

Diffusion Monte-Carlo, nodal surfaces and pairing functions

Variational principle for Fermions

$$E_0 = \min_{\Psi} \left\{ \frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} \right\}$$

Variational principle for Fermions

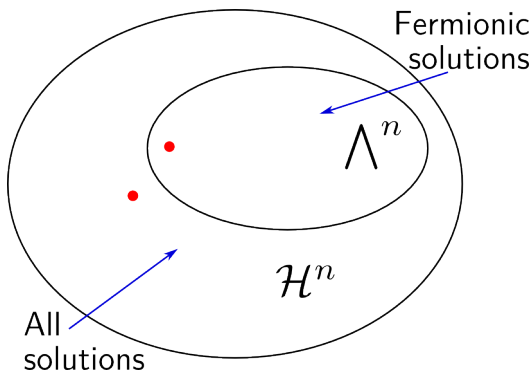
$$E_0 = \min_{\Psi} \left\{ \frac{\langle \Psi | \hat{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} \right\}$$

Subject to...

Ψ must remain anti-symmetric.

Variational principle for Fermions

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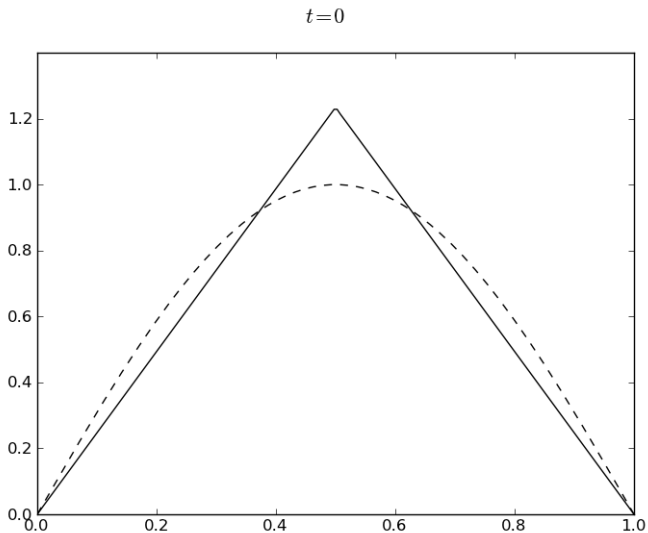


Diffusion Monte-Carlo

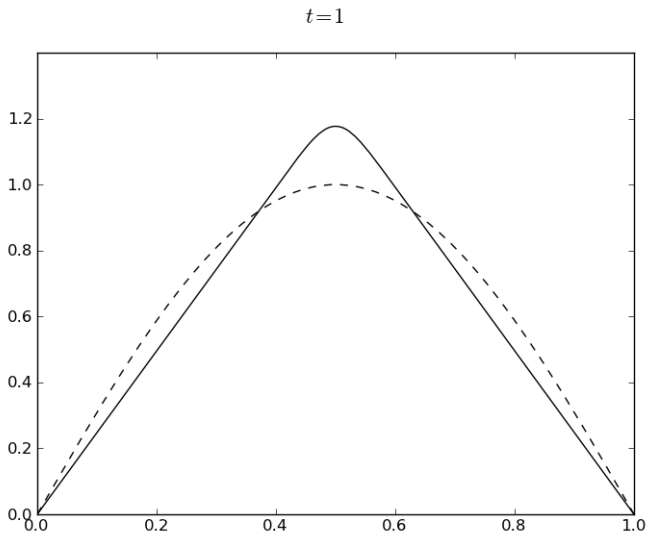
$$\lim_{t \rightarrow \infty} \left(\exp^{-\hat{H}t} \Psi_T \right) = \Psi_0$$

where Ψ_T is *any* n -electron wavefunction and Ψ_0 is the wavefunction that minimizes the energy in \mathcal{H}^n .

