

An adaptive Langevin thermostat for non-equilibrium molecular dynamics simulations

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21st May 2008

Outline

Molecular Dynamics

Thermostats

Adaptive Langevin thermostat

Conclusions and Further Work

Why do we do MD simulations?

- To calculate observables - static, dynamic
- To see what happens

Changing Ensemble

- MD is nominally energy conserving - NVE ensemble
- Usually more interested in NVT or NPT ensemble
- Need temperature regulation - Thermostat

Thermostatted MD

- A thermostat alters the forces and/or velocities
- These alterations can be deterministic or stochastic

Non-equilibrium MD

Non-equilibrium simulations or bad equilibrium simulations generate heat:

- Non-Hamiltonian
- QM/MM with discontinuous force calculation
- $F \neq -\nabla U$

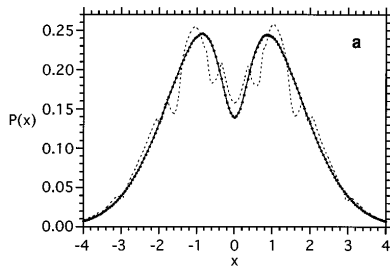
Here we can only hope a thermostat gives the correct average temperature.

Deterministic	Stochastic
<p>Rescaling Rescale velocities such that</p> $E_k = \frac{1}{2} N_{\text{dof}} k_B T$	<p>Andersen Every M steps choose a particle and reassign its velocity from Maxwellian distribution</p>
<p>Nosé-Hoover Hoover, PhysRevA 31 1695</p> $\dot{r}_i = \frac{p_i}{m_i}$ $\dot{p}_i = F_i - \xi p_i$ $Q\dot{\xi} = \sum_i \frac{p_i^2}{m_i} - N_{\text{dof}} k_B T$	<p>Langevin Quigley, Probert, JChemPhys 120(24) 11432</p> $\dot{r}_i = \frac{p_i}{m_i}$ $\dot{p}_i = F_i - \gamma p_i + \sqrt{\Gamma} \tilde{A}_i(t)$ $\Gamma = 2\gamma m_i k_B T \quad \langle \tilde{A}_i(t) \rangle = 0$ $\langle \tilde{A}_i(t) \tilde{A}_j(t') \rangle = \delta_{ij} \delta(t - t')$

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<p>Rescaling Temp. gradients Rescale velocities such that</p> $E_k = \frac{1}{2} N_{\text{dof}} k_B T$	<p>Andersen Unphysical (MC) Every M steps choose a particle and reassign its velocity from Maxwellian distribution</p>
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Nosé-Hoover Thermostat

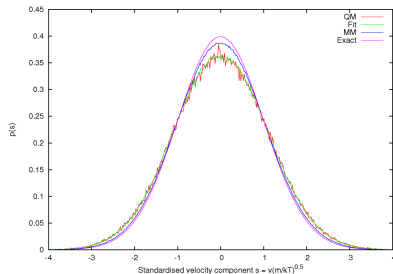
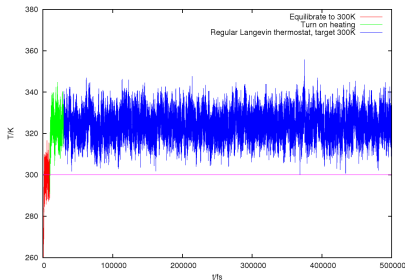
The Nosé-Hoover thermostat relies on chaotic trajectories



R.G. Winkler et. al., JChemPhys **102**(22) 9018

Langevin Thermostat

The Langevin thermostat cannot be used when there is heating - incorrect average temperature:



1728 Si atoms, with Stillinger-Weber potential

Requirements

- Feedback - Nosé-Hoover
- Stochastic - Langevin
- Modified Fluctuation-Dissipation relation

Kühne et al., PRL **98** 066401 (2007)

Equations of Motion

$$\dot{r}_i = \frac{p_i}{m_i}$$

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$$\langle f(t) \rangle_\tau = \frac{1}{\tau} \int_{-\infty}^t e^{\frac{t'-t}{\tau}} f(t') dt'$$

Choose β such that oscillations in s , and so $\langle T_k \rangle$, are critically damped (approximately)

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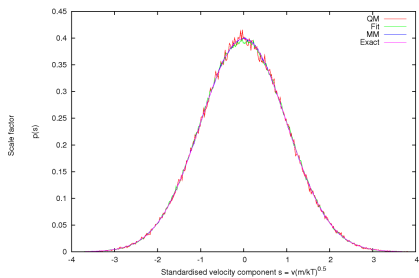
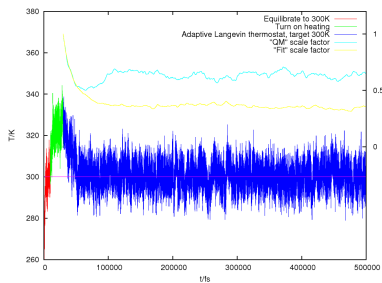
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Caveat: $s < 0$ makes no sense (\tilde{A} and $-\tilde{A}$ have same properties)

Results



Conclusions and Further Work

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Further Work:

- Recover Newtonian dynamics when temperature is OK
Leimkuhler et al., JChemPhys **128** 074105
- Re-derive using Fokker-Planck equation in extended phase-space

Thank-you for listening!

Any questions?

Nosé-Hoover-Langevin Thermostat?

$$\dot{r}_i = \frac{p_i}{m_i}$$

$$\dot{p}_i = F_i - \xi p_i$$

$$Q\dot{\xi} = \left[\sum_i \frac{p_i^2}{m_i} - Nk_B T \right] - \gamma Q\xi + \sqrt{\Gamma}\tilde{A}(t)$$

$$\Gamma = 2\gamma Qk_B T$$

- Gives canonical probability density in equilibrium simulations
- Has feedback to deal with non-equilibrium simulations
- When temperature has stabilised ξ decays - dynamics is more Newtonian