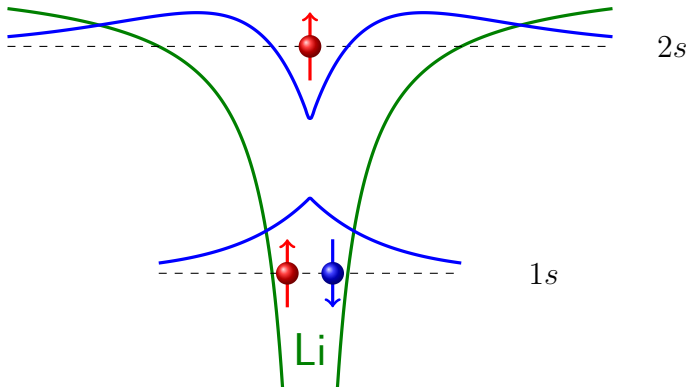


Pseudopotentials for cold atoms

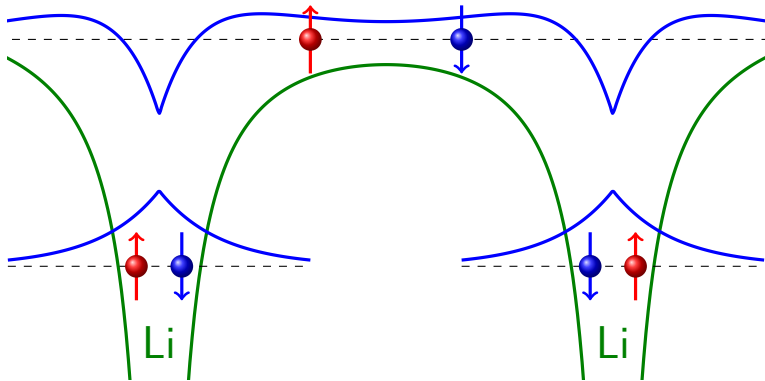
Pascal Bagnion, Gareth Conduit and Richard Needs

