



The quantum Monte Carlo method can determine the total energy of an assembly of quantum particles to a high degree of accuracy.

In variational Monte Carlo (VMC), a trial wave function $\Psi_T(\mathbf{R})$ is used to evaluate an upper bound to the exact ground-state energy E_0 .

$$E_{VMC} = \sum_{i=1}^{M} \frac{\hat{H}(\boldsymbol{R}_{i}) \boldsymbol{\Psi}_{T}(\boldsymbol{R}_{i})}{\boldsymbol{\Psi}_{T}(\boldsymbol{R}_{i})} \simeq \frac{\int |\boldsymbol{\Psi}_{T}(\boldsymbol{R})|^{2} \frac{\hat{H}(\boldsymbol{R}) \boldsymbol{\Psi}_{T}(\boldsymbol{R})}{\boldsymbol{\Psi}_{T}(\boldsymbol{R})} d\boldsymbol{R}}{\int |\boldsymbol{\Psi}_{T}(\boldsymbol{R})|^{2} d\boldsymbol{R}} \ge E_{0}$$

In diffusion Monte Carlo (DMC), imaginary-time propagation is used to project out the higher-energy components of $\Psi_T(\mathbf{R})$.

♦ For fermions, the *fixed-node approximation* is required. It is equivalent to constraining the *nodal surface* of the DMC wave function to equal that of $\Psi_{\tau}(\mathbf{R})$.

The following are typical fermionic trial wave functions:

 $\Psi_{T}(\mathbf{R}) = \Psi_{S}(\mathbf{R})$ (Slater or Hartree-Fock type)

where $\Psi_{s}(\mathbf{R})$ is a Slater determinant (or a multi-determinant expansion) of suitable one-particle orbitals. With a single determinant, there are **no correlation effects**.

 $\Psi_{T}(\mathbf{R}) = e^{J(\mathbf{R})} \Psi_{S}(\mathbf{R})$ (Slater-Jastrow type)

where $e^{J(\mathbf{R})}$ is a Jastrow correlation factor containing optimizable parameters.

$$\Psi_{T}(\mathbf{R}) = e^{J(\mathbf{R})} \Psi_{S}[\mathbf{X}(\mathbf{R})]$$

(Slater-Jastrow-backflow type)

where $X(\mathbf{R})$ is a set of collective coordinates containing optimizable parameters.

Key points:

- The quality of $\Psi_T(\mathbf{R})$, and in particular of its nodal surface, is paramount for the accuracy of the results.
- The Jastrow factor is the most important tool to improve upon $\Psi_{s}(\mathbf{R})$, but it does not change its nodes.
- There are ways to modify the nodal surface of $\Psi_{s}(\mathbf{R})$: orbital optimization, backflow transformations, multideterminant expansions, multi-Pfaffian wfns, etc. This is a very active area of research.

The electron-hole system is a model for excited semiconductors.

It is the simplest model system after the homogeneous electron gas, yet its phase diagram remains largely unknown.

Some electron-hole systems, such as the two-dimensional bilayer, can be recreated experimentally and display very interesting properties.



The aim of this work is to determine the phase diagram of electron-hole systems using QMC, to a greater degree of accuracy than previous studies.

Previous QMC studies have made use of the following scheme:



"**Nodal argument**": different orbitals give different nodes give different energies. Hence there is a correspondence between phase/orbitals (input) and energies (output).

However, in the limit of a perfect wave function:



An appropriate approach is:



And on the plus side, this requires fewer calculations.

The wave functions are formed using products of Slater determinants:

$$D[\phi] = \begin{vmatrix} \phi_{11} & \phi_{12} & \cdots & \phi_{1N} \\ \phi_{21} & \phi_{22} & \cdots & \phi_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ \phi_{NI} & \phi_{N2} & \cdots & \phi_{NN} \end{vmatrix}$$

with:

Pairing:

Fluid:

Crystal:

$$\phi_{ij} = \phi_L(\boldsymbol{e}_i - \boldsymbol{h}_j)$$

$$\Psi_S = D_{e\uparrow h\downarrow} D_{e\downarrow h\uparrow}$$

$$\phi_{ij} = e^{i\boldsymbol{k}_j \cdot \boldsymbol{r}_i}$$

$$\Psi_S = D_{e\uparrow} D_{e\downarrow} D_{h\uparrow} D_{h\downarrow}$$

$$\phi_{ij} = \phi_C(\boldsymbol{r}_i - \boldsymbol{R}_j)$$

$$\Psi_S = D_{e\uparrow} D_{e\downarrow} D_{h\uparrow} D_{h\downarrow}$$

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The phase diagram of the symmetric 2D electron-hole bilayer has already been studied.



S. de Palo et al, Phys. Rev. Lett. 88, 206401 (2002)



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Fluid, E=-0.17879(3) a.u.



Exponential, E=-0.18937(3) a.u.









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VMC results [variance minimization]





VMC results [energy minimization]





BF-VMC results [energy minimization]





BF-VMC results [energy minimization]















How about the crystalline phase?



S. de Palo et al, Phys. Rev. Lett. 88, 206401 (2002)



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with:
Pairing:
Fluid:

$$\phi_{ij} = \sum_{p=0}^{P} c_p e^{ik_p \cdot (e_i - h_j)} + \phi_L(e_i - h_j)$$

$$\Psi_S = D_{e^{\uparrow}h^{\downarrow}} D_{e^{\downarrow}h^{\uparrow}}$$
Fluid:

$$\psi_{ij} = \phi_C(\mathbf{r}_i - \mathbf{R}_j)$$

$$\Psi_S = D_{e^{\uparrow}} D_{e^{\downarrow}} D_{h^{\uparrow}} D_{h^{\downarrow}}$$

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$$\phi_{ij} = \sum_{p=0}^{P} c_p e^{ik_p \cdot (e_i - h_j)} + \phi_L(e_i - h_j)$$
Fluid:
Fluid:

$$\psi_{ij} = \sum_{q=0}^{Q} d_q e^{ik_q \cdot (r_i - R_j)} + \phi_C(r_i - R_j)$$

$$\Psi_S = D_{e\uparrow} D_{e\downarrow} D_{h\uparrow} D_{h\downarrow}$$

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$$\Gamma_{ij} = \phi(\boldsymbol{e}_{i}, \boldsymbol{h}_{j}) + \sum_{l=0}^{L} \psi(\boldsymbol{e}_{i} - \boldsymbol{R}_{l}) \psi(\boldsymbol{h}_{j} - \boldsymbol{R}_{l})$$
with:
Pairing:

$$\phi_{ij} = \sum_{p=0}^{P} c_{p} e^{i\boldsymbol{k}_{p} \cdot (\boldsymbol{e}_{i} - \boldsymbol{h}_{j})} + \phi_{L}(\boldsymbol{e}_{i} - \boldsymbol{h}_{j})$$
Fluid:
Crystal:

$$\psi_{ij} = \sum_{q=0}^{Q} d_{q} e^{i\boldsymbol{k}_{q} \cdot (\boldsymbol{r}_{i} - \boldsymbol{R}_{j})} + \phi_{C}(\boldsymbol{r}_{i} - \boldsymbol{R}_{j})$$

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Conclusions and further work:

Improvements over previous QMC calculations are possible. Need a different approach involving:

More general wave functions.

Expectation values of density-matrix related objects.

Robust wave-function optimization is a key element to get results efficiently.

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Real test: compare against experiment.

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