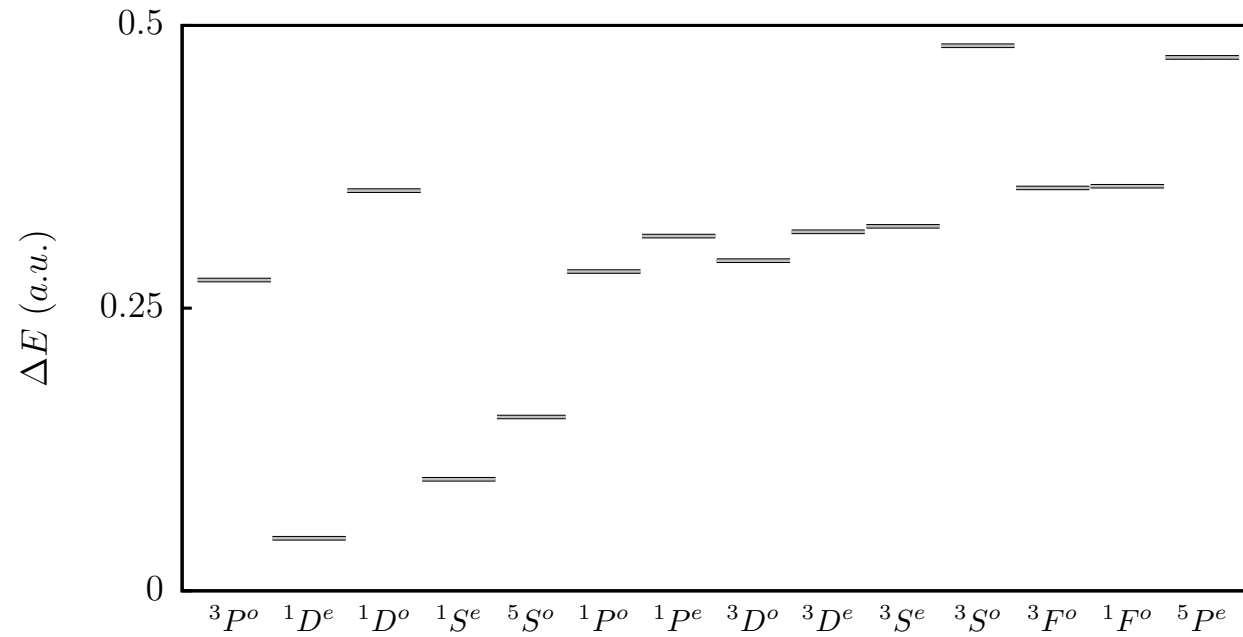
Sufficient conditions for this estimate to be Normal are

- All moments exist
- All means in the denominator are non-zero
- At least one mean in each product on numerator is non-zero