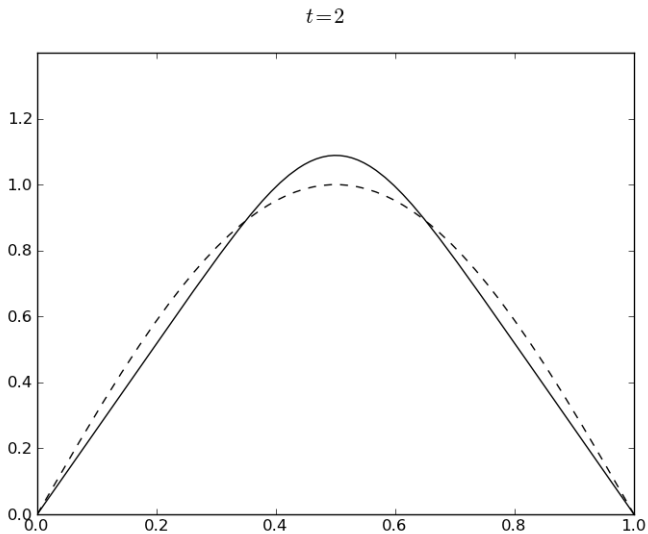
DMC for the particle in a box



DMC for the particle in a box

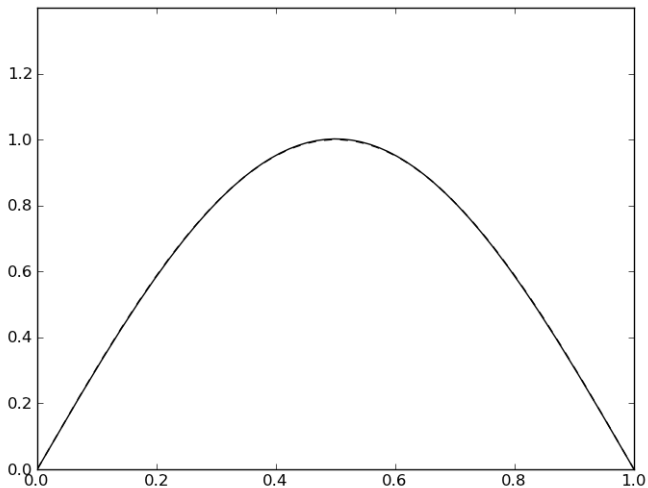


DMC for the particle in a box



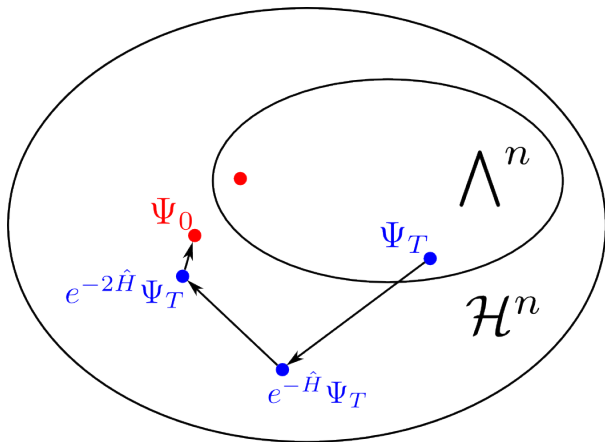
DMC for the particle in a box

$t=3$



Fermion sign problem

Repeated application of $\exp^{-\hat{H}}$ will eventually give the lowest-energy eigenfunction in \mathcal{H}^n .



Nodal surface

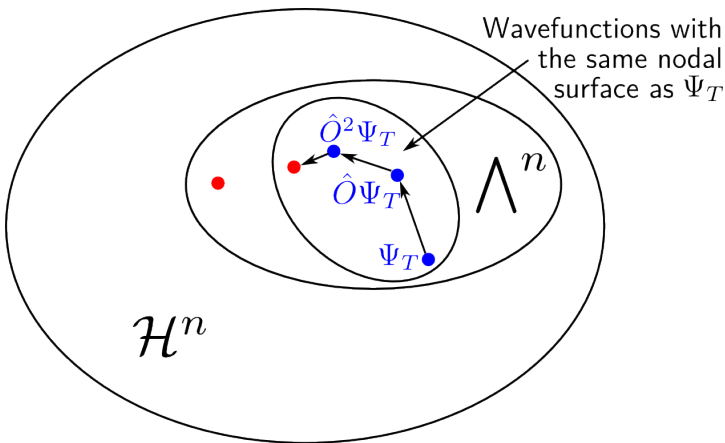
The Hamiltonian operator acts in the same manner on all particles.

Ensuring that the wavefunction is 0 when any two electrons coincide is sufficient to guarantee anti-symmetry.

The locus of points on which the wavefunction is 0 is called the **nodal surface**.

Fixed node approximation

DMC within the **fixed-node approximation** allows the optimization of a trial wavefunction without changing the nodal surface.



Trial functions

Traditionally, the Slater-Jastrow wavefunction is used as a trial wavefunction:

$$\Psi_T = e^J D$$

where e^J is positive everywhere and D is the Hartree-Fock wavefunction for the system.

New trial wavefunctions?

The nodal surface of the trial wavefunction can be improved by replacing the Slater determinant with some more complex anti-symmetric function Ψ_A :

$$\Psi_T = e^J \Psi_A$$

But what functional form can we use for Ψ_A ?

Beyond one-electron orbitals

Electrons in the real world do not exist in one-electron orbitals. Could we build **two-electron orbitals**?