ESDG, 27th November 2013

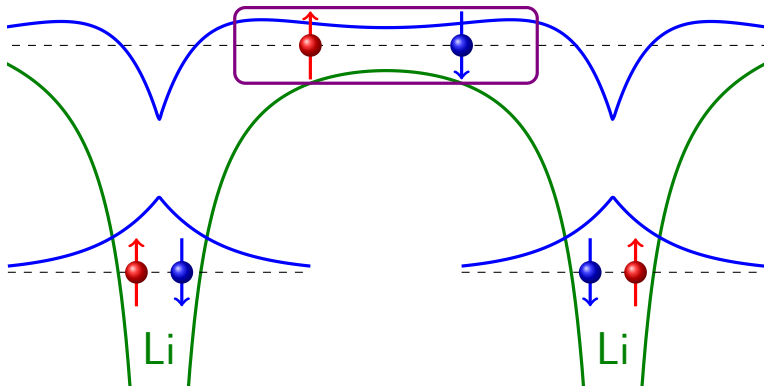
Pseudopotentials 101



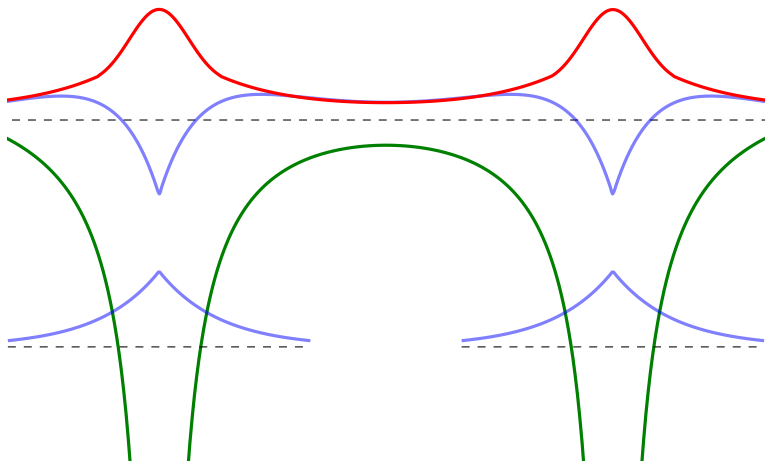
Pseudopotentials 101



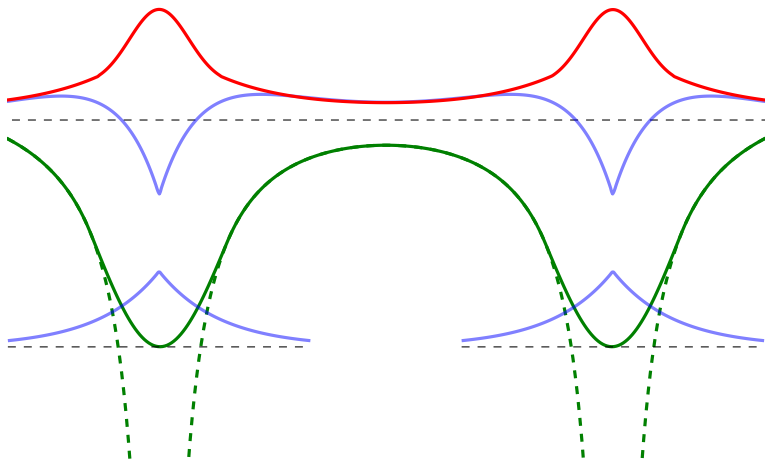
Pseudopotentials 101



Pseudopotentials 101

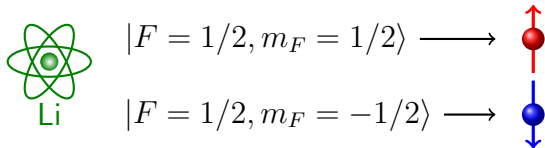


Pseudopotentials 101



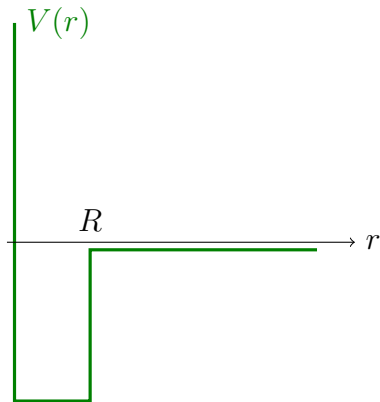
Why cold atoms?

Simulate complex many body systems that are often difficult to observe in the solid state.



The inter-atom interactions can be **tuned** by changing an external magnetic field.