Carbon atom excitation energies

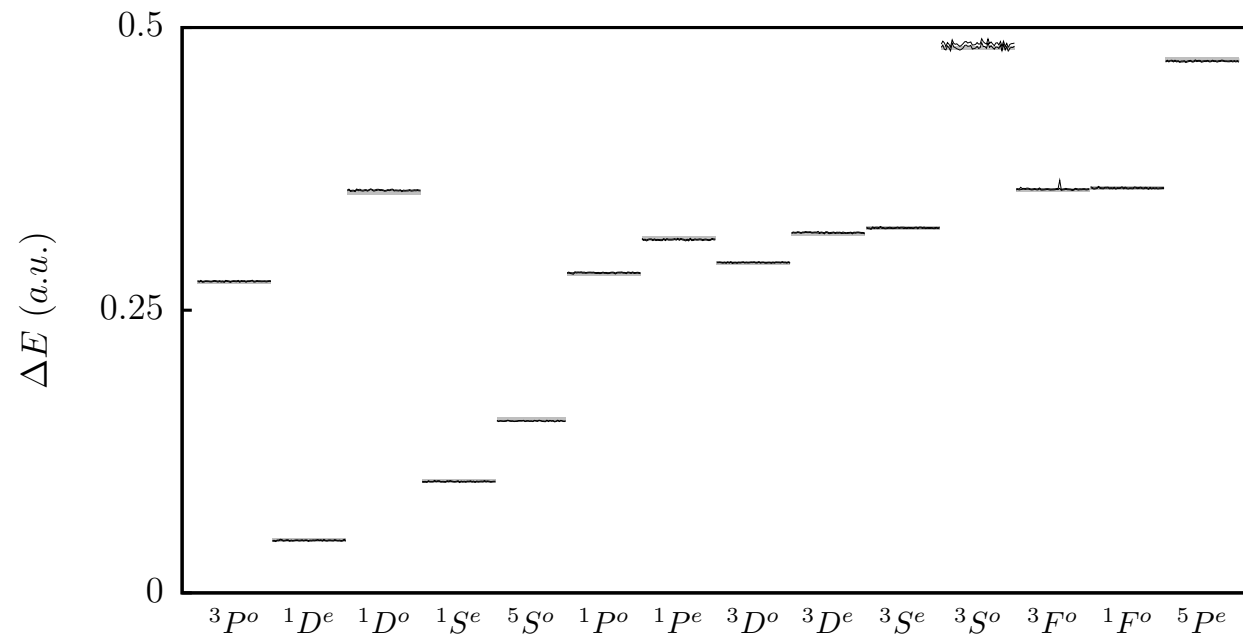
- Calculation as for GS (Multideterminant and numerical orbitals from MCSCF \sim 300 parameters)
- $E_{tot} [{}^{2S+1}L] - E_{tot} [{}^3P]$
- Lowest two eigenstates for each Term
- Inversion symmetry conserved by Jastrow and Backflow
- Term *approximately* conserved by introduction of Jastrow and Backflow
- Energy minimisation approximately valid for lowest two energies for each Term
- Sample with $P = D_1 [{}^{2S+1}L]^2 + D_2 [{}^{2S+1}L]^2 + D_1 [{}^3P]^2 + D_2 [{}^3P]^2$

Carbon atom excitation energies



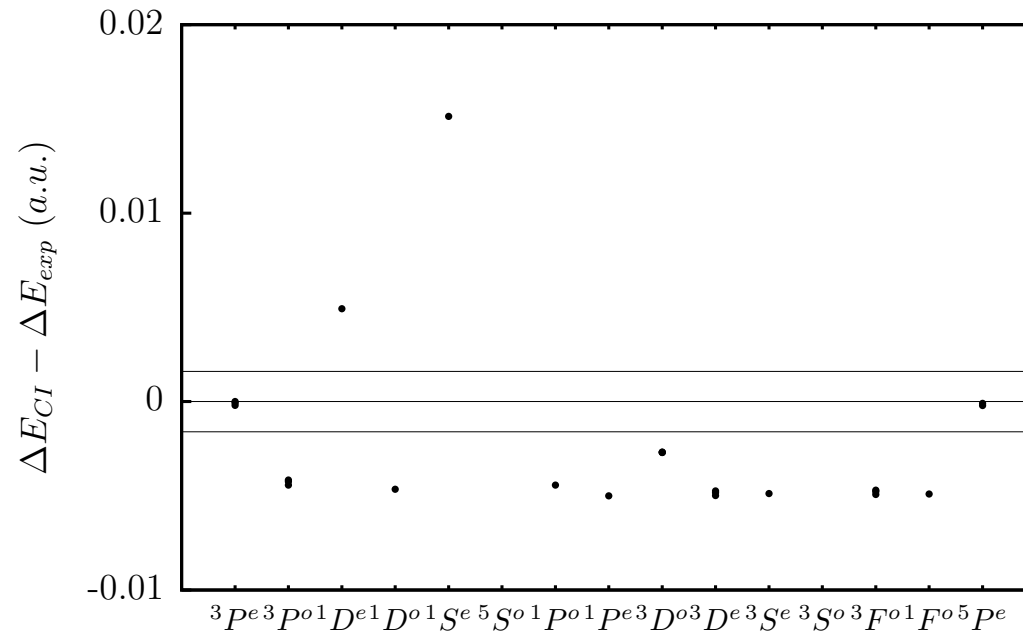
- $C^{2S+1}L \rightarrow C^3P$
- Grey: experimental Spectroscopic values \pm 'chemical accuracy'

