Orbital-dependent backflow transformations in quantum Monte Carlo

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VMC and DMC Optimization Wave functions

Variational and diffusion Monte Carlo

VMC and DMC are stochastic methods to solve the Schrödinger equation.

Aspects in common:

- Sample real-space configurations $\{ R \}$ using guiding wave function $\Psi_{\it T}(R)$
- Evaluate $E_L(\mathbf{R}) = \Psi_T^{-1}(\mathbf{R})\hat{H}(\mathbf{R})\Psi_T(\mathbf{R})$ to give variational estimate of ground-state energy E_0

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Variational and diffusion Monte Carlo

VMC works like this:

- {**R**} distributed according to $|\Psi_T(\mathbf{R})|^2$
- $E_{\rm VMC} = \langle \Psi_T | \hat{H} | \Psi_T \rangle$
- Quality of $E_{\rm VMC}$ depends on quality of Ψ_T

DMC works like this:

- Define Φ(**R**) as the exact solution of Schrödinger equation subject to: Φ(**R**) = 0 ⇔ Ψ_T(**R**) = 0
- $\{\mathbf{R}\}$ distributed according to $|\Phi(\mathbf{R})\Psi_T(\mathbf{R})|$
- $E_{\rm DMC} = \langle \Phi | \hat{H} | \Phi \rangle$
- Quality of E_{DMC} depends on quality of **nodes** of Ψ_T

VMC and DMC Optimization Wave functions

Wave function optimization

Optimization procedure:

- Generate set of VMC configurations $\{\mathbf{R}\}(\boldsymbol{\alpha})$
- Optimize target function with respect to $\pmb{\alpha}$ using fixed configurations

Target function can be:

- Energy, since VMC is variational
- Measure of spread of local energies (e.g., variance), since eigenstates of \hat{H} yield constant local energies

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VMC and DMC Optimization Wave functions

Slater-Jastrow wave function

For electronic systems one typically uses Slater determinants with HF orbitals times a Jastrow correlation factor:

Slater-Jastrow trial wave function

 $\Psi_{\rm SJ}(\mathbf{R}) = \exp\left[J(\mathbf{R})\right] \det\left[\phi_i(\mathbf{r}_j)\right]$

where $J(\mathbf{R})$ consists of e–e, e–n, e–e–n, ... terms which contain optimizable parameters

- J typically recovers 70 90% of correlation energy in VMC
- J does not change DMC result since it does not affect nodes of Ψ_T

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VMC and DMC Optimization Wave functions

Beyond Slater-Jastrow

One can use:

- Multi-determinant expansions
- Geminal-/Pfaffian-based wave functions
- Multi-Jastrow wave function
- Backflow transformation:

Slater-Jastrow-backflow trial wave function

 $\Psi_{\rm BF}(\mathbf{R}) = \exp[J(\mathbf{R})] \det \{\phi_i[\mathbf{x}_j(\mathbf{R})]\}$

where $\mathbf{x}_j(\mathbf{R}) = \mathbf{r}_j + \boldsymbol{\xi}_j(\mathbf{R})$, and $\boldsymbol{\xi}$ consists of e-e, e-n, e-e-n, terms containing optimizable parameters

- ... terms containing optimizable parameters.
 - Can be thought of as a correlated orbital correction method
 - $\bullet\,$ In VMC recovers 50% of the correlation energy remaining at SJ
 - Same accuracy as system size increases (fixed # parameters)

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Definition of ODBF wave function

Could use different backflow transformations for different orbitals:

ODBF trial wave function

$$\Psi_{\text{ODBF}}(\mathbf{R}) = \exp[J(\mathbf{R})] \det \left\{ \phi_i \left[\mathbf{x}_j^{t_i}(\mathbf{R}) \right] \right\}$$

where $\mathbf{x}_{j}^{t_{i}}(\mathbf{R}) = \mathbf{r}_{j} + \boldsymbol{\xi}_{j}^{t_{i}}(\mathbf{R})$

- Potentially very good since different orbitals in principle require different corrections
- Questions to answer:
 - Is it worth it? (it'd better be!)
 - Does it break size-extensivity?