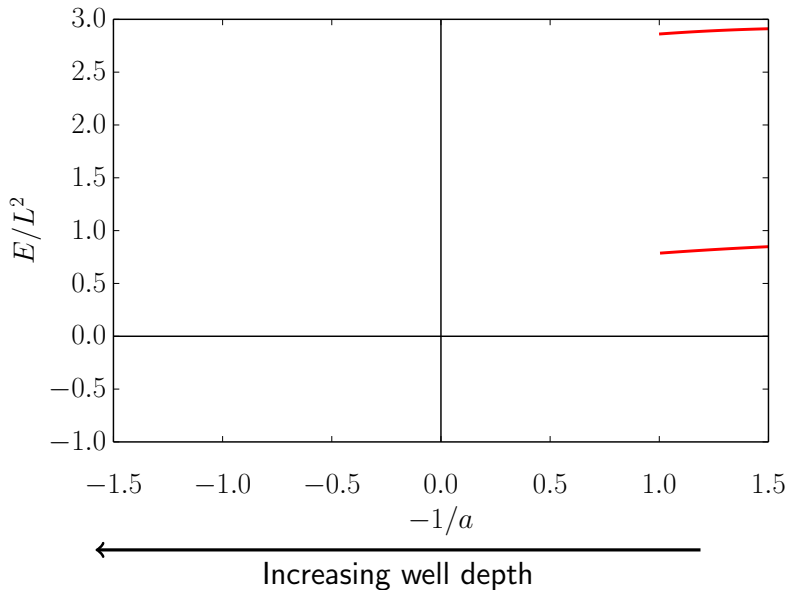
Inter-atomic potential



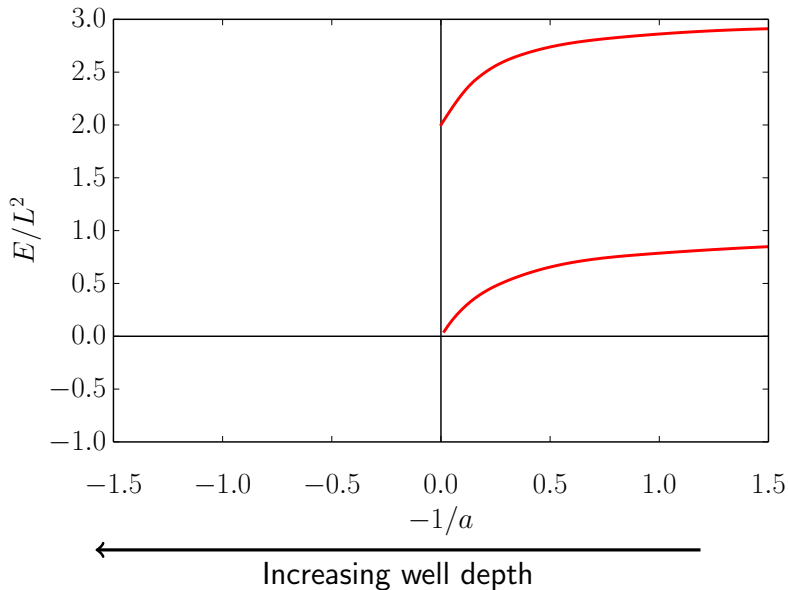
$$V(r) = \lim_{R \rightarrow 0} \begin{cases} -V_0 & r \leq R \\ 0 & \text{otherwise} \end{cases}$$