Carbon atom excitation energies



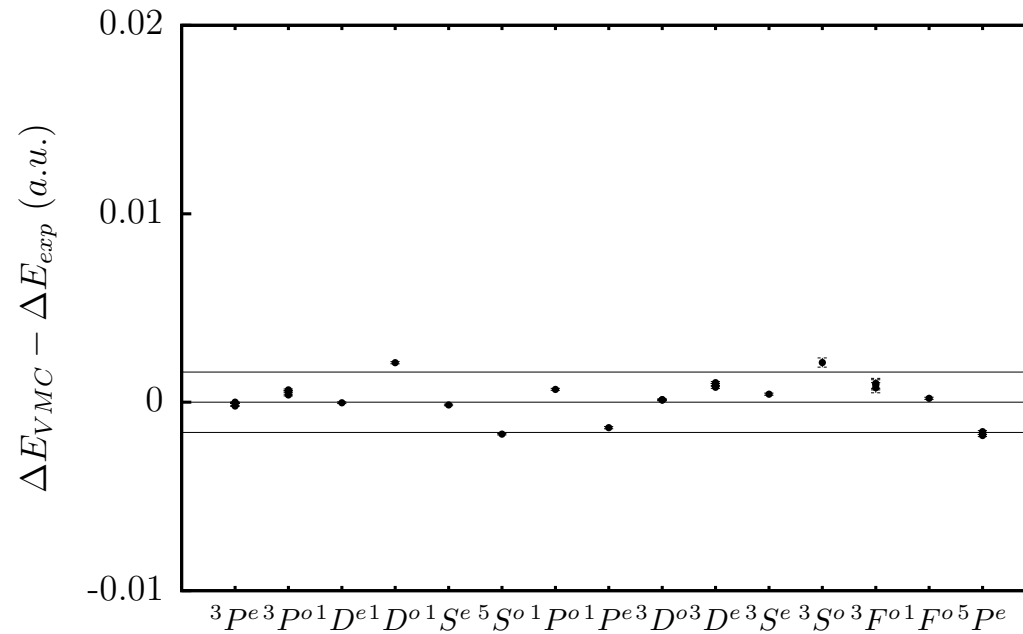
- $C^{2S+1}L \rightarrow C^3P$
- Grey: experimental Spectroscopic values \pm 'chemical accuracy'
- 48 energy difference estimates for each excitation

Error in Carbon atom excitation energies



- CI results used for allocating transitions to lines
- CI AS and orbitals chosen empirically to reproduce spectroscopy

Error in Carbon atom excitation energies



- Chemical accuracy from VMC + efficient sampling
- *NOT* spectroscopic accuracy
- Estimate of difference *not* difference of estimates
- Correlation reduces error by 10 – 70%

Ionization energies

- How do we deal with changes in electron number?

$$\begin{aligned}
 E_{ion} &= E_n - E_{n-1} \\
 &= \frac{\int \psi_1^2 E_1 d\mathbf{R}_n}{\int \psi_1^2 d\mathbf{R}_n} - \frac{\int \psi_2^2 E_2 d\mathbf{R}_{n-1}}{\int \psi_2^2 d\mathbf{R}_{n-1}}
 \end{aligned}$$

- Sampled $P_n(\mathbf{R}_n) = D_1(n)^2 + D_2(n)^2$
- For estimating E_{n-1} ignore *one* 3d sample vector, so $\mathbf{R}_n \rightarrow \mathbf{R}_{n-1}$
- Distribution of $(w_2 E_2, w_2)$ given by integrating P_n analytically

$$\begin{aligned}
 P_{n-1} &= \int D_1(n)^2 + D_2(n)^2 d^3 \mathbf{r}_n \\
 &= \int [\phi_1(\mathbf{r}_n) \cdot C_1(1, n) + \phi_2(\mathbf{r}_n) \cdot C_1(2, n) + \dots]^2 + [\phi_1(\mathbf{r}_n) \cdot C_2(1, n) + \phi_2(\mathbf{r}_n) \cdot C_2(2, n) + \dots]^2 d^3 \mathbf{r}_n \\
 &= [C_1(1, n)^2 + C_1(2, n)^2 + \dots] + [C_2(1, n)^2 + C_2(2, n)^2 + \dots]
 \end{aligned}$$