Beyond one-electron orbitals

Electrons in the real world do not exist in one-electron orbitals. Could we build **two-electron orbitals**?

$$\Phi(\mathbf{r}, \bar{\mathbf{r}}) = \sum_{i,j} g_{ij} \phi_i(\mathbf{r}) \phi_j(\bar{\mathbf{r}})$$

These two-electron orbitals are called **geminals** in the literature.

From orbitals to wavefunction

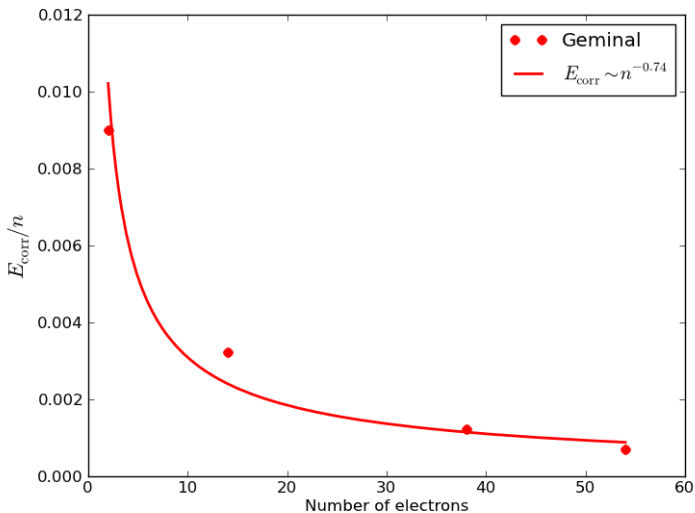
$$\Psi_A = \hat{\mathcal{A}}[\Phi(\mathbf{r}_1, \bar{\mathbf{r}}_1)\Phi(\mathbf{r}_2, \bar{\mathbf{r}}_2) \dots \Phi(\mathbf{r}_n, \bar{\mathbf{r}}_n)]$$

From orbitals to wavefunction

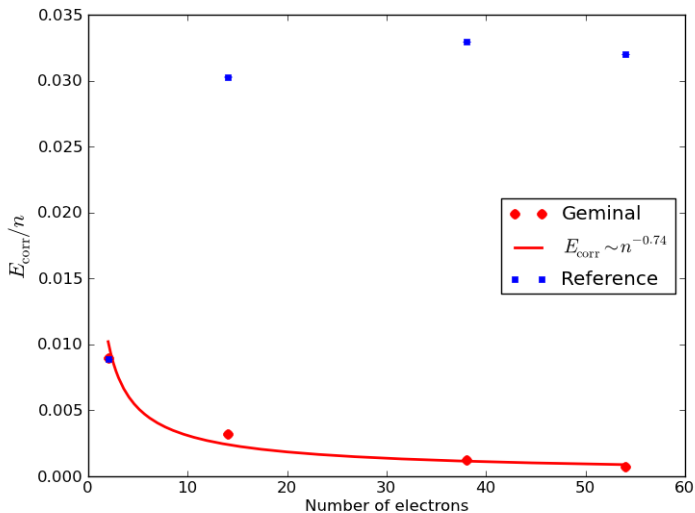
$$\Psi_A = \hat{\mathcal{A}}[\Phi(\mathbf{r}_1, \bar{\mathbf{r}}_1)\Phi(\mathbf{r}_2, \bar{\mathbf{r}}_2) \dots \Phi(\mathbf{r}_n, \bar{\mathbf{r}}_n)]$$

$$\Psi_A = \begin{vmatrix} \Phi(\mathbf{r}_1, \bar{\mathbf{r}}_1) & \Phi(\mathbf{r}_1, \bar{\mathbf{r}}_2) & \cdots & \Phi(\mathbf{r}_1, \bar{\mathbf{r}}_N) \\ \Phi(\mathbf{r}_2, \bar{\mathbf{r}}_1) & \Phi(\mathbf{r}_2, \bar{\mathbf{r}}_2) & \cdots & \Phi(\mathbf{r}_2, \bar{\mathbf{r}}_N) \\ \vdots & \vdots & \ddots & \vdots \\ \Phi(\mathbf{r}_N, \bar{\mathbf{r}}_1) & \Phi(\mathbf{r}_N, \bar{\mathbf{r}}_2) & \cdots & \Phi(\mathbf{r}_N, \bar{\mathbf{r}}_N) \end{vmatrix}$$

