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The quantum Monte Carlo method can determine the total energy of an assembly of quantum particles to a high degree of accuracy.

In variational Monte Carlo (VMC), a trial wave function $\Psi_T(\mathbf{R})$ is used to evaluate an upper bound to the exact ground-state energy E_{ϱ} .

$$E_{VMC} = \sum_{i=1}^{M} \frac{\hat{H}(\boldsymbol{R}_{i}) \boldsymbol{\Psi}_{T}(\boldsymbol{R}_{i})}{\boldsymbol{\Psi}_{T}(\boldsymbol{R}_{i})} \simeq \frac{\int \left|\boldsymbol{\Psi}_{T}(\boldsymbol{R})\right|^{2} \frac{\hat{H}(\boldsymbol{R}) \boldsymbol{\Psi}_{T}(\boldsymbol{R})}{\boldsymbol{\Psi}_{T}(\boldsymbol{R})} d \boldsymbol{R}}{\int \left|\boldsymbol{\Psi}_{T}(\boldsymbol{R})\right|^{2} d \boldsymbol{R}} \ge E_{0}$$

- \diamond In diffusion Monte Carlo (DMC), imaginary-time propagation is used to project out the higher-energy components of $\Psi_{\tau}(\mathbf{R})$.
- igle For fermions, the *fixed-node approximation* is required. It is equivalent to constraining the *nodal surface* of the DMC wave function to equal that of $\Psi_{\tau}(\mathbf{R})$.

ESDG - 6/06/07 2/26

The following are typical fermionic trial wave functions:

(Slater or Hartree-Fock type)

where $\Psi_{S}(\mathbf{R})$ is a Slater determinant (or a multi-determinant expansion) of suitable one-particle orbitals. With a single determinant, there are **no correlation effects**.

(Slater-Jastrow type)

where $e^{J(\mathbf{R})}$ is a Jastrow correlation factor containing optimizable parameters.

(Slater-Jastrow-backflow type)

where X(R) is a set of collective coordinates containing optimizable parameters.

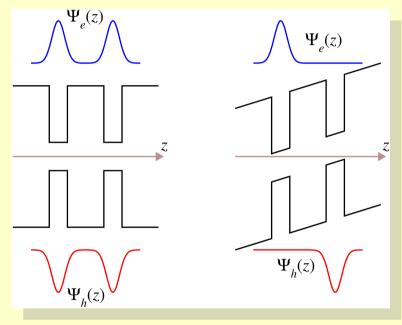
Key points:

- \diamond The quality of $\Psi_T(\mathbf{R})$, and in particular of its nodal surface, is paramount for the accuracy of the results.
- \diamond The Jastrow factor is the most important tool to improve upon $\Psi_{S}(\mathbf{R})$, but it does not change its nodes.
- There are ways to modify the nodal surface of $\Psi_{S}(\mathbf{R})$: orbital optimization, backflow transformations, multideterminant expansions, multi-Pfaffian wfns, etc. This is a very active area of research.

ESDG - 6/06/07 4/26

The electron-hole system is a model for excited semiconductors.

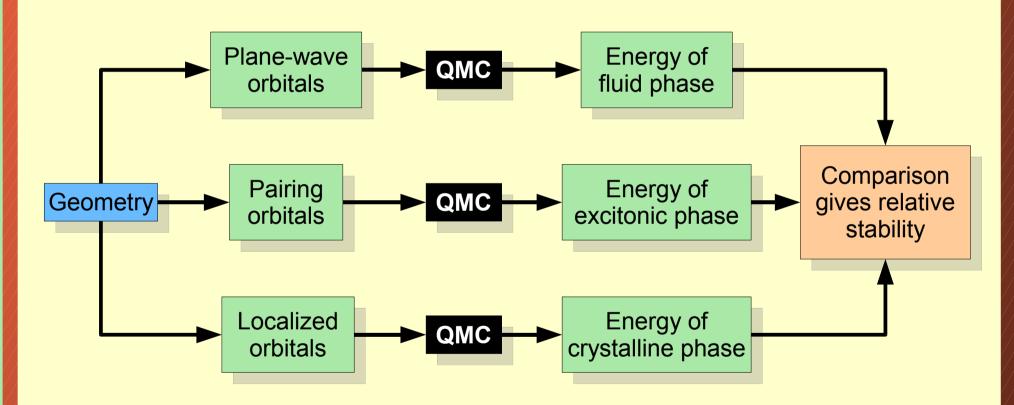
- ◆ It is the simplest model system after the homogeneous electron gas, yet its phase diagram remains largely unknown.
- Some electron-hole systems, such as the two-dimensional bilayer, can be recreated experimentally and display very interesting properties.