$$V_0 R^3 = \text{const.}$$

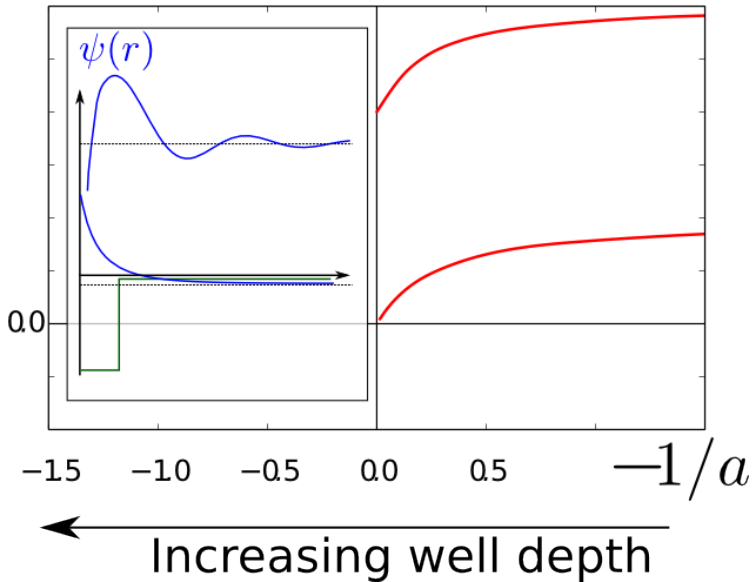
Inter-atomic potential



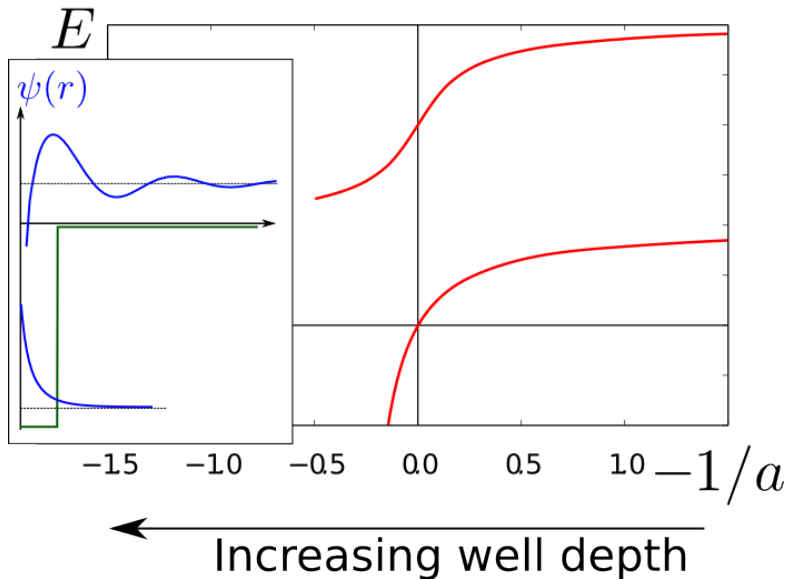
Inter-atomic potential



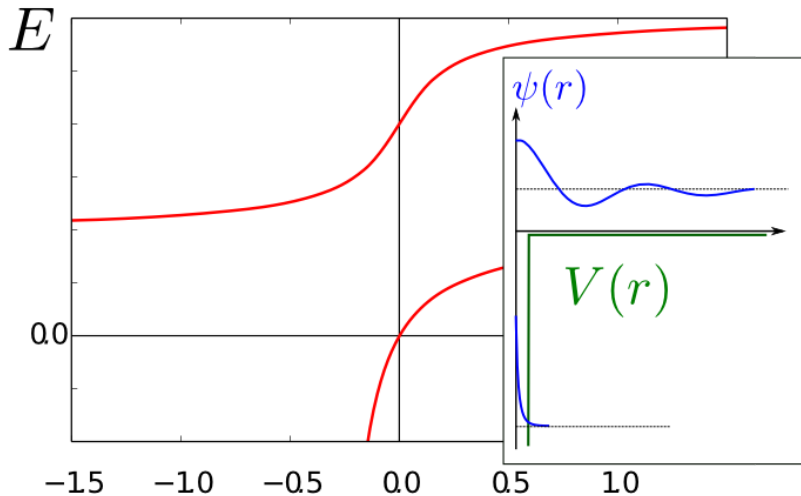
Inter-atomic potential



Inter-atomic potential

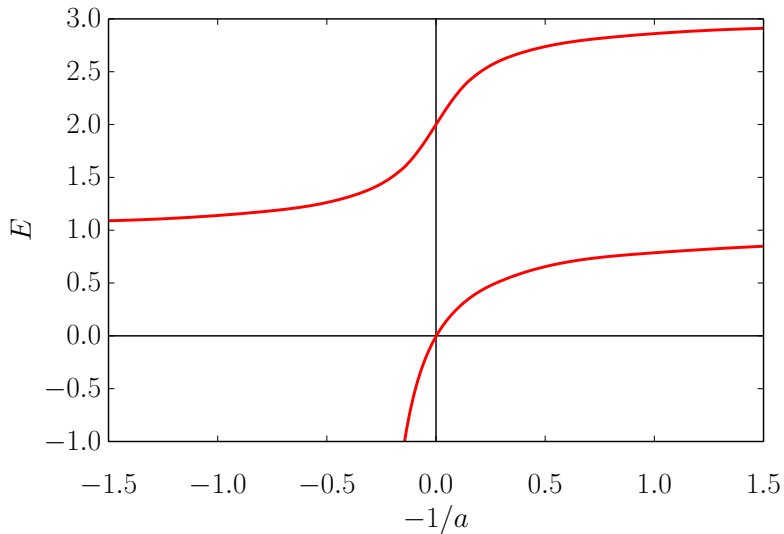


Inter-atomic potential

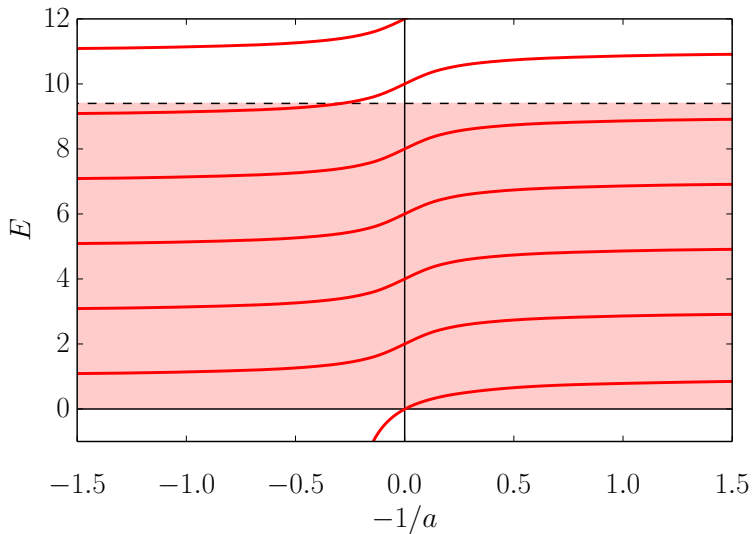


← Increasing well depth

Inter-atomic potential



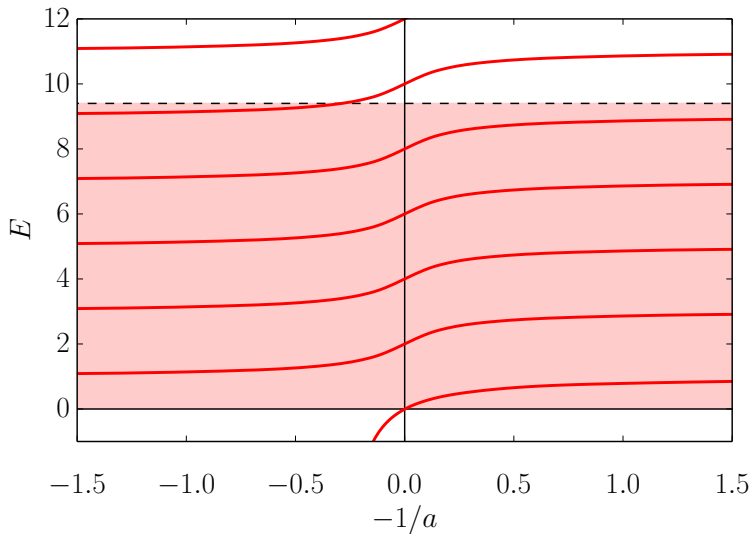
Pseudopotential requirements