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For a paramegnetic HEG consisting of 54 electrons at $r_s = 2$:

Ψ_T	<i>E</i> (a.u.)	V (a.u.)
$\Psi_{\rm HF}$	0.01876736	
Ψ_{SJ}	-0.01245(1)	0.148(1)
Ψ_{BF}	-0.013801(7)	0.0344(5)
$\Psi_{ m OBDF}$ (by k^2)	-0.013815(6)	0.0346(6)

Pablo López Ríos Orbital-dependent backflow

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For an all-electron carbon atom (single-determinant):

Ψ_T	<i>E</i> (a.u.)	V (a.u.)	CE (%)
$\Psi_{ m HF}$	-37.688619		0.0
Ψ_{SJ}	-37.8086(6)	0.139(1)	76.7(4)
Ψ_{BF}	-37.8286(4)	0.067(2)	89.5(3)
Ψ_{OBDF} (by l)	-37.8301(3)	0.0561(8)	90.5(2)
Ψ_{OBDF} (by n)	-37.8312(3)	0.053(1)	91.2(2)
$\Psi_{ ext{OBDF}}$ (by n,l)	-37.8330(3)	0.0451(9)	92.3(2)
Exact	-37.8450		100.0

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For an all-electron neon atom (single-determinant):

Ψ_T	<i>E</i> (a.u.)	V (a.u.)	CE (%)
$\Psi_{\rm HF}$	-128.547098		0.0
Ψ_{SJ}	-128.8975(6)	1.048(9)	89.7(2)
Ψ_{BF}	-128.9224(4)	0.341(2)	96.1(1)
Ψ_{OBDF} (by l)	-128.9263(3)	0.247(4)	97.11(8)
Ψ_{OBDF} (by n)	-128.9269(3)	0.239(3)	97.26(8)
Ψ_{OBDF} (by n, l)	-128.9278(3)	0.221(2)	97.49(8)
Exact	-128.9376		100.0

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The freedom to form linear combinations of orbitals

Replacing the set of orbitals by linear combinations of them leaves the wave function unchanged up to a multiplicative factor:

Determinant property

$$\begin{vmatrix} \phi_1(\mathbf{r}_1) & \phi_2(\mathbf{r}_1) & \phi_3(\mathbf{r}_1) \\ \phi_1(\mathbf{r}_2) & \phi_2(\mathbf{r}_2) & \phi_3(\mathbf{r}_2) \\ \phi_1(\mathbf{r}_3) & \phi_2(\mathbf{r}_3) & \phi_3(\mathbf{r}_3) \end{vmatrix} \propto \begin{vmatrix} \{\phi_1 + k\phi_2\}(\mathbf{r}_1) & \phi_2(\mathbf{r}_1) & \phi_3(\mathbf{r}_1) \\ \{\phi_1 + k\phi_2\}(\mathbf{r}_2) & \phi_2(\mathbf{r}_2) & \phi_3(\mathbf{r}_2) \\ \{\phi_1 + k\phi_2\}(\mathbf{r}_3) & \phi_2(\mathbf{r}_3) & \phi_3(\mathbf{r}_3) \end{vmatrix}$$

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The freedom to form linear combinations of orbitals

However if one applies orbital-dependent backflow to both sides of the previous equation the results differ:

This means that we have additional variational freedom which should be exploited.

Orbital mixing destroys properties of original orbitals (eigenvalues, quantum numbers, etc.), so:

- How do we assign different backflow transformations to the new orbitals?
- Can we still use fewer transformations than there are orbitals?
- Can the mixing be sensibly restricted? E.g., can we only mix same-*n* orbitals and get good results?

I don't know the answers, but I suspect that orbital mixing is more important in the HEG than in the atoms.

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Pablo López Ríos Orbital-dependent backflow



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