◆ The aim of this work is to determine the phase diagram of electron-hole systems using QMC, to a greater degree of accuracy than previous studies.

ESDG - 6/06/07 5/26

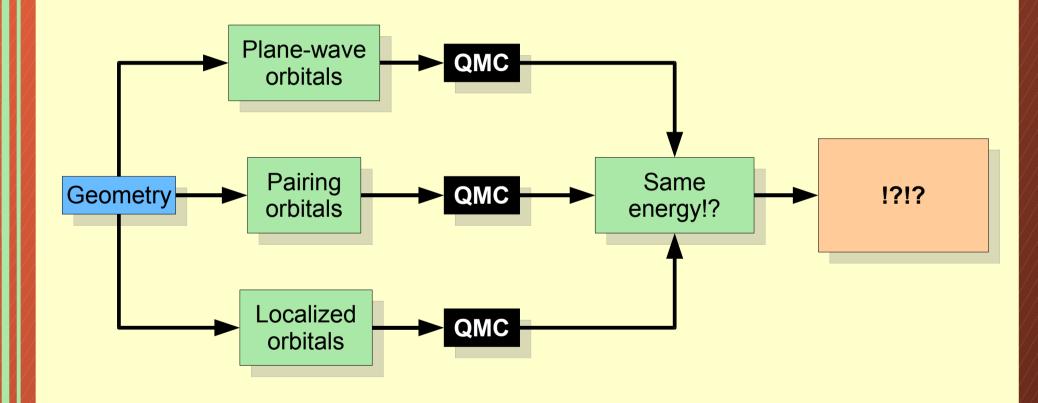
Previous QMC studies have made use of the following scheme:



"Nodal argument": different orbitals give different nodes give different energies. Hence there is a correspondence between phase/orbitals (input) and energies (output).

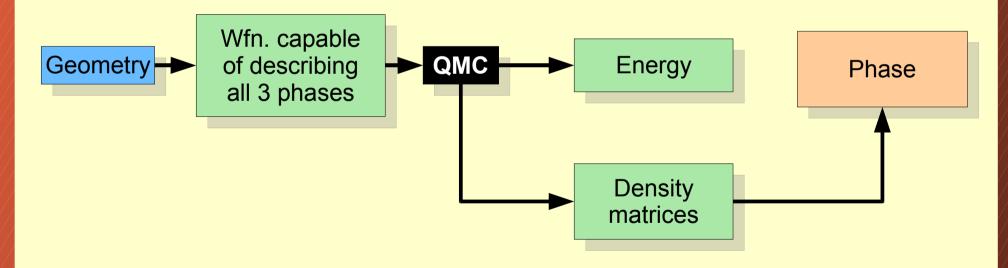
ESDG - 6/06/07 6/26

However, in the limit of a perfect wave function:



ESDG - 6/06/07 7/26

An appropriate approach is:



And on the plus side, this requires fewer calculations.

ESDG - 6/06/07 8/26

The wave functions are formed using products of Slater determinants:

$$D[\phi] = \begin{vmatrix} \phi_{11} & \phi_{12} & \cdots & \phi_{1N} \\ \phi_{21} & \phi_{22} & \cdots & \phi_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ \phi_{NI} & \phi_{N2} & \cdots & \phi_{NN} \end{vmatrix}$$

with:

$$\phi_{ij} = \phi_L(\boldsymbol{e}_i - \boldsymbol{h}_j)$$

Crystal:

$$\phi_{ij} = e^{i \mathbf{k}_j \cdot \mathbf{r}_i}$$

$$\phi_{ij} = e^{i \mathbf{k}_j \cdot \mathbf{r}_i}$$

$$\phi_{ij} = \phi_C(\mathbf{r}_i - \mathbf{R}_j)$$

$$\Psi_{S} = D_{e \uparrow h \downarrow} D_{e \downarrow h \uparrow}$$

$$\Psi_{S} = D_{e\uparrow} D_{e\downarrow} D_{h\uparrow} D_{h\downarrow}$$

$$\rightarrow$$

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with:

$$\boldsymbol{\phi}_{ij} = \boldsymbol{\phi}_L(\boldsymbol{e}_i - \boldsymbol{h}_j)$$

$$\phi_{ij} = \sum_{l=0}^{N} c_l e^{i \mathbf{k}_l \cdot (\mathbf{e}_i - \mathbf{h}_j)}$$