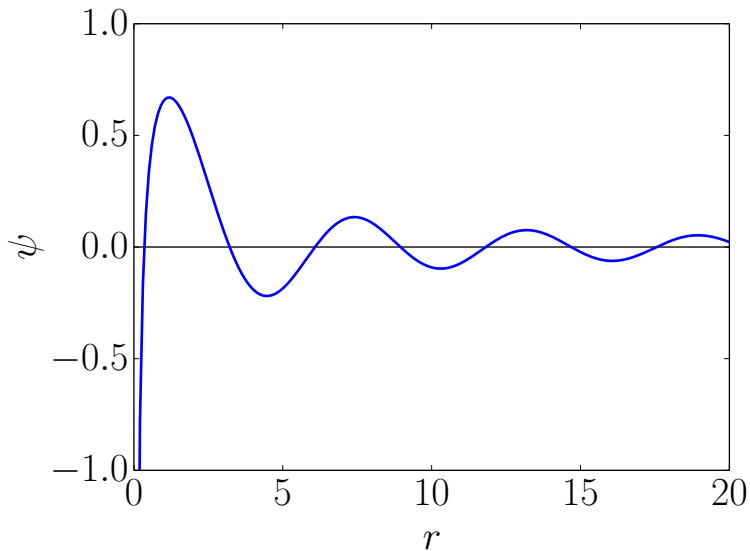
Building the pseudopotential

N. Trouiller and J.L Martins, *Efficient pseudopotentials for plane wave calculations*, PRB **43**, 1993 (1991)

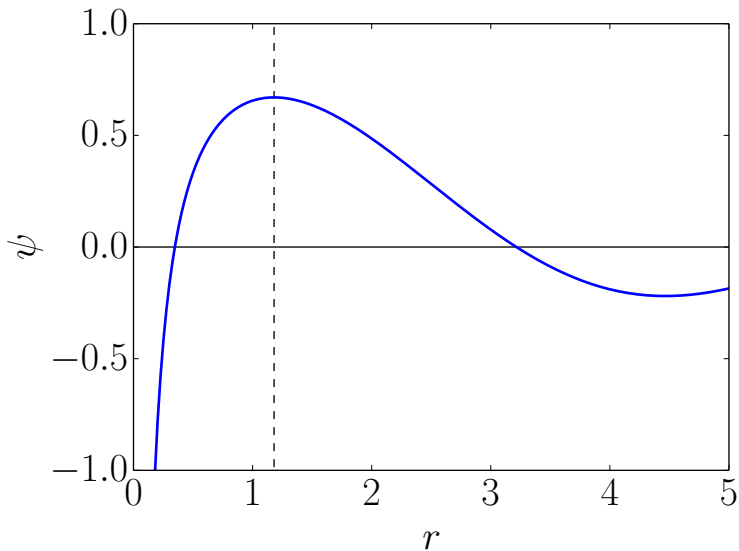
Pseudopotential requirements



Building the pseudo-wavefunction



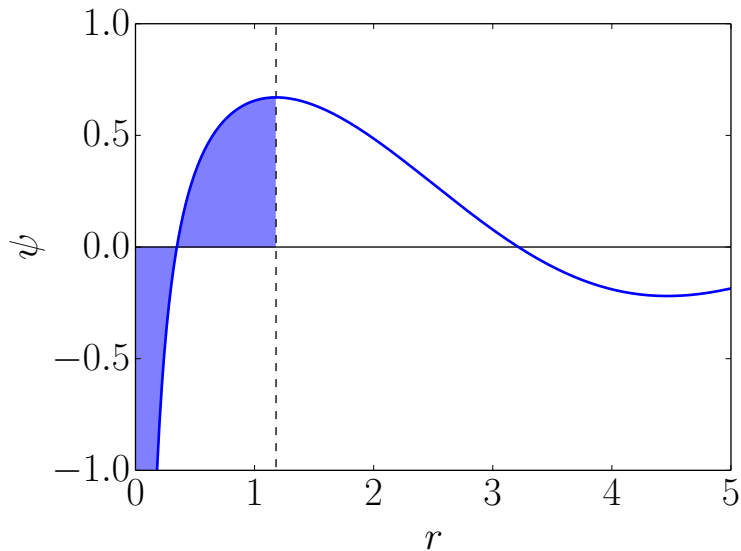
Building the pseudo-wavefunction



Norm conservation

$$\left. \frac{\partial \tan(\delta_{PP})}{\partial E} \right|_{r=c} = \left. \frac{\partial \tan(\delta_{\text{true}})}{\partial E} \right|_{r=c}$$