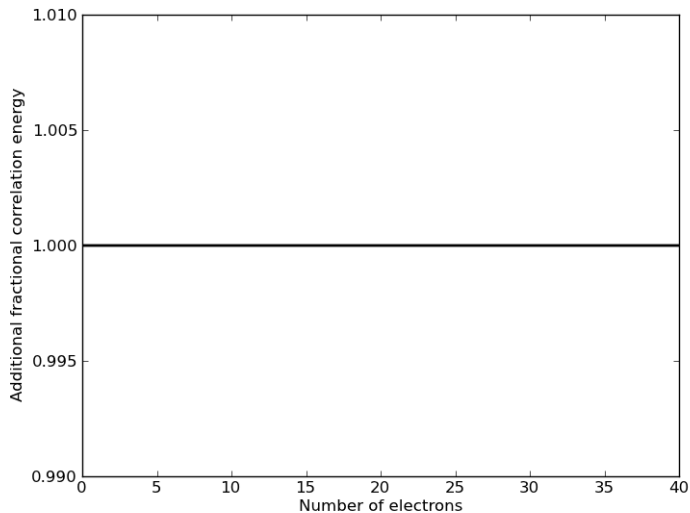
Results - Geminals alone



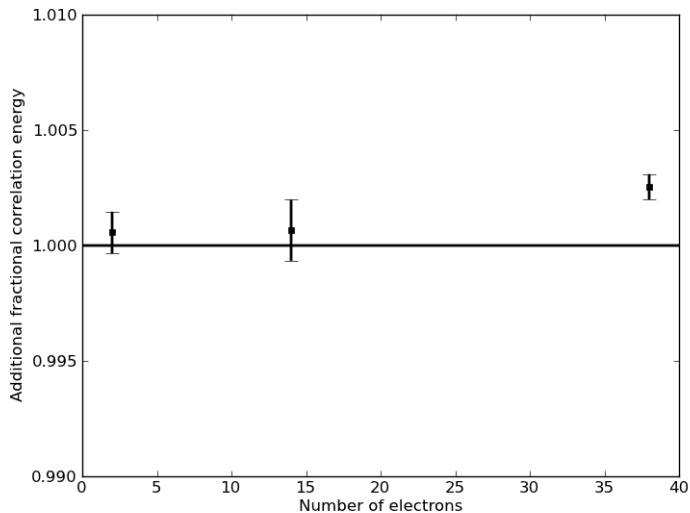
Results - Geminals alone



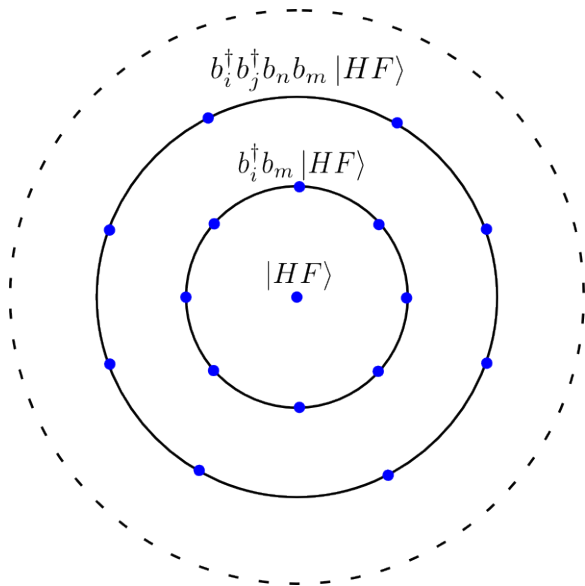
Results - Geminals in DMC



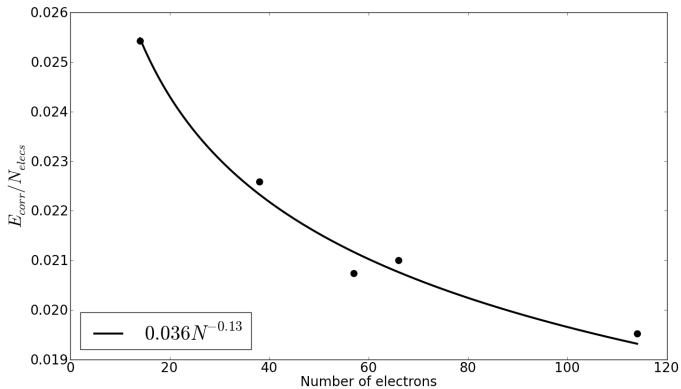
Results - Geminals in DMC



The shape of Λ^n



Multi-geminals?



Conclusions

- The nodal surface of the trial wavefunction controls the ultimate accuracy of a DMC calculation.
- Single geminals scale very poorly in bulk systems.
- Sums of geminals may be an (expensive) answer?

Geminal wavefunctions with Jastrow correlation: a first application to atoms

M. Casula and S. Sorella, J. Chem. Phys. 119, 6500 (2003)

Thank you for listening!