Ionization energies

- Provides distributions of $(w_1 E_1, w_1, w_2 E_2, w_2)$ in terms of weights

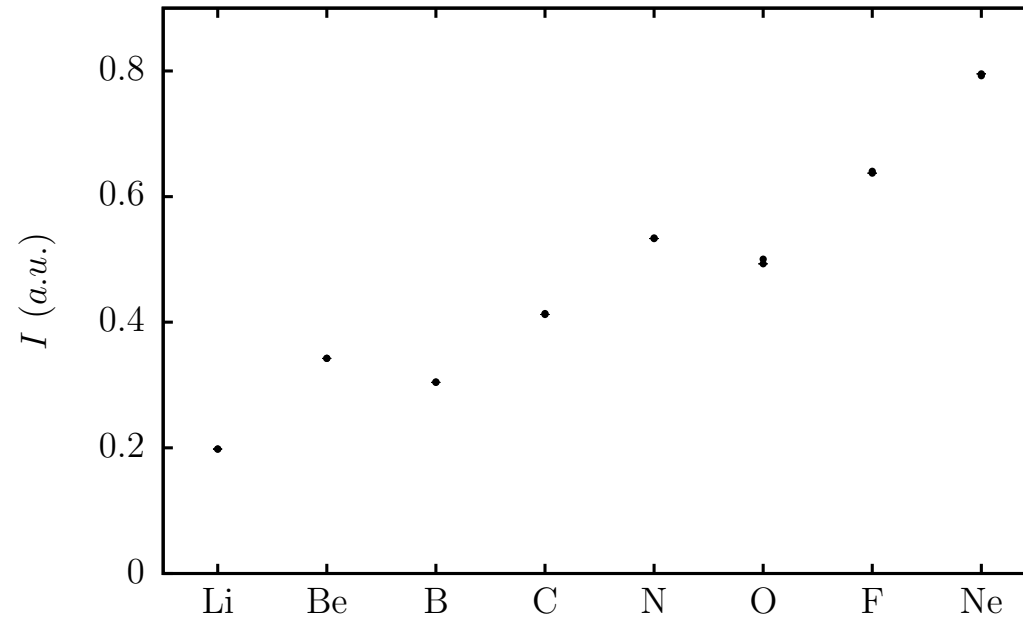
$$\begin{aligned}w_1(\mathbf{R}_n) &= \psi_1^2(\mathbf{R}_n)/P_n(\mathbf{R}_n) \\w_2(\mathbf{R}_{n-1}) &= \psi_2^2(\mathbf{R}_{n-1})/P_{n-1}(\mathbf{R}_{n-1})\end{aligned}$$

- P_n is zero on coalescence planes only
- P_{n-1} is zero on coalescence planes only
- All sums are Normal and linearly correlated \rightarrow Normal estimate

Ionization energies

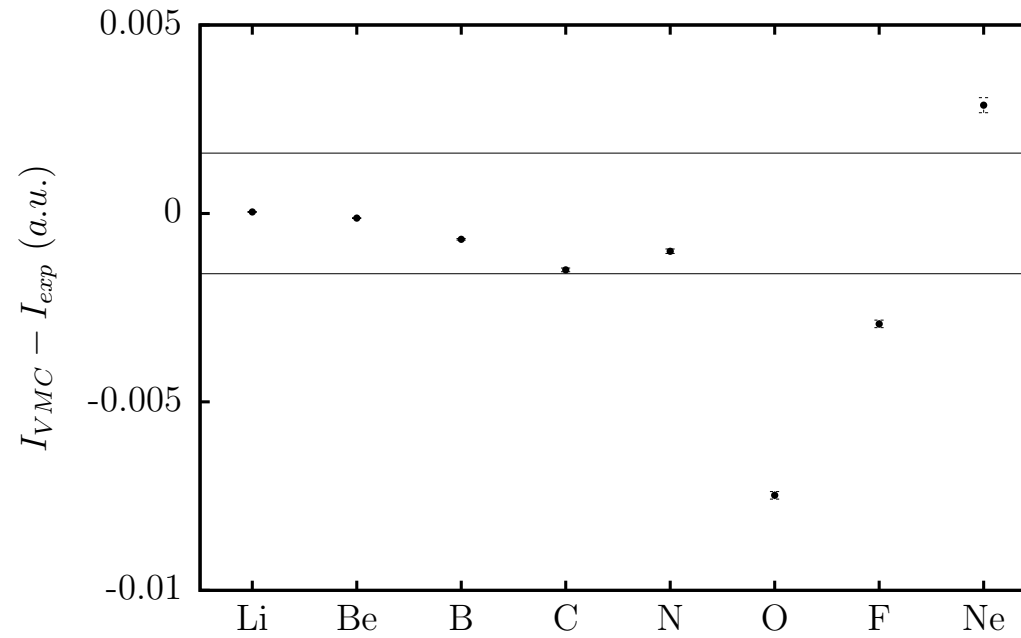
- First row neutral atoms and ions
- Orbitals from same source as before
- Jastrow/Backflow and computational cost as before
- Optimise neutral atom and ion separately, with energy minimisation
- Estimate energy difference/error as for excitation energies