Norm conservation



Building the pseudo-wavefunction

- Norm conservation
- The pseudo-wavefunction must join smoothly onto the wavefunction at the cutoff: at least two derivatives must match.

Building the pseudo-wavefunction

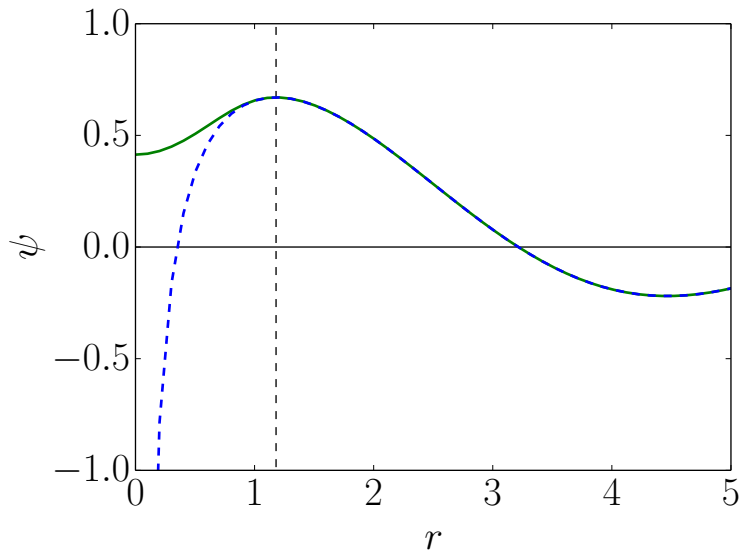
- Norm conservation
- The pseudo-wavefunction must join smoothly onto the wavefunction at the cutoff: at least two derivatives must match.