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with:

$$\phi_{ij} = \sum_{p=0}^{P} c_p e^{i \mathbf{k}_p \cdot (\mathbf{e}_i - \mathbf{h}_j)} + \phi_L(\mathbf{e}_i - \mathbf{h}_j)$$

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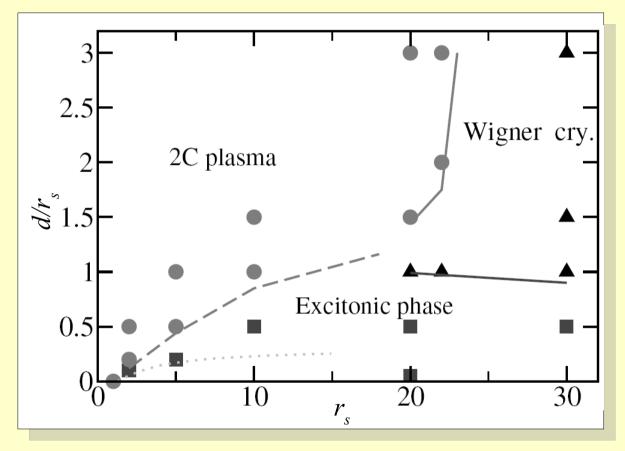
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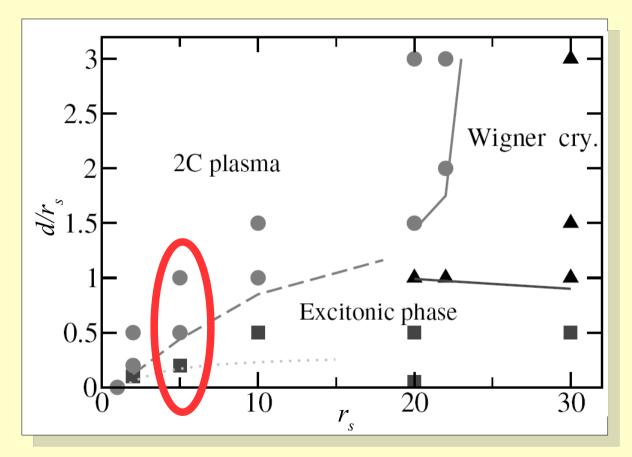
The phase diagram of the symmetric 2D electron-hole bilayer has already been studied.



S. de Palo et al, Phys. Rev. Lett. 88, 206401 (2002)

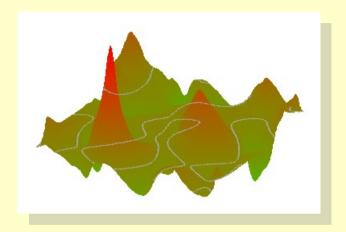
ESDG - 6/06/07 12/26

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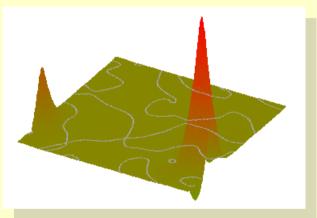
S. de Palo et al, Phys. Rev. Lett. 88, 206401 (2002)

ESDG - 6/06/07 13/26



Fluid, E=-0.17879(3) a.u.