Ionization energies



- Experimental ionization energies
- VMC estimated energy difference

Error in Ionization energies



- Close but *not* chemical accuracy
- Maybe we need T/Q excitations ?
- Correlation reduces error by 0 – 40%

Transition moments

Spectroscopic line widths characterised by transition dipole moments

$$\Delta_\omega = \frac{3}{2} \Delta E \sum_m |\langle \psi_0 | \mathbf{R} | \psi_m \rangle|^2$$

with sum over total angular momentum eigenstates.

Not done yet ... start with estimates for transition dipole moments:

$$t_{12} = |\langle \psi_1 | \sum_i \mathbf{r}_i | \psi_2 \rangle|^2$$

- Sample using $P = D_1(1)^2 + D_2(1)^2 + D_1(2)^2 + D_2(2)^2$
- Get an error estimate from random variable algebra
- Not zero variance...

Transition moments

Estimate:

$$\begin{aligned}
 t_{12} &= |\langle \psi_1 | \sum \mathbf{r}_i | \psi_2 \rangle|^2 \\
 &= \frac{[\int \psi_1 \psi_2 (x_1 + \dots + x_n) d\mathbf{R}]^2 + [\int \psi_1 \psi_2 (y_1 + \dots + y_n) d\mathbf{R}]^2 + [\int \psi_1 \psi_2 (z_1 + \dots + z_n) d\mathbf{R}]^2}{\int \psi_1^2 d\mathbf{R} \int \psi_2^2 d\mathbf{R}}
 \end{aligned}$$

- Define weights $w_1 = \psi_1^2/P, w_2 = \psi_2^2/P$
- Analyse as before ...

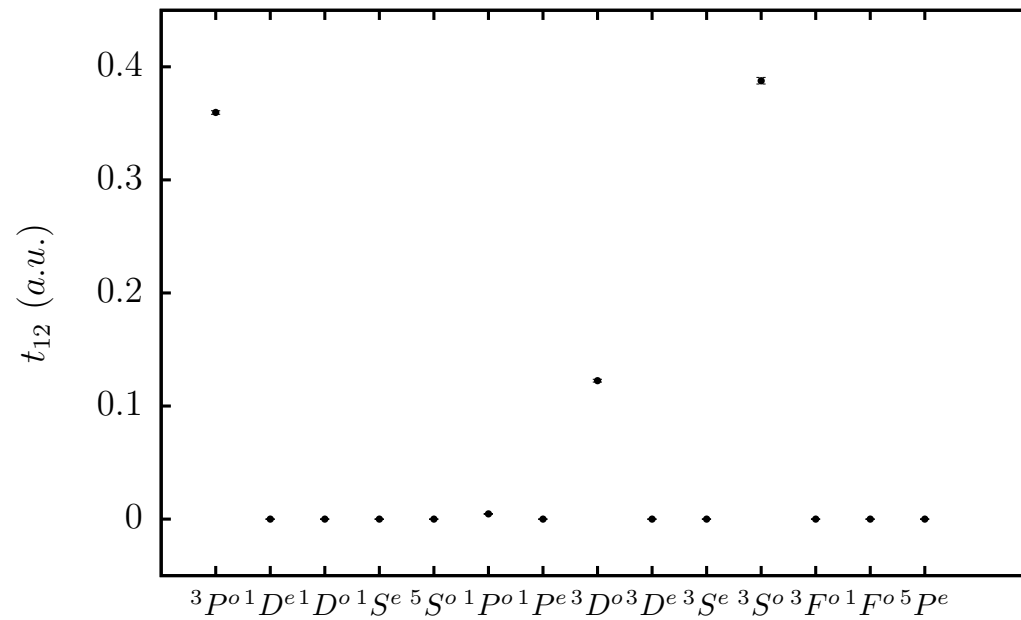
$$\text{Est} [t_{12}] = \frac{S_1^2 + S_2^2 + S_3^2}{S_4 \cdot S_5}$$