Functional form of the pseudo-wavefunction

$$\psi_{PP} = \begin{cases} e^{p(r)} & r \leq c \\ \psi_{\text{true}} & r > c \end{cases}$$

$$p(r) = a_0 + a_2 r^2 + a_4 r^4 + \sum_{i=6} a_i r^i$$

Building the pseudo-wavefunction

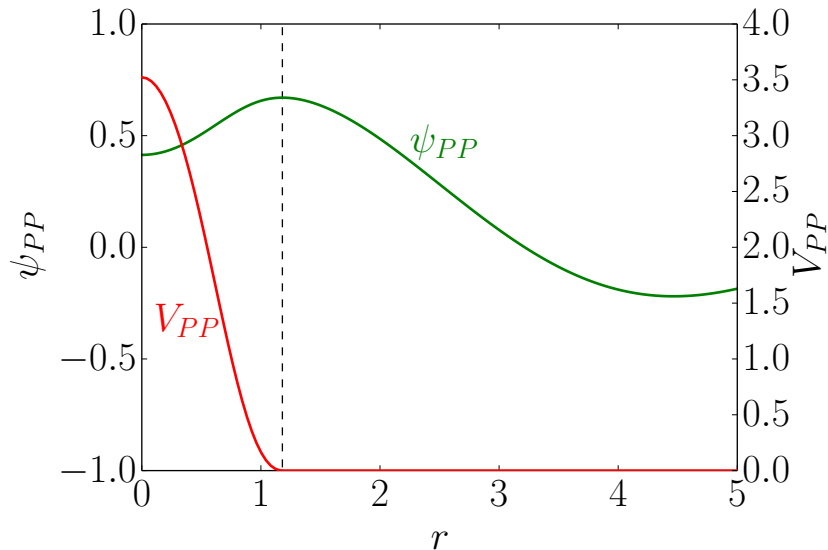


Building the pseudopotential

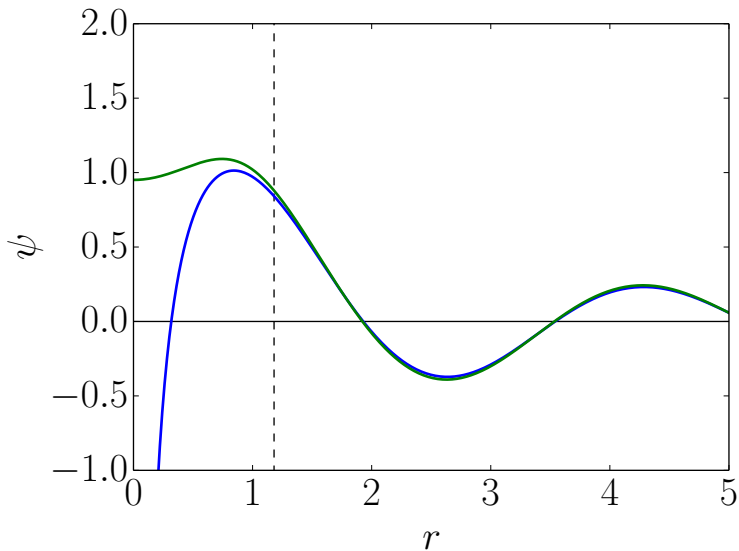
Invert the Schrödinger equation

$$V(r) = E + \frac{1}{2} \cdot \frac{\nabla^2 \psi}{\psi}$$

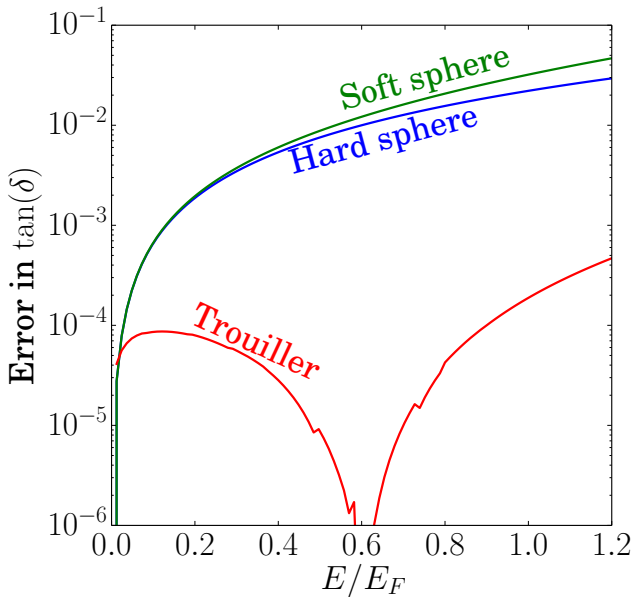
Building the pseudopotential



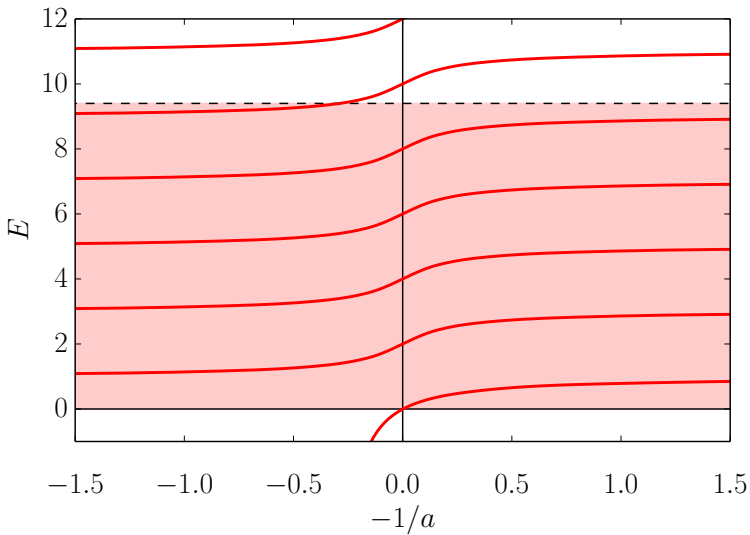
Breaking the pseudopotential: $E = 2E_F$



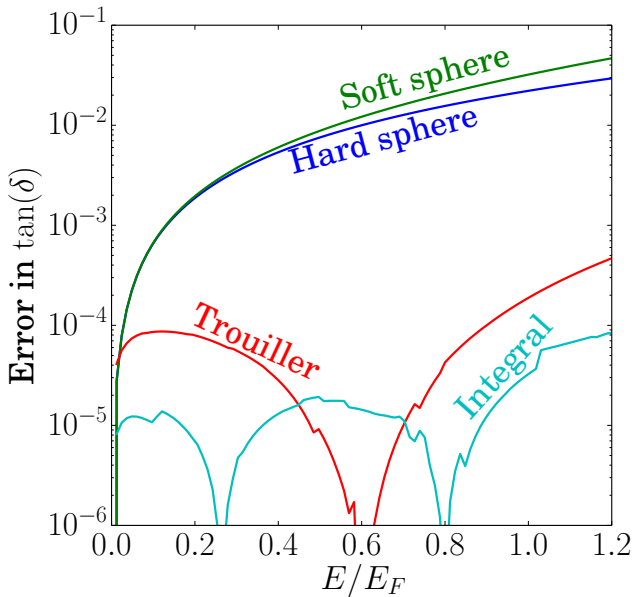
Phase shifts



Pseudopotentials++



Phase shifts



Ferromagnetism

We expect *itinerant ferromagnetism* to appear at sufficiently strong interactions.

Applications

Ferromagnetism

We expect **itinerant ferromagnetism** to appear at sufficiently strong interactions.

p-wave superconductivity

It is postulated that same-spin atoms feel a weak effective **attraction**.

Like-spin atoms may therefore form **Cooper pairs**.

Conclusion

- A good repulsive pseudopotential is necessary to model ultracold atomic gases at positive scattering length.

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- There is a lot of useful pseudopotential lore in the electronic structure literature.
- New pseudopotentials offer unprecedented opportunities to model exotic phenomena.