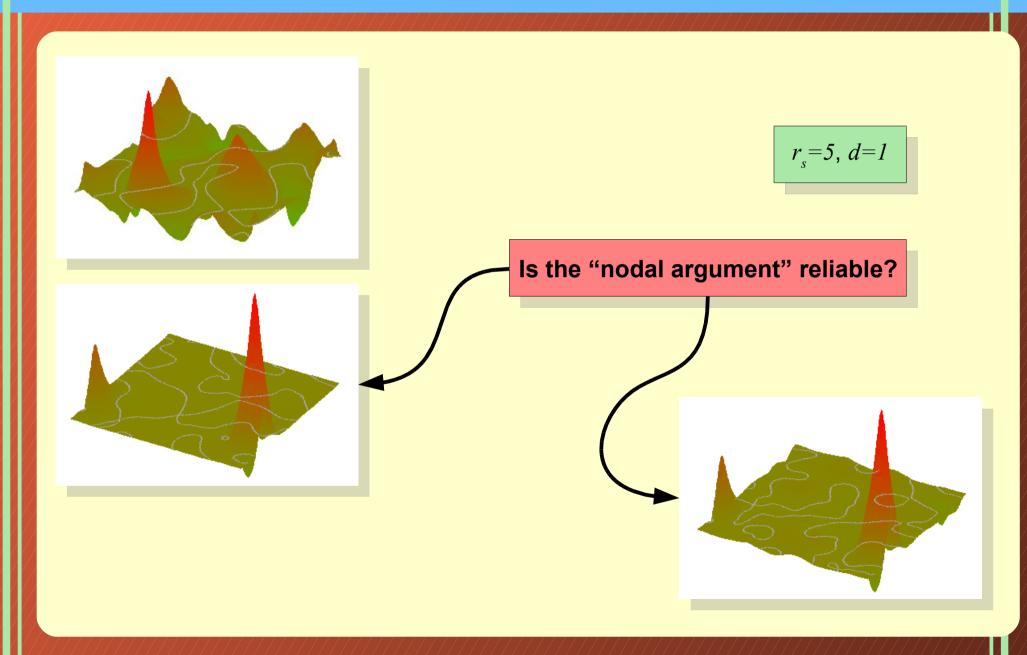
$$r_{s}=5, d=1$$



Exponential, E=-0.18937(3) a.u.

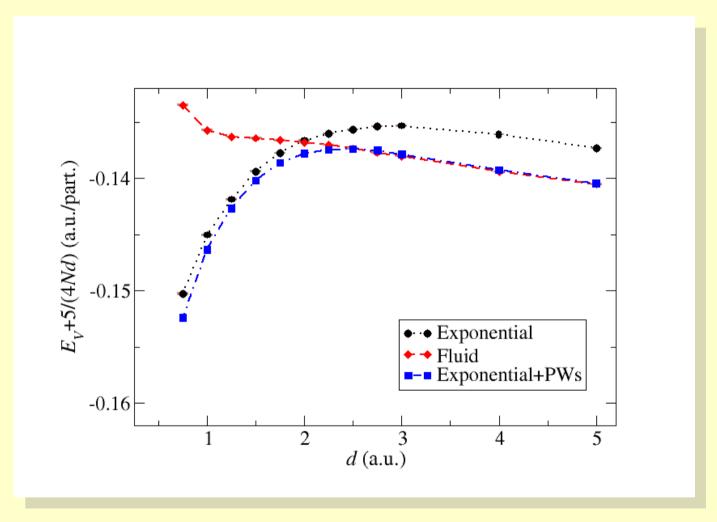


Exponential+PWs, E=-0.18979(3) a.u.



