Backflow transformations in QMC: extension to inhomogeneous systems

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Theory of Condensed Matter group Cavendish Laboratory. University of Cambridge The backflow transformation

Application and results

Conclusions

The backflow transformation

Motivation:

 Backflow can change nodal surface; may overcome the fixed-node approximation

 Application of backflow to homogeneous systems successful

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Application and results

Conclusions

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• The Slater-Jastrow wave-function is:

 $\Psi_{SJ} = e^{J(\boldsymbol{R})} \Psi_{S}(\boldsymbol{R})$

with Ψ_s = Slater (multi-) determinant part, and J = Jastrow factor

• The backflow wave-function is:

$$\Psi_{BF} = e^{J(\mathbf{R})} \Psi_S(\mathbf{X})$$
, $\mathbf{x}_i = \mathbf{r}_i + \mathbf{\xi}_i$

ξ_i = backflow displacement of i-th electron

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Application and results

Conclusions

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Generalized backflow transformation:

Consider set of "preferred directions" seen by each electron



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Application and results

Conclusions

The backflow transformation

Generalized backflow transformation:

Write most general (isotropic) vector field in terms of such vectors

$$\boldsymbol{\xi}_{i} = \boldsymbol{\xi}_{i}^{e-e} + \boldsymbol{\xi}_{i}^{e-N} + \boldsymbol{\xi}_{i}^{e-e-N}$$

$$\boldsymbol{\xi}_{i}^{e-e} = \sum_{j \neq i}^{N} \boldsymbol{\eta}(r_{ij}) \boldsymbol{r}_{ij}$$

$$\boldsymbol{\xi}_{i}^{e-e-N} = \sum_{I}^{N} \sum_{j \neq i}^{N_{ion}} \left[\boldsymbol{\Phi}(r_{ij}, r_{iI}, r_{jI}) \boldsymbol{r}_{ij} + \boldsymbol{\Theta}(r_{ij}, r_{iI}, r_{jI}) \boldsymbol{r}_{iI} \right]$$
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Conclusions

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The backflow transformation

Generalized backflow transformation:

- Use power expansions to parametrize the backflow functions
- Smoothly cut off the backflow functions at optimizable distances
- Apply cusp conditions (choice: backflow not to alter conditions applied by Jastrow and orbitals)
- All-electron atoms delicate as orbitals are to cancel divergencies of potential; this must not be modified

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Application and results

Conclusions

The backflow transformation

Backflow complementary to Jastrow:





EFFECT OF A JASTROW FACTOR

COMPLEMENTARY EFFECT

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Application and results

Conclusions

The backflow transformation



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Changing the nodal surface:



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Changing the nodal surface:



CARBON ATOM, PROJECTED HF NODAL SURFACE



CARBON ATOM, PROJECTED BF NODAL SURFACE The backflow transformation

Application and results

Conclusions

The backflow transformation

Changing the nodal surface:



BERILLIUM ATOM, PROJECTED MD-BF NODAL SURFACE The backflow transformation

Application and results

Conclusions

The backflow transformation

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Application and results

Conclusions



DOTTED LINES: HOLZMANN, CEPERLEY, PIERLEONI, ESLER; PHYS REV E 68, 046707 (2003)

Application and results



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		Si diamor	nd 2x2x2	
Method	Wfn	Energy (au)	σ² (au)	
VMC	SJ	-7.87026(9)	0.591(2)	
	BF	-7.8875(3)	0.237(3)	
DMC	SJ	-7.8878(1)	-	
	BF	(running)	-	

		C diamond 2x2x2	
Method	Wfn	Energy (au)	σ² (au)
	SJ	-11.3708(2)	1.51(8)
VIVIC	BF	-11.3970(3)	0.897(8)
	SJ	-11.40717(8)	-
DIVIC	BF	-11.4141(3)	-

			Si betatin 2x2x2	
	Method	Wfn	Energy (au)	σ² (au)
	VMC	SJ	-62.0063(3)	0.74(2)
		BF	-62.180(5)	0.346(6)
		SJ	-62.175(1)	-
	DIVIC	BF	(running)	-

CRYSTALLINE SYSTEMS (PSEUDOPOTENTIALS)

The backflow transformation

Application and results

Conclusions

Application and results

- Implementation detail: electron-by-electron sampling is <u>much</u> <u>cheaper</u> than configuration-by-configuration sampling.
- Timing for fixed number of configs is increased.
- Lower errorbars provide large compensation.
- It is actually cheaper to get to a fixed errorbar with BF than SJ in the HEG and lithium atom!

Other systems: 2 to 8 times more expensive

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The closer VMC and DMC are for a given Ψ, the lower the DMC variance is, as:

$$\sigma_{DMC}^2 \sim |E_{VMC} - E_{DMC}|$$

Backflow generally lowers VMC more than DMC

 Hence DMC statistics are better: good for DMC even if the energy is unchanged
 The backflow transformation

Application and results

Conclusions

Application and results

- **Other systems successfully treated so far include:**
- Water molecule and water clusters (I. García de Gurtubay, UNPUBLISHED)
- Electron-hole systems (P. López Ríos, UNPUBLISHED)
- Neon and Neon+ (N. D. DRUMMOND ET AL, J. CHEM. PHYS. 124, 224104 (2006).)
- HEG in full (G. SPINK ET AL, UNPUBLISHED)

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Application and results

Conclusions

Conclusions

- Excellent results, good scaling properties, wide applicability
- Successfully combined with MD (not a substitute in all cases)
- Gets on well with orbital optimization too
- Combine with Pfaffians?
- BF-VMC is a powerful level of theory in many cases.
 BF-DMC statistics are hugely improved as a result

 Optimization method focused on nodes would help a lot [1]

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Conclusions

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