- Normally distributed with estimateable mean and variance:

$$\begin{aligned}
 \bar{t}_{12} &= \frac{\mu_1^2 + \mu_2^2 + \mu_3^2}{\mu_4 \mu_5} \\
 r \mu_4^2 \mu_5^2 \text{Var}[\bar{t}_{12}] &= [4\mu_1^2 C_{11} + 4\mu_2^2 C_{22} + 4\mu_3^2 C_{33} + 8\mu_1 \mu_2 C_{12} + 8\mu_1 \mu_3 C_{13} + 8\mu_2 \mu_3 C_{23}] \\
 &\quad - 2\bar{t}_{12} [2\mu_1 \mu_5 C_{14} + 2\mu_2 \mu_5 C_{24} + 2\mu_3 \mu_5 C_{34} + 2\mu_1 \mu_4 C_{15} + 2\mu_2 \mu_4 C_{25} + 2\mu_3 \mu_4 C_{35}] \\
 &\quad + \bar{t}_{12}^2 [\mu_5^2 C_{44} + 2\mu_4 \mu_5 C_{45} + \mu_4^2 C_{55}]
 \end{aligned}$$

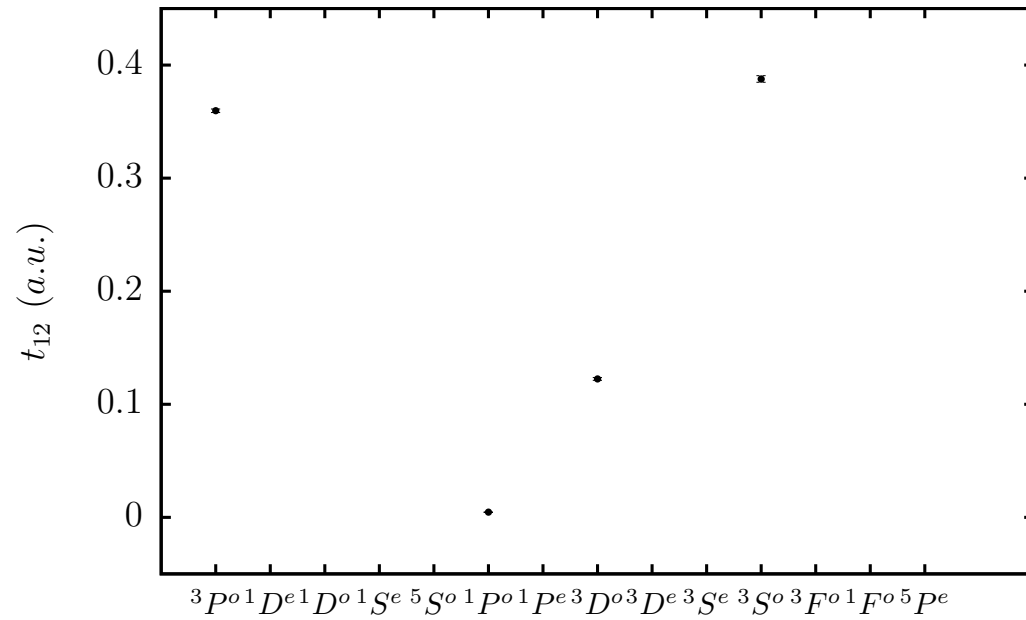
- Replace C 's and μ 's with unbiased estimates

Transition moments



- Many are non-normal and zero from symmetry considerations...

Transition moments



- ... so we drop them
- Random error $\sim 1\%$

Conclusions

- Normal errors can be reintroduced
- More computationally efficient than standard sampling
- Optimisation is on a Normal surface, unlike standard sampling
- Distribution of random error can be derived for general estimates and sampling

Normal and efficient estimates implemented for:

- Total energies
- Energy differences and Ionization energies
- Transition moments

Next?

- Spectroscopic line widths
- Electron affinities
- Optimise orthogonalised trial wavefunctions for more excited states
- Normal force estimates and geometry optimisation surfaces
- Generalised DMC