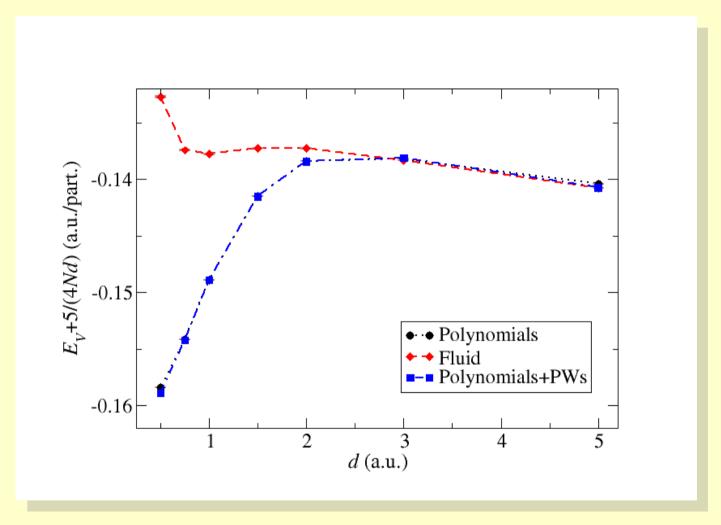
ESDG - 6/06/07

15/26



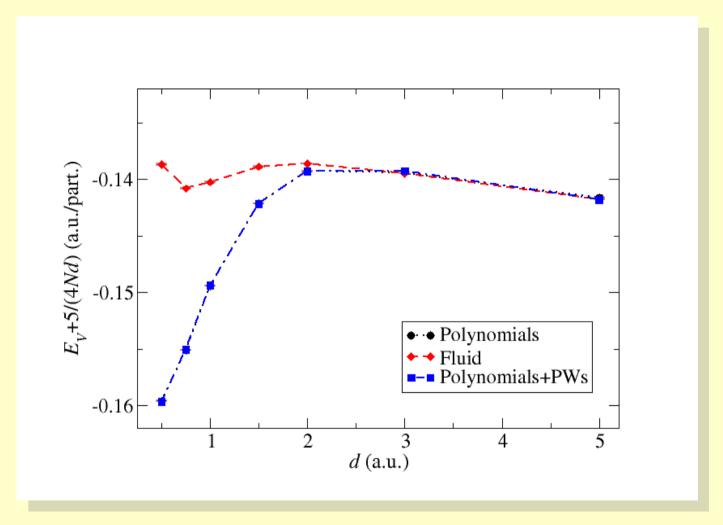
VMC results [variance minimization]

ESDG - 6/06/07 16/26



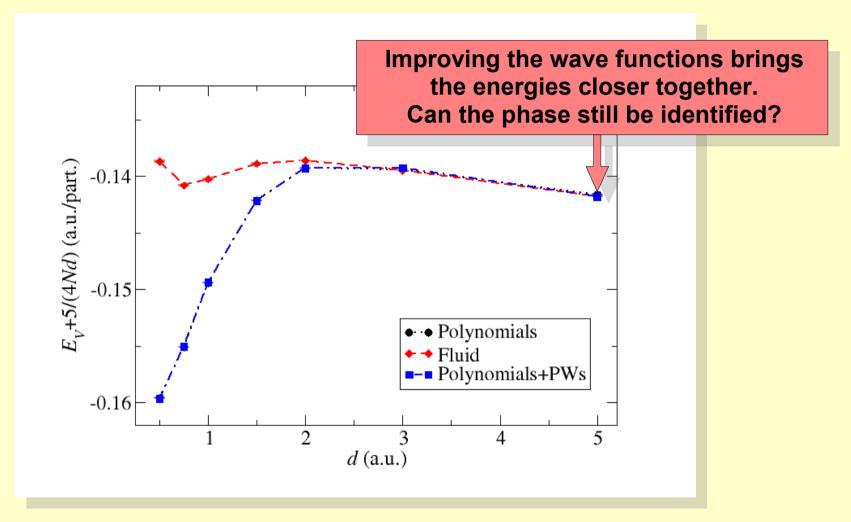
VMC results [energy minimization]

ESDG - 6/06/07 17/26



BF-VMC results [energy minimization]

ESDG - 6/06/07 18/26

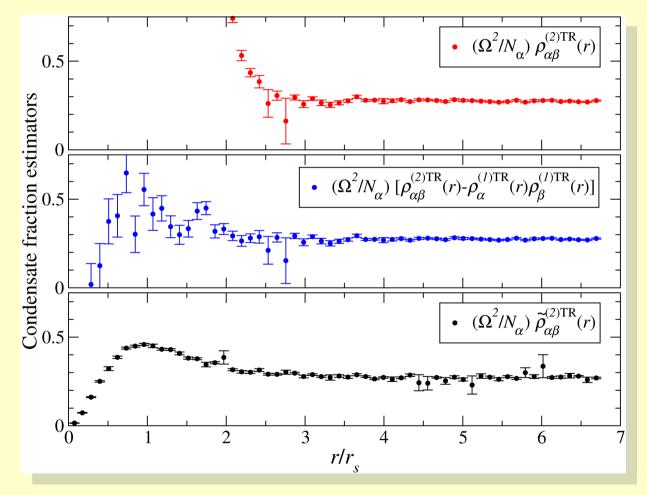


BF-VMC results [energy minimization]

ESDG - 6/06/07 19/26

The asymptotic behaviour of the two-body density matrix can be used to determine

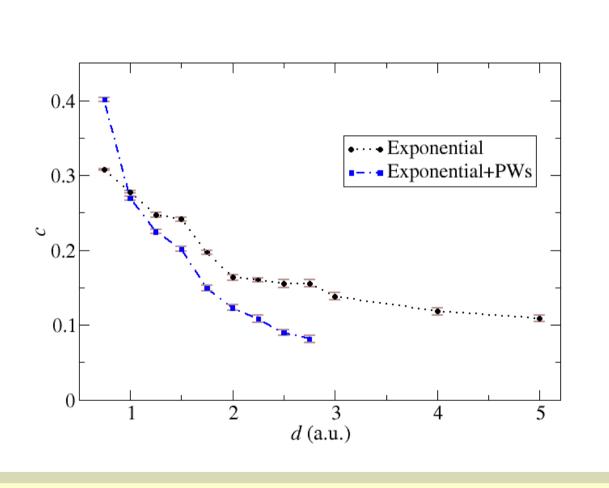
the phase:



ESDG - 6/06/07 20/26

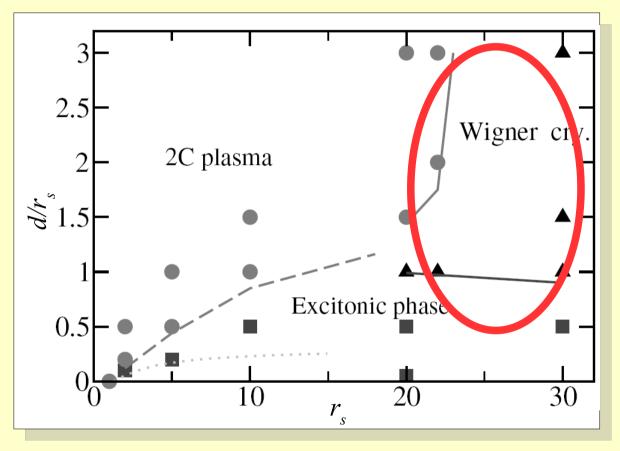
The asymptotic behaviour of the two-body density matrix can be used to determine

the phase:



ESDG - 6/06/07 21/26

How about the crystalline phase?



S. de Palo *et al*, Phys. Rev. Lett. **88**, 206401 (2002)

ESDG - 6/06/07 22/26

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with:

$$\phi_{ij} = \sum_{p=0}^{P} c_p e^{i \mathbf{k}_p \cdot (\mathbf{e}_i - \mathbf{h}_j)} + \phi_L(\mathbf{e}_i - \mathbf{h}_j)$$

$$\Psi_S = D_{e \uparrow h \downarrow} D_{e \downarrow h \uparrow}$$

$$\Psi_{S} = D_{e \uparrow h \downarrow} D_{e \downarrow h \uparrow}$$

Fluid:

$$\phi_{ij} = \phi_C(\boldsymbol{r}_i - \boldsymbol{R}_j)$$

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with:

Pairing:

$$\phi_{ij} = \sum_{p=0}^{P} c_p e^{i \mathbf{k}_p \cdot (\mathbf{e}_i - \mathbf{h}_j)} + \phi_L(\mathbf{e}_i - \mathbf{h}_j)$$

$$\Psi_S = D_{e \uparrow h \downarrow} D_{e \downarrow h \uparrow}$$

$$\Psi_{\scriptscriptstyle S} = D_{e^{\uparrow}h^{\downarrow}} D_{e^{\downarrow}h^{\uparrow}}$$

Fluid:

Crystal:
$$\psi_{ij} = \sum_{q=0}^{Q} d_q e^{i \mathbf{k}_q \cdot (\mathbf{r}_i - \mathbf{R}_j)} + \phi_C(\mathbf{r}_i - \mathbf{R}_j)$$

$$\psi_S = D_{e\uparrow} D_{e\downarrow} D_{h\uparrow} D_{h\downarrow}$$

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$$\Gamma_{ij} = \phi(\boldsymbol{e}_i, \boldsymbol{h}_j) + \sum_{l=0}^{L} \psi(\boldsymbol{e}_i - \boldsymbol{R}_l) \psi(\boldsymbol{h}_j - \boldsymbol{R}_l)$$

with:

Pairing:

$$\phi_{ij} = \sum_{p=0}^{P} c_p e^{i \boldsymbol{k}_p \cdot (\boldsymbol{e}_i - \boldsymbol{h}_j)} + \phi_L(\boldsymbol{e}_i - \boldsymbol{h}_j)$$

Fluid:

$$\psi_{ij} = \sum_{q=0}^{Q} d_q e^{i \mathbf{k}_q \cdot (\mathbf{r}_i - \mathbf{R}_j)} + \phi_C(\mathbf{r}_i - \mathbf{R}_j)$$



$$\Psi_{S} = D_{e \uparrow h \downarrow} D_{e \downarrow h \uparrow}$$

Conclusions and further work:

- Improvements over previous QMC calculations are possible. Need a different approach involving:
 - More general wave functions.
 - Expectation values of density-matrix related objects.
- Robust wave-function optimization is a key element to get results efficiently.
- Real test: compare against experiment.

ESDG - 6/06/